

Generalizing Polylogarithms to Riemann Surfaces of Arbitrary Genus

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*Knut and Alice
Wallenberg
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Organization of the Talk

1. Introduction
2. Review of polylogarithms at genus zero and one
3. A brief overview of the geometry of higher-genus Riemann surfaces
4. Construction of higher-genus polylogarithms
5. Conclusion

Introduction

Introduction

- **Polylogarithms** play a significant role in scattering amplitudes for LHC processes, SYM theory, supergravity, and string theory.
- Suitable **generalizations of classical** polylogarithms are defined by considering **iterated integrals** on closed Riemann surfaces.
- Much of the literature on polylogarithms has focused on **genus zero and genus one** Riemann surfaces, with **higher-genus surfaces** less understood.
 - Proposals for higher-genus polylogarithm function spaces exist, but without explicit **formulas** for use in physics. [Enriquez, 1112.0864]
[Enriquez, Zerbini, 2110.09341] [Enriquez, Zerbini, 2212.03119]
- Today, we will explore a **new** construction of **higher-genus polylogarithms**.
- Our method includes two key steps:
 - We create a new set of **integration kernels** using **convolutions** of certain functions defined on higher-genus Riemann surfaces.
 - With these kernels, we build a **generating function**, which helps define our **higher-genus polylogarithms** which are **closed under taking primitives**.

String amplitudes motivation

- String perturbation theory involves expanding in the **string coupling constant** g_s , which in turn is an expansion based on the **genus** of the string world-sheet.

[Figure taken from PhD thesis of J. Gerken]

$$\mathcal{A}_{\text{closed}} = g_s^{-2} \int_{\mathcal{M}_{0,4}} \text{diagram} + \int_{\mathcal{M}_{1,4}} \text{diagram} + g_s^2 \int_{\mathcal{M}_{2,4}} \text{diagram} + \dots$$

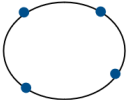
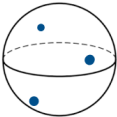
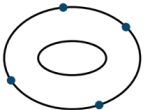
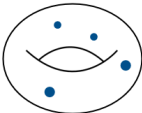
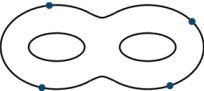
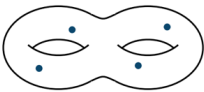
$$\mathcal{A}_{\text{open}} = g_s^{-1} \int_{\mathcal{M}_{0,4}} \text{diagram} + \int_{\mathcal{M}_{1,4}} \text{diagram} + g_s \int_{\mathcal{M}_{2,4}} \text{diagram} + \dots$$

MPL's eMPL's Higher-genus polylogs

- Furthermore, typically we also expand in the **inverse string tension** α' , which corresponds to low energy and weak coupling regimes.
- The resulting function space of these expansions is that of **polylogarithms**, (or single-valued combinations thereof.)

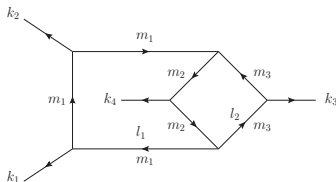
String amplitudes and special functions

- Different types of special functions emerge depending on whether we are considering **open/closed** strings, and depending on the **genus**:

	Open string	Closed string
$g = 0$	 <p>(MPL's)</p>	 <p>(<u>sv.</u> MPL's)</p>
$g = 1$	 <p>(<u>e</u>MPL's)</p>	 <p><u>e</u>MGF's (\approx <u>sv.</u> <u>e</u>MPL's)</p>
$g = 2,$ $g \geq 2$	 <p>Higher-genus polylogs (this talk)</p>	 <p>Single-valued analogues: To be explored</p>

Higher genus curves in Feynman integrals

- The appearance of **hyperelliptic curves** in Feynman integrals has also been observed in a number of publications. See for example:
- *R. Huang and Y. Zhang, "On Genera of Curves from High-loop Generalized Unitarity Cuts," JHEP 04 (2013), 080 [arXiv:1302.1023 [hep-ph]].*
- *A. Georgoudis and Y. Zhang, "Two-loop Integral Reduction from Elliptic and Hyperelliptic Curves," JHEP 12 (2015), 086 [arXiv:1507.06310 [hep-th]].*



The maximal cut of this diagram yields a hyperelliptic curve. Figure taken from [1507.06310].

- *C. F. Doran, A. Harder, E. Pichon-Pharabod and P. Vanhove, "Motivic geometry of two-loop Feynman integrals," [arXiv:2302.14840 [math.AG]].*
- *R. Marzucca, A. J. McLeod, B. Page, S. Pögel, S. Weinzierl, "Genus Drop in Hyperelliptic Feynman Integrals," [arXiv:2307.11497 [hep-th]].*

Review of polylogarithms at genus zero and one

Building Polylogarithms as Iterated Integrals

- We want to construct **polylogarithms** in terms of iterated integrals on a **compact Riemann surface, Σ , with genus h** .
- The polylogarithms we construct should have these qualities:
 1. **Homotopy Invariance**: The polylogarithms should retain their value when we smoothly change the path of integration, keeping the endpoints constant.
 2. **Logarithmic Branch-Cuts**: The integration kernels should only have simple poles, meaning our integrals should show just logarithmic irregularities at branch points.
 3. **Closed Under Integration**: Our function space should remain intact under integration, and form a basis for all iterated integrals on Σ .

Homotopy-Invariant Iterated Integrals on a Surface

- Let's consider the differential equation: $d\Gamma = \mathcal{J}\Gamma$.
- If we want the equation to be **integrable**, we need $d^2 = 0$. This leads us to the **Maurer-Cartan** equation for the connection \mathcal{J} :

$$d\mathcal{J} - \mathcal{J} \wedge \mathcal{J} = 0$$

- Such a connection is called **flat**. The solution Γ to our differential equation can be obtained by the path-ordered exponential (POE):

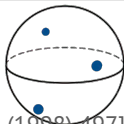
$$\Gamma(\mathcal{C}) = \text{P exp} \int_{\mathcal{C}} \mathcal{J}(\cdot) = \text{P exp} \int_0^1 dt J(t)$$

- Let's denote $\mathcal{J} = J(t)dt$, following a path \mathcal{C} where $t \in [0, 1]$, $\mathcal{C}(0) = z_0$, and $\mathcal{C}(1) = z$. Using **physics conventions**, we position $J(t)$ to the **left** of $J(t')$ for $t > t'$:

$$\text{P exp} \int_{\mathcal{C}} \mathcal{J}(\cdot) = 1 + \int_0^1 dt J(t) + \int_0^1 dt \int_0^t dt' J(t)J(t') + \dots$$

- The flatness \mathcal{J} leads to **homotopy-invariant** integrals over \mathcal{C} , (though results can differ for z_0 and z when the path circles around poles on Σ .)

Genus 0: MPLs and Generating Series



- Multiple polylogarithms (MPLs) are **iterated integrals** of rational forms $dz/(z - s)$ with $z, s \in \mathbb{C}$, on the Riemann sphere $\mathbb{C}P^1$.

[A.B. Goncharov, Math. Res. Lett. 5 (1998) 497]

- They are **defined recursively** by:

[A.B. Goncharov, math.AG/0103059]

$$G(s_1, s_2, \dots, s_n; z) = \int_0^z \frac{dt}{t - s_1} G(s_2, \dots, s_n; t)$$

where we have the special case $G(\emptyset; z) = 1$. The integer $n \geq 0$ is referred to as the **transcendental weight**.

- As a nested sum MPL's can be represented as:

$$\text{Li}_{m_1, \dots, m_k}(z_1, \dots, z_k) = \sum_{0 < n_1 < n_2 < \dots < n_k} \frac{z_1^{n_1} z_2^{n_2} \dots z_k^{n_k}}{n_1^{m_1} n_2^{m_2} \dots n_k^{m_k}}$$

- Multiple zeta values** (MZV's) are defined by the special case

$$\zeta_{m_1, \dots, m_k} \equiv \text{Li}_{m_k, \dots, m_1}(1, \dots, 1)$$

(where the ordering of the indices is subject to conventions.)

Closure of MPLs Under Integration

- Any integral of a **rational function** times a **multiple polylogarithm** (MPL) can be expressed in terms of MPLs.
- This is achieved by **partial fractioning** the rational function and/or using **integration by parts** (IBP) identities. For example:

$$\frac{1}{(x-s_1)(x-s_2)} = \frac{1}{(s_1-s_2)} \left(\frac{1}{(x-s_1)} - \frac{1}{(x-s_2)} \right)$$

- After partial fractioning, we distinguish the following cases:

$$\int_0^z dt \frac{1}{(t-b)^k} G(\vec{s}; t), \quad \int_0^z dt G(\vec{s}; t), \quad \int_0^z dt t^k G(\vec{s}; t)$$

where $0 < k \neq 1$. We then use **IBP identities** to **iteratively reduce** the value of k . For example:

$$\int_0^z dt \frac{1}{(t+1)^2} G(0; t) = \frac{z}{1+z} G(0; z) - G(-1; z)$$

Shuffle Algebra for Multiple Polylogarithms

- Multiple polylogarithms satisfy a **shuffle algebra**, which is expressed as:

$$G(s_1, s_2, \dots, s_k; z) \cdot G(s_{k+1}, \dots, s_r; z) = \sum_{\text{shuffles } \sigma} G(s_{\sigma(1)}, s_{\sigma(2)}, \dots, s_{\sigma(r)}; z),$$

- The sum runs over all permutations σ which are **shuffles** of $(1, \dots, k)$ and $(k+1, \dots, r)$.
- These permutations **preserve the relative order** of the two partitions.
- A **simple example** of the shuffle product of two multiple polylogarithms is:

$$G(s_1; z) \cdot G(s_2; z) = G(s_1, s_2; z) + G(s_2, s_1; z).$$

- The proof of the shuffle product formula relies on the integral representation of multiple polylogarithms.
- A shuffle algebra structure holds for **all the homotopy-invariant iterated integrals** which we consider.

Regularization

- Multiple polylogarithms with **trailing zeroes** do **not** have a Taylor expansion in z around $z = 0$, but **logarithmic singularities** at $z = 0$.
- We can use the shuffle product to **remove trailing zeros, separating** these logarithmic terms, such that the rest has a regular expansion around $z = 0$.
- For example, for $G(s_1, 0; z)$ with $s_1 \neq 0$, we have:

$$G(s_1, 0; z) = G(0; z) G(s_1; z) - G(0, s_1; z).$$

- Both $G(s_1; z)$ and $G(0, s_1; z)$ are **free** of trailing zeros. We then define the **special cases**:

$$G(0; z) = \log(z) \qquad G(\vec{0}_n; z) = \frac{1}{n!} \log(z)^n,$$

where $\vec{0}_n$ denotes a sequence of n zeros. These definitions follow the **tangential basepoint prescription**:

$$\int_{0+\epsilon}^x \frac{dt}{t} = \log(x) - \log(\epsilon) \rightarrow \log(x)$$

for a prescribed tangent vector (in \mathbb{C}) with $|\epsilon| \ll 1$.

Single-Valued Multiple Polylogarithms (sv. MPLs)

- Polylogarithms can be paired with their **complex conjugates** to eliminate branch cuts, thereby creating **single-valued functions**.
- For each $G(a, b, \dots; z)$, we have a corresponding $G^{\text{sv}}(a, b, \dots; z)$.
- Generally, we can **maintain the holomorphic derivative in the same form**:

$$\partial_z G^{\text{sv}}(a, b, \dots; z) = \frac{1}{z-a} G^{\text{sv}}(b, \dots; z) \quad [\text{Brown, Schnetz}]$$

and **modify the antiholomorphic** derivative to ensure single-valuedness.

- Examples: $G^{\text{sv}}(a; z) = G(a; z) + G(\bar{a}; \bar{z}) = \log \left| 1 - \frac{z}{a} \right|^2$
 $G^{\text{sv}}(0, 0, 1, 1; z) = G(0, 0, 1, 1; z) + G(0, 0, 1; z)G(1; \bar{z}) + G(0, 0; z)G(1, 1; \bar{z})$
 $\quad + G(0; z)G(1, 1, 0; \bar{z}) + G(1, 1, 0, 0; \bar{z}) + 2\zeta_3 G(1; \bar{z})$

- We can also straightforwardly define 'single-valued' MZV's:

$$\zeta_{k_1, k_2, \dots, k_r}^{\text{sv}} = \text{sv}(\zeta_{k_1, k_2, \dots, k_r}) = (-1)^r G^{\text{sv}}(\vec{0}_{k_r-1}, 1, \dots, \vec{0}_{k_1-1}, 1; 1)$$

- For instance: $\zeta_{2n}^{\text{sv}} = 0$, $\zeta_{2n+1}^{\text{sv}} = 2\zeta_{2n+1}$, $\zeta_{5,3}^{\text{sv}} = 14\zeta_3\zeta_5$.

Generating Series

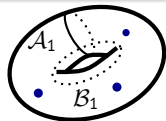
- A **generating series** for the polylogarithms can be constructed from the **Knizhnik-Zamolodchikov** (KZ) connection:

$$\mathcal{J}_{\text{KZ}}(z) = \sum_{i=1}^m \frac{dz}{z - s_i} e_i$$

- The elements e_1, \dots, e_m are generators of a free Lie algebra \mathcal{L} associated with the **marked points** s_1, \dots, s_m .
- Choosing endpoints $z_0 = 0$ and $z_1 = z$, we can **organize** the expansion of the **path-ordered exponential** in terms of the **generators** e_1, \dots, e_m :

$$\begin{aligned} \text{P exp} \int_0^z \mathcal{J}_{\text{KZ}}(\cdot) &= 1 + \sum_{i=1}^m e_i G(s_i; z) + \sum_{i=1}^m \sum_{j=1}^m e_i e_j G(s_i s_j; z) \\ &+ \sum_{i=1}^m \sum_{j=1}^m \sum_{k=1}^m e_i e_j e_k G(s_i s_j s_k; z) + \dots \end{aligned}$$

Genus 1: Elliptic Multiple Polylogarithms



- Next, consider a compact **genus-one** surface, Σ , with modulus τ , denoted as a lattice by $\Sigma = \mathbb{C}/(\mathbb{Z} + \tau\mathbb{Z})$.
 - For a surface with genus $h \geq 1$, there are multiple options for constructing a connection:
 - [Brown, Levin, arXiv:1110.6917]
 - [Broedel, Mafra, Matthes, Schlotterer, arXiv:1412.5535]
 - [Broedel, Duhr, Dulat, Tancredi, arXiv:1712.07089]
 - 1. A connection that is **single-valued** on Σ , but **non-meromorphic** (due to \bar{z} -dependence), with at most **simple poles**.
 - 2. A **meromorphic** connection that has at most **simple poles**, but is **not single-valued** (and lives on the universal cover of Σ). This can be obtained with a minor tweak of the first construction.
 - 3. A connection which is **meromorphic** and **single-valued** but has poles of **arbitrary order**. [Enriquez, Zerbini, 2110.09341]
- The **Brown-Levin construction** opts for the first choice.
 - Note that there is an **infinite set of integration kernels** at genus one, even for **a single marked point z** .

The Brown-Levin Construction

- Brown and Levin pioneered a method of **homotopy-invariant iterated integrals** at genus one. [Brown, Levin, arXiv:1110.6917]
- The key element to their construction is the so-called **Kronecker-Eisenstein (KE-) series**:

$$\Omega(z, \alpha|\tau) = \exp\left(2\pi i \alpha \frac{\operatorname{Im} z}{\operatorname{Im} \tau}\right) \frac{\vartheta_1'(0|\tau)\vartheta_1(z+\alpha|\tau)}{\vartheta_1(z|\tau)\vartheta_1(\alpha|\tau)} = \sum_{n=0}^{\infty} \alpha^{n-1} f^{(n)}(z|\tau)$$

- The KE-series is **single-valued on the torus**, has a **simple pole at $z = 0$** and satisfies the following **differential relation** (for $z \neq 0$):

$$\partial_{\bar{z}} \Omega(z, \alpha|\tau) = -\frac{\pi \alpha}{\operatorname{Im} \tau} \Omega(z, \alpha|\tau)$$

- They then constructed the **flat connection** $\mathcal{J}_{\text{BL}}(z|\tau)$, which is valued in the Lie algebra \mathcal{L} , generated by elements a, b :

$$\mathcal{J}_{\text{BL}}(z|\tau) = \frac{\pi}{\operatorname{Im} \tau} (dz - d\bar{z}) b + dz \operatorname{ad}_b \Omega(z, \operatorname{ad}_b|\tau) a$$

- Note that we have put $\alpha \rightarrow \operatorname{ad}_b = [b, \circ]$. **Flatness** can be proven using that $d_z = dz \partial_z + d\bar{z} \partial_{\bar{z}}$, and using the above differential equation.

Homotopy-Invariant Iterated Integrals

- We may write down **homotopy-invariant iterated integrals** on the torus by expanding the path-ordered exponential in terms of words in a, b :

$$\begin{aligned} \text{P exp} \int_0^z \mathcal{J}_{\text{BL}}(\cdot|\tau) &= 1 + a \Gamma(a; z|\tau) + b \Gamma(b; z|\tau) \\ &\quad + ab \Gamma(ab; z|\tau) + ba \Gamma(ba; z|\tau) + \dots \end{aligned}$$

- The resulting coefficient functions $\Gamma(\mathfrak{w}; z|\tau)$ are referred to as **elliptic polylogarithms**.
- While the connection is single-valued on the torus, the integrals are **not** and have monodromies along the \mathfrak{A} - and \mathfrak{B} -cycles.
- **Note:** In the physics literature we typically see the following functions:

$$\tilde{\Gamma}\left(\begin{matrix} n_1 & n_2 & \dots & n_r \\ w_1 & w_2 & \dots & w_r \end{matrix}; z|\tau\right) = \int_0^z dz_1 g^{(n_1)}(z_1 - w_1|\tau) \tilde{\Gamma}\left(\begin{matrix} n_2 & \dots & n_r \\ w_2 & \dots & w_r \end{matrix}; z_1|\tau\right)$$

- We review the connection of these functions to the above construction next.

Meromorphic Variant

- We can define a **meromorphic counterpart** of the doubly-periodic Kronecker-Eisenstein series and its expansion coefficients $g^{(n)}(z|\tau)$:

$$\frac{\vartheta_1'(0|\tau)\vartheta_1(z+\alpha|\tau)}{\vartheta_1(z|\tau)\vartheta_1(\alpha|\tau)} = \sum_{n=0}^{\infty} \alpha^{n-1} g^{(n)}(z|\tau)$$

- The meromorphic integration kernels $g^{(n)}(z|\tau)$ are **multiple-valued on the torus**, and actually **live on the universal covering space**, which is \mathbb{C} .
- Brown-Levin polylogarithms associated with **words** $w \rightarrow ab \cdots b$ **reduce to a single integral over the meromorphic kernels**. For example:

$$\Gamma(ab; z|\tau) = \int_0^z dt \left(2\pi i \frac{\operatorname{Im} t}{\operatorname{Im} \tau} - f^{(1)}(t|\tau) \right) = - \int_0^z dt g^{(1)}(t|\tau) = -\tilde{\Gamma}\left(\frac{1}{0}; z|\tau\right)$$

- **More generally**, $\Gamma(\underbrace{ab \cdots b}_n; z|\tau)$ can be expressed as:

$$\Gamma(\underbrace{ab \cdots b}_n; z|\tau) = (-1)^n \int_0^z dt g^{(n)}(t|\tau) = (-1)^n \tilde{\Gamma}\left(\frac{n}{0}; z|\tau\right)$$

Closure under integration

- For the MPLs, we saw that partial fraction identities were essential for splitting up a product of integration kernels.
- We need similar identities for the **function space to close under integration** at genus one. For example, we might encounter an integral of the type:

$$\int_0^z dt f^{(n_1)}(t - a_1) f^{(n_2)}(t - a_2)$$

[Broedel, Mafra, Matthes, Schlotterer, arXiv:1412.5535]

- The so-called **Fay identities** generalize the partial fraction relations. They are generated by:

$$\begin{aligned} \Omega(z_1, \alpha_1, \tau) \Omega(z_2, \alpha_2, \tau) &= \Omega(z_1, \alpha_1 + \alpha_2, \tau) \Omega(z_2 - z_1, \alpha_2, \tau) \\ &\quad + \Omega(z_2, \alpha_1 + \alpha_2, \tau) \Omega(z_1 - z_2, \alpha_1, \tau) \end{aligned}$$

- For example we have that:

$$\begin{aligned} f^{(1)}(t - x) f^{(1)}(t) &= f^{(1)}(t - x) f^{(1)}(x) - f^{(1)}(t) f^{(1)}(x) \\ &\quad + f^{(2)}(t) + f^{(2)}(x) + f^{(2)}(t - x) \end{aligned}$$

Alternative Construction via Convolutions

- An **alternative construction** of the functions $f^{(k)}(z|\tau)$ is in terms of the **scalar Green function** $g(z|\tau)$ on Σ . The Green function is defined by:

$$\partial_{\bar{z}}\partial_z g(z|\tau) = -\pi\delta(z) + \frac{\pi}{\text{Im } \tau}, \quad \int_{\Sigma} d^2z g(z|\tau) = 0$$

- It can be expressed in terms of the Jacobi theta function ϑ_1 and the Dedekind eta-function η as follows:

$$g(z|\tau) = -\ln \left| \frac{\vartheta_1(z|\tau)}{\eta(\tau)} \right|^2 - \pi \frac{(z-\bar{z})^2}{2 \text{Im } \tau}$$

- We define the function $f^{(1)}(z|\tau)$ as the derivative of the Green's function:

$$f^{(1)}(z|\tau) = -\partial_z g(z|\tau)$$

- Subsequently, we can define **higher dimensional convolutions** of f recursively as follows:

$$f^{(k)}(z|\tau) = - \int_{\Sigma} \frac{d^2x}{\text{Im } \tau} \partial_x g(x|\tau) f^{(k-1)}(x-z|\tau), \quad k \geq 2$$

- We will see in the following that **similar convolutions underlie** our higher-genus generalizations of these kernels.

Modular Properties of the Brown-Levin Construction

- Let us consider the **modular properties** of the Brown-Levin construction.
- We take a modular transformation on the modulus τ , z , and α :

$$\tau \rightarrow \tilde{\tau} = \frac{A\tau + B}{C\tau + D}, \quad z \rightarrow \tilde{z} = \frac{z}{C\tau + D}, \quad \alpha \rightarrow \tilde{\alpha} = \frac{\alpha}{C\tau + D}$$

where $A, B, C, D \in \mathbb{Z}$ with $AD - BC = 1$.

- The Kronecker-Eisenstein series Ω and the functions $f^{(n)}$ transform as **modular forms of weight $(1, 0)$ and $(n, 0)$** , respectively:

$$\Omega(\tilde{z}, \tilde{\alpha}|\tilde{\tau}) = (C\tau + D)\Omega(z, \alpha|\tau), \quad f^{(n)}(\tilde{z}|\tilde{\tau}) = (C\tau + D)^n f^{(n)}(z|\tau)$$

- The connection \mathcal{J}_{BL} can be made **modular invariant** by assigning the following transformation to the generators a, b :

$$a \rightarrow \tilde{a} = (C\tau + D)a + 2\pi i C b, \quad b \rightarrow \tilde{b} = \frac{b}{C\tau + D}$$

- The **extra contribution** $2\pi i C b$ to \tilde{a} is engineered so that:

$$\frac{\pi d\tilde{z}}{\text{Im } \tilde{\tau}} \tilde{b} = \frac{C\tilde{\tau} + D}{C\tau + D} \frac{\pi dz}{\text{Im } \tau} b$$

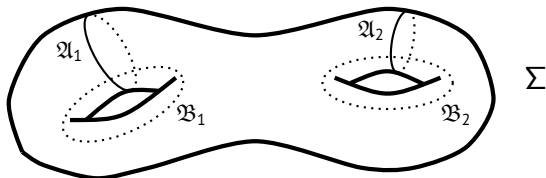
Summary: The Brown-Levin construction

Step	Brown-Levin construction	Higher-genus construction
1. Integration kernels	$f^{(k)}(z \tau) = - \int_{\Sigma} \frac{d^2x}{\text{Im } \tau} \partial_x g(x \tau) f^{(k-1)}(x-z \tau)$	$\begin{aligned} \Phi^{l_1 \cdots l_r}_J(x) &= \int_{\Sigma} d^2z \mathcal{G}(x, z) \bar{\omega}^{l_1}(z) \partial_z \Phi^{l_2 \cdots l_r}_J(z) \quad (r \geq 2) \\ \mathcal{G}^{l_1 \cdots l_s}(x, y) &= \int_{\Sigma} d^2z \mathcal{G}(x, z) \bar{\omega}^{l_1}(z) \partial_z \mathcal{G}^{l_2 \cdots l_s}(z, y) \quad (s \geq 1) \end{aligned}$
2. Generating series	$\alpha \Omega(z, \alpha \tau) = \sum_{n=0}^{\infty} \alpha^n f^{(n)}(z \tau)$	$\begin{aligned} \Psi_J(x, \rho; B) &= \omega_J(x) + \left(\partial_x \Phi^{l_1}_J(x) - \partial_x \mathcal{G}(x, \rho) \delta^{l_1}_J \right) B_{l_1} \\ &+ \sum_{l=2}^{\infty} \left(\partial_x \Phi^{l_1 l_2 \cdots l_r}_J(x) - \partial_x \mathcal{G}^{l_1 l_2 \cdots l_r-1}(x, \rho) \delta^{l_r}_J \right) \\ &\times B_{l_1} B_{l_2} \cdots B_{l_r} \end{aligned}$
3. Flat connection ($d\mathcal{J} - \mathcal{J} \wedge \mathcal{J} = 0$)	$\begin{aligned} \mathcal{J}_{\text{BL}}(x \tau) &= -d\bar{x} b \\ &+ \frac{\pi}{\text{Im } \tau} dx b + dx \text{ad}_b \Omega(x, \text{ad}_b \tau) a \end{aligned}$	$\begin{aligned} \mathcal{J}(x, \rho) &= -\pi d\bar{x} \bar{\omega}^l(x) b_l \\ &+ \pi dx \mathcal{H}^l(x; B) b_l + dx \Psi_J(x, \rho; B) d^l \end{aligned}$
4. Path-ordered exponential	$\begin{aligned} \text{P exp} \int_0^x \mathcal{J}_{\text{BL}}(\cdot \tau) &= \\ &1 + a \Gamma(a; x \tau) + b \Gamma(b; x \tau) \\ &+ ab \Gamma(ab; x \tau) + ba \Gamma(ba; x \tau) + \dots \end{aligned}$	$\begin{aligned} \text{P exp} \int_{\gamma}^x \mathcal{J}(t, \rho) &= \\ &1 + d^l \Gamma_l(x, y; \rho) + b_l \Gamma^l(x, y; \rho) \\ &+ d^l d^l \Gamma_{ll}(x, y; \rho) + b_l b_l \Gamma^{ll}(x, y; \rho) \\ &+ d^l b_l \Gamma_l^l(x, y; \rho) + b_l d^l \Gamma^l_l(x, y; \rho) + \dots \end{aligned}$
5. Polylogs	<p>e.g.</p> $\Gamma(ab; x \tau) = \int_0^x dt \left(2\pi i \frac{\text{Im } t}{\text{Im } \tau} - f^{(1)}(t \tau) \right)$	<p>e.g.</p> $\begin{aligned} \Gamma^l(x, y; \rho) &= \pi \int_{\gamma}^x \left(\partial_t (\partial_t \Phi^l_K(t) \gamma^{Kl} - \partial_t \Phi^l_K(t) \gamma^{Kl}) \right. \\ &\left. + \pi (\omega^l(t) - \bar{\omega}^l(t)) \int_{\gamma}^t (\omega^l - \bar{\omega}^l) \right) \end{aligned}$

Brief overview of higher-genus Riemann surfaces

Topology of a Compact Riemann Surface Σ

- The **topology** of a **compact** Riemann surface Σ without boundary is specified by its **genus h** .
- The **homology group** $H_1(\Sigma, \mathbb{Z})$ is isomorphic to \mathbb{Z}^{2h} and supports an **anti-symmetric non-degenerate intersection pairing** denoted by \mathfrak{J} .



A choice of canonical homology basis on a compact **genus-two** Riemann surface Σ .

- A **canonical homology basis** of cycles \mathfrak{A}_l and \mathfrak{B}_l with $l, j = 1, \dots, h$ has symplectic intersection matrix $\mathfrak{J}(\mathfrak{A}_l, \mathfrak{B}_j) = -\mathfrak{J}(\mathfrak{B}_j, \mathfrak{A}_l) = \delta_{lj}$, and $\mathfrak{J}(\mathfrak{A}_l, \mathfrak{A}_j) = \mathfrak{J}(\mathfrak{B}_l, \mathfrak{B}_j) = 0$.

Canonical Basis of Holomorphic Abelian Differentials

- A **canonical basis** of **holomorphic Abelian differentials** ω_j may be normalized on \mathfrak{A} -cycles:

$$\oint_{\mathfrak{A}_I} \omega_j = \delta_{IJ} \quad \oint_{\mathfrak{B}_I} \omega_j = \Omega_{IJ}$$

- The complex variables Ω_{IJ} denote the components of the **period matrix** Ω of the surface Σ .
- By the **Riemann relations**, Ω is **symmetric**, and has **positive definite imaginary part**:

$$\Omega^t = \Omega \quad Y = \text{Im } \Omega > 0$$

- We will use the matrix $Y_{IJ} = \text{Im } \Omega_{IJ}$ and its **inverse** $Y^{IJ} = ((\text{Im } \Omega)^{-1})^{IJ}$ to **raise and lower** indices:

$$\omega^I = Y^{IJ} \omega_j \quad \bar{\omega}^I = Y^{IJ} \bar{\omega}_j \quad Y^{IK} Y_{KJ} = \delta^I_J$$

Modular Transformations and Modular Tensors

- A **new canonical basis** $\tilde{\mathfrak{X}}$ and $\tilde{\mathfrak{Y}}$ is obtained by applying a **modular transformation** $M \in Sp(2h, \mathbb{Z})$, such that $M^t \mathfrak{Y} M = \mathfrak{X}$.
- Under a modular transformation, we have:

$$\tilde{\omega} = \omega (C\Omega + D)^{-1}, \quad \tilde{\Omega} = (A\Omega + B)(C\Omega + D)^{-1}$$

$$\tilde{Y} = (\tilde{\Omega} C^t + D^t)^{-1} Y (C\Omega + D)^{-1}$$

- **Modular tensors**, generalize modular forms, replacing $(C\tau + D)$ of $SL(2, \mathbb{Z})$ with $Q = C\Omega + D$ and $R = (C\Omega + D)^{-1}$. It holds that:

$$\tilde{\omega}_I = \omega_{I'} R^{I'}_I \qquad \tilde{Y}_{IJ} = Y_{I'J'} \bar{R}^{I'}_I R^{J'}_J$$

$$\tilde{\omega}^J = \bar{Q}^{J'}_J \omega^{J'} \qquad \tilde{Y}^{IJ} = Q^{I'}_{I'} \bar{Q}^{J'}_{J'} Y^{I'J'}$$

- A modular tensor \mathcal{T} of arbitrary rank **transforms** as follows:

$$\tilde{\mathcal{T}}^{I_1, \dots, I_n; J_1, \dots, J_n}(\tilde{\Omega}) = Q^{I_1}_{I'_1} \cdots Q^{I_n}_{I'_n} \bar{Q}^{J_1}_{J'_1} \cdots \bar{Q}^{J_n}_{J'_n} \mathcal{T}^{I'_1, \dots, I'_n; J'_1, \dots, J'_n}(\Omega)$$

The Arakelov Green Function

- The **Arakelov Green function** $\mathcal{G}(x, y|\Omega)$ on $\Sigma \times \Sigma$ is a **single-valued version** of the Green function, defined by: [D'Hoker, Green, Pioline, arXiv:1712.06135]
[G. Faltings, Ann. Math., 119(2), 1984]

$$\partial_{\bar{x}}\partial_x\mathcal{G}(x, y|\Omega) = -\pi\delta(x, y) + \pi\kappa(x), \quad \int_{\Sigma} \kappa(x)\mathcal{G}(x, y|\Omega) = 0$$

where the **Kähler form** κ is given by:

$$\kappa = \frac{i}{2h}\omega_I \wedge \bar{\omega}^I = \kappa(z) d^2z \quad \int_{\Sigma} \kappa = 1$$

- The Arakelov Green function also obeys the following derivatives:

$$\partial_x\partial_y\mathcal{G}(x, y) = -\partial_x\partial_y \ln E(x, y) + \pi\omega_I(x)\omega^I(y)$$

$$\partial_x\partial_{\bar{y}}\mathcal{G}(x, y) = \pi\delta(x, y) - \pi\omega_I(x)\bar{\omega}^I(y)$$

- The prime form $E(x, y)$ is a unique form that is **holomorphic** in x and y and **vanishes linearly** as x approaches y .
- In what follows we will not write the explicit dependence on the moduli Ω .

The Arakelov Green Function

- An **explicit formula** for $\mathcal{G}(x, y)$ may be given in terms of the non-conformally invariant string Green function $G(x, y)$:

$$\mathcal{G}(x, y) = G(x, y) - \gamma(x) - \gamma(y) + \gamma_0$$

- The **string Green function** is given in terms of the **prime form** $E(x, y)$ by:

$$G(x, y) = -\log |E(x, y)|^2 + 2\pi \left(\operatorname{Im} \int_y^x \omega_l \right) \left(\operatorname{Im} \int_y^x \omega'_l \right)$$

- The functions $\gamma(x)$ and γ_0 are given by:

$$\gamma(x) = \int_{\Sigma} \kappa(z) G(x, z) \quad \gamma_0 = \int_{\Sigma} \kappa \gamma$$

- Both κ and $\mathcal{G}(x, y)$ are **conformally invariant**.
- At genus one the (Arakelov) Green function only depends on a difference of points $\mathcal{G}(x, y)|_{h=1} = \mathcal{G}(x - y)|_{h=1}$.
- However, this **translation invariance** is **absent** on a Riemann surface Σ of genus $h > 1$.

The Interchange Lemma

- The tensor $\Phi^I_J(x)$, introduced by Kawazumi, compensates for the lack of translation invariance at higher genus: [Kawazumi, MCM2016] [Kawazumi, 2017]

$$\Phi^I_J(x) = \int_{\Sigma} d^2z \mathcal{G}(x, z) \bar{\omega}^I(z) \omega_J(z)$$

- Note that the **trace** of $\Phi^I_J(x)$ **vanishes** by the definition of the Arakelov Green function.
- In particular, the so-called **interchange lemma** provides a substitute for the absence of translation invariance:

$$\partial_x \mathcal{G}(x, y) \omega_J(y) + \partial_y \mathcal{G}(x, y) \omega_J(x) - \partial_x \Phi^I_J(x) \omega_I(y) - \partial_y \Phi^I_J(y) \omega_I(x) = 0$$

[E. D'Hoker, C. R. Mafra, B. Pioline, O. Schlotterer, arXiv:2008.08687]

Construction of higher-genus polylogarithms

Higher Convolution of the Arakelov Green Function

- Inspired by the alternative construction of the Kronecker-Eisenstein kernels through convolutions, we define the **tensors** $\Phi^{l_1 \cdots l_r}_J(x)$ and $\mathcal{G}^{l_1 \cdots l_s}(x, y)$:

$$\Phi^{l_1 \cdots l_r}_J(x) = \int_{\Sigma} d^2z \mathcal{G}(x, z) \bar{\omega}^{l_1}(z) \partial_z \Phi^{l_2 \cdots l_r}_J(z) \quad (r \geq 2)$$

$$\mathcal{G}^{l_1 \cdots l_s}(x, y) = \int_{\Sigma} d^2z \mathcal{G}(x, z) \bar{\omega}^{l_1}(z) \partial_z \mathcal{G}^{l_2 \cdots l_s}(z, y) \quad (s \geq 1)$$

- (We also encounter these tensors while decomposing cyclic products of Szegő kernels, see [D'Hoker, MH, Schlotterer, arXiv:2308.05044]).
- At genus one**, the derivatives of the tensor $\mathcal{G}^{l_1 \cdots l_s}$ for $l_1 = \cdots = l_s = 1$ equal the Kronecker-Eisenstein integration kernels $f^{(s+1)}$:

$$\partial_x \mathcal{G}^{l_1 \cdots l_s}(x, y) \Big|_{h=1} = -f^{(s+1)}(x-y|\tau)$$

- The trace $\Phi^{l_1 \cdots l_r}_J = 0$ for arbitrary genus implies that Φ -tensors for arbitrary $r \geq 1$ **vanish** identically for **genus one**.
- In the next part:** we will construct **generating functions** of our kernels, and combine them into a flat connection.

Generating Functions

- Let us introduce a **non-commutative algebra freely generated by B_l** for $l = 1, \dots, h$ (loosely inspired by the approach of Enriquez and Zerbini arXiv:2110.09341).
- Next, we fix an arbitrary **auxiliary marked point p** on the Riemann surface Σ and introduce the following **generating functions**:

$$\mathcal{H}(x, p; B) = \partial_x \mathcal{G}(x, p) + \sum_{r=1}^{\infty} \partial_x \mathcal{G}^{l_1 l_2 \dots l_r}(x, p) B_{l_1} B_{l_2} \cdots B_{l_r}$$

$$\mathcal{H}_J(x; B) = \omega_J(x) + \sum_{r=1}^{\infty} \partial_x \Phi^{l_1 l_2 \dots l_r}_J(x) B_{l_1} B_{l_2} \cdots B_{l_r}$$

- By forming the **combination $\Psi_J(x, p; B) = \mathcal{H}_J(x; B) - \mathcal{H}(x, p; B)B_J$** , we obtain a compact antiholomorphic derivative:

$$\partial_{\bar{x}} \Psi_J(x, p; B) = -\pi \bar{\omega}^l(x) B_l \Psi_J(x, p; B)$$

for $x \neq p$, which **generalizes the genus-one differential relation for Ω** .

The Flat Connection

- Next, we **extend** to a Lie algebra \mathcal{L} **freely generated** by elements a^l and b_l for $l = 1, \dots, h$ and set $B_l = \text{ad}_{b_l} = [b_l, \cdot]$.
- Our **connection** $\mathcal{J}(x, p)$, on a **Riemann surface** Σ of arbitrary **genus** h with a **marked point** $p \in \Sigma$ and valued in the **Lie algebra** \mathcal{L} is then given by:

$$\mathcal{J}(x, p) = -\pi d\bar{x} \bar{\omega}^l(x) b_l + \pi dx \mathcal{H}^l(x; B) b_l + dx \Psi_l(x, p; B) a^l$$

- Working out $d_x = dx \partial_x + d\bar{x} \partial_{\bar{x}}$, we may show that:

$$d_x \mathcal{J}(x, p) - \mathcal{J}(x, p) \wedge \mathcal{J}(x, p) = \pi d\bar{x} \wedge dx \delta(x, p) [b_l, a^l]$$

proving that the connection is **flat** (away from $x = p$).

- At **genus one**, $\mathcal{J}(x, p)$ **reduces** to the **Brown-Levin connection**, upon relabeling $a^1 = a$ and $b_1 = b$. In particular:

$$\Psi_1(x, p; B) \Big|_{h=1} = \text{ad}_b \Omega(x-p, \text{ad}_b | \tau)$$

Expansion of the Connection

- The connection \mathcal{J} may be **expanded in words** in the basis (a^l, b_l) :

$$\begin{aligned}\mathcal{J}(x, p) &= \pi(dx \omega^l(x) - d\bar{x} \bar{\omega}^l(x))b_l + \pi dx \sum_{r=1}^{\infty} \partial_x \Phi^{l_1 \cdots l_r}(x) Y^{JK} B_{l_1} \cdots B_{l_r} b_K \\ &+ dx \sum_{r=1}^{\infty} \left(\partial_x \Phi^{l_1 \cdots l_r}(x) - \partial_x \mathcal{G}^{l_1 \cdots l_{r-1}}(x, p) \delta_J^{l_r} \right) B_{l_1} \cdots B_{l_r} a^J\end{aligned}$$

- Like before, the flat connection $\mathcal{J}(x, p)$ **integrates** to a homotopy-invariant path-ordered exponential $\Gamma(x, y; p)$:

$$\Gamma(x, y; p) = \text{P exp} \int_y^x \mathcal{J}(t, p)$$

- For example, for words with at most two letters in the basis (a^l, b_l) :

$$\begin{aligned}\Gamma(x, y; p) &= 1 + a^l \Gamma_l(x, y; p) + b_l \Gamma^l(x, y; p) \\ &+ a^l a^J \Gamma_{JJ}(x, y; p) + b_l b_J \Gamma^{JJ}(x, y; p) \\ &+ a^l b_J \Gamma_l^J(x, y; p) + b_l a^J \Gamma^J_l(x, y; p) + \cdots\end{aligned}$$

Summary: Construction of higher-genus polylogs

Step	Brown-Levin construction	Higher-genus construction
1. Integration kernels	$f^{(k)}(z \tau) =$ $- \int_{\Sigma} \frac{d^2x}{\text{Im}\tau} \partial_x g(x \tau) f^{(k-1)}(x-z \tau)$	$\Phi^{l_1 \cdots l_r}_J(x) =$ $\int_{\Sigma} d^2z \mathcal{G}(x, z) \tilde{\omega}^{l_1}(z) \partial_z \Phi^{l_2 \cdots l_r}_J(z) \quad (r \geq 2)$ $\mathcal{G}^{l_1 \cdots l_s}(x, y) =$ $\int_{\Sigma} d^2z \mathcal{G}(x, z) \tilde{\omega}^{l_1}(z) \partial_z \mathcal{G}^{l_2 \cdots l_s}(z, y) \quad (s \geq 1)$
2. Generating series	$\alpha \Omega(z, \alpha \tau) = \sum_{n=0}^{\infty} \alpha^n f^{(n)}(z \tau)$	$\Psi_J(x, \rho; B) = \omega_J(x) + \left(\partial_x \Phi^{l_1}_J(x) - \partial_x \mathcal{G}(x, \rho) \delta^{l_1}_J \right) B_{l_1}$ $+ \sum_{r=2}^{\infty} \left(\partial_x \Phi^{l_1 l_2 \cdots l_r}_J(x) - \partial_x \mathcal{G}^{l_1 l_2 \cdots l_r-1}(x, \rho) \delta^{l_r}_J \right)$ $\times B_{l_1} B_{l_2} \cdots B_{l_r}$
3. Flat connection	$\mathcal{J}_{\text{BL}}(x \tau) = -d\bar{x} b$ $+ \frac{\pi}{\text{Im}\tau} dx b + dx \text{ad}_b \Omega(x, \text{ad}_b \tau) a$	$\mathcal{J}(x, \rho) = -\pi d\bar{x} \tilde{\omega}^l(x) b_l$ $+ \pi dx \mathcal{H}^l(x; B) b_l + dx \Psi_J(x, \rho; B) d^l$
4. Path-ordered exponential	$\text{P exp} \int_0^x \mathcal{J}_{\text{BL}}(\cdot \tau) =$ $1 + a \Gamma(a; x \tau) + b \Gamma(b; x \tau)$ $+ ab \Gamma(ab; x \tau) + ba \Gamma(ba; x \tau) + \dots$	$\text{P exp} \int_Y^X \mathcal{J}(t, \rho) =$ $1 + d^l \Gamma_l(x, y; \rho) + b_l \Gamma^l(x, y; \rho)$ $+ d^l d^l \Gamma_{ll}(x, y; \rho) + b_l b_l \Gamma^{ll}(x, y; \rho)$ $+ d^l b_l \Gamma_l^l(x, y; \rho) + b_l d^l \Gamma^l_l(x, y; \rho) + \dots$
5. Polylogs	<p>e.g.</p> $\Gamma(ab; x \tau) =$ $\int_0^x dt \left(2\pi i \frac{\text{Im}t}{\text{Im}\tau} - f^{(1)}(t \tau) \right)$	<p>e.g.</p> $\Gamma^l(x, y; \rho) = \pi \int_Y^X \left(dt (\partial_t \Phi^l_K(t) y^{Kl} - \partial_t \Phi^l_K(t) y^{Kl}) \right.$ $\left. + \pi (\omega^l(t) - \tilde{\omega}^l(t)) \int_Y^t (\omega^l - \tilde{\omega}^l) \right)$

Polylogarithms for Words without b_l

- The polylogarithms associated with words w that do not involve any of the letters b_l are given by the following simple formula:

$$\Gamma_{l_1 l_2 \dots l_r}(x, y; p) = \int_y^x \omega_{l_1}(t_1) \int_y^{t_1} \omega_{l_2}(t_2) \cdots \int_y^{t_{r-1}} \omega_{l_r}(t_r)$$

which we'll refer to as iterated Abelian integrals.

- These polylogarithms are **independent of the marked point p** .
- They obey the differential equations:

$$\partial_x \Gamma_{l_1 l_2 \dots l_r}(x, y; p) = \omega_{l_1}(x) \Gamma_{l_2 \dots l_r}(x, y; p)$$

- For the case $h = 1$, we simply obtain:

$$\Gamma_{\underbrace{11 \dots 1}_r}(x, y; z) \Big|_{h=1} = \frac{1}{r!} (x-y)^r$$

Low Letter Count Polylogarithms

- Next let us consider some cases involving the letters b_l . For the **single-letter word** b_l , we obtain:

$$\Gamma^l(x, y; p) = \pi \int_y^x (\omega^l - \bar{\omega}^l)$$

- For **double-letter words** with **at least one letter** b_l , we obtain:

$$\Gamma^{ll}(x, y; p) = \pi \int_y^x \left(dt (\partial_t \Phi^l_{Kl}(t) Y^{Kl} - \partial_t \Phi^l_{Kl}(t) Y^{Kl}) + \pi (\omega^l(t) - \bar{\omega}^l(t)) \int_y^t (\omega^l - \bar{\omega}^l) \right)$$

$$\Gamma^l_l(x, y; p) = \int_y^x \left(dt \partial_t \Phi^l_l(t) - dt \partial_t \mathcal{G}(t, p) \delta^l_l + \pi (\omega^l(t) - \bar{\omega}^l(t)) \int_y^t \omega_l \right)$$

$$\Gamma^l_l(x, y; p) = \int_y^x \left(-dt \partial_t \Phi^l_l(t) + dt \partial_t \mathcal{G}(t, p) \delta^l_l + \pi \omega_l(t) \int_y^t (\omega^l - \bar{\omega}^l) \right)$$

Meromorphic Variants of Polylogarithms

- Lastly, let's explore an instance showcasing where the **meromorphic variants** of polylogarithms live in our function space.
- Consider again the following higher-genus polylogarithm:

$$\Gamma_I^J(x, y; p) = \int_y^x dt \left(-\partial_t \Phi_I^J(t) + \delta_I^J \partial_t \mathcal{G}(t, p) + \pi \omega_I(t) Y^{JK} (\Gamma_K(t, y; p) - \overline{\Gamma_K(t, y; p)}) \right)$$

- Upon specializing to genus $h = 1$ and setting $p = y = 0$, this reproduces the Brown-Levin polylogarithm $\Gamma(ab; p|\tau) = -\tilde{\Gamma}\left(\frac{1}{0}; p|\tau\right)$.
- The integrand with respect to t in the equation above can be viewed as a **higher-genus uplift** of the Kronecker-Eisenstein kernel $g^{(1)}(t|\tau)$:

$$g_I^J(t, y; p) = \partial_t \Phi_I^J(t) - \delta_I^J \partial_t \mathcal{G}(t, p) - 2\pi i \omega_I(t) Y^{JK} \operatorname{Im} \int_y^t \omega_K$$

- One may verify that indeed (for $t \neq p$):

$$\partial_t g_I^J(t, y; p) = 0$$

Modular Invariance and Hatted Basis

- To investigate modular properties, let us define an **alternative basis** (\hat{a}^I, b_I) of generators of the Lie algebra \mathcal{L} :

$$\hat{a}^I = a^I + \pi Y^I b_I$$

- In this basis, the connection $\mathcal{J}(x, p)$ takes on a **simplified form**:

$$\mathcal{J}(x, p) = -\pi d\bar{x} \bar{\omega}^I(x) b_I + dx \Psi_I(x, p; B) \hat{a}^I$$

- A **modular transformation** $M \in Sp(2h, \mathbb{Z})$, acts on $\bar{\omega}^I$, B_I , \mathcal{H}_I , and Ψ_I , and on the Lie algebra generators a^I and b_I by:

$$a^I \rightarrow \tilde{a}^I = Q^I_J a^J + 2\pi i C^I_J b_J$$

$$b_I \rightarrow \tilde{b}_I = b_J R^J_I$$

- Then also

$$\hat{a}^I \rightarrow \tilde{\hat{a}}^I = Q^I_J \hat{a}^J$$

- The connection $\mathcal{J}(x, p)$ is seen to be **manifestly invariant** under $Sp(2h, \mathbb{Z})$.

Polylogarithms In The Hatted Basis

- In the basis (\hat{a}^I, b_I) , the expansion is given by:

$$\begin{aligned}\Gamma(x, y; p) &= 1 + \hat{a}^I \hat{\Gamma}_I(x, y; p) + b_I \hat{\Gamma}^I(x, y; p) \\ &\quad + \hat{a}^I \hat{a}^J \hat{\Gamma}_{IJ}(x, y; p) + b_I b_J \hat{\Gamma}^{IJ}(x, y; p) \\ &\quad + \hat{a}^I b_J \hat{\Gamma}_I^J(x, y; p) + b_I \hat{a}^J \hat{\Gamma}_J^I(x, y; p) + \dots\end{aligned}$$

- The polylogarithms $\hat{\Gamma}(x, y; p)$ in the basis (\hat{a}^I, b_I) are **modular tensors** by the $Sp(2h, \mathbb{Z})$ **invariance** of the connection $\mathcal{J}(x, p)$.

$$\tilde{\Gamma}^{\dots I \dots \dots J \dots}(x, y; p) = \dots R^I{}_I \dots Q^J{}_{J'} \dots \hat{\Gamma}^{\dots I' \dots \dots J' \dots}(x, y; p)$$

- Identifying term by term in both expansions gives the relations $\Gamma_I = \hat{\Gamma}_I$ and $\Gamma_{IJ} = \hat{\Gamma}_{IJ}$, as well as the following relations:

$$\hat{\Gamma}^I = \Gamma^I - \pi Y^{IJ} \Gamma_J$$

$$\hat{\Gamma}^{IJ} = \Gamma^{IJ} - \pi Y^{IK} \Gamma_{KJ}$$

$$\hat{\Gamma}_I^J = \Gamma_I^J - \pi \Gamma_{IK} Y^{KJ}$$

$$\hat{\Gamma}^{IJ} = \Gamma^{IJ} - \pi Y^{IK} \Gamma_{KJ} - \pi \Gamma_K^I Y^{KJ} + \pi^2 Y^{IK} \Gamma_{KL} Y^{LJ}$$

Low Letter Count Polylogarithms in the Hatted Basis

- Let us write the expansion of the generating function $\Psi_J(x, p; B)$ in the following way:

$$\Psi_J(x, p; B) = \omega_J(x) + \sum_{r=1}^{\infty} B_{l_1} \cdots B_{l_r} f^{l_1 \cdots l_r}_J(x, p)$$
$$f^{l_1 \cdots l_r}_J(x, p) = \partial_x \Phi^{l_1 \cdots l_r}_J(x) - \partial_x \mathcal{G}^{l_1 \cdots l_{r-1}}(x, p) \delta_J^{l_r}$$

- The polylogarithms for **one- and two-letter words, starting with b_I** , are:

$$\hat{\Gamma}^I(x, y; p) = -\pi \int_y^x \bar{\omega}^I = -\pi Y^{IK} \overline{\Gamma_K(x, y; p)}$$

$$\hat{\Gamma}^{IJ}(x, y; p) = \pi^2 \int_y^x \bar{\omega}^I(t_1) \int_y^{t_1} \bar{\omega}^J = \pi^2 Y^{IK} Y^{JL} \overline{\Gamma_{KL}(x, y; p)}$$

$$\hat{\Gamma}_I^J(x, y; p) = -\int_y^x dt \left(f^J_I(t, p) + \pi \omega_I(t) \int_y^t \bar{\omega}^J \right)$$

$$\hat{\Gamma}_J^I(x, y; p) = \int_y^x dt \left(f^I_J(t, p) + \pi \omega_J(t) \int_y^t \bar{\omega}^I \right) - \pi Y^{IK} \overline{\Gamma_K(x, y; p)} \Gamma_J(x, y; p)$$

- The expressions are more **compact** compared to the previous case.

Conclusion

Conclusion

- We have presented an explicit construction of **polylogarithms** on **higher-genus** compact Riemann surfaces.
- Our construction relies on a **flat connection** whose **path-ordered exponential** plays the role of a **generating series** for **higher-genus polylogarithms**.
- The flat connection takes values in the **freely-generated Lie algebra generated by elements a^l and b_l** for $l = 1, \dots, h$, introduced by Enriquez and Zerbini.
- Although we have strong evidence the function space of our polylogarithms is closed under integration, we have not yet proven this conjecture.
- Our construction provides the first **explicit proposal** for a **complete** set of **integration kernels beyond genus one**.

Future Directions

1. Obtaining the **separating and non-separating degenerations** of the polylogarithms for arbitrary genera.
2. Determining the **differential relations with respect to moduli variations** satisfied by higher-genus polylogarithms.
3. Identifying generalizations of the **higher-genus modular graph tensors** that close under complex-structure variations and degenerations.
4. **Re-formulation** of higher-genus string amplitudes in terms of the integration kernels and polylogarithms constructed in this work.

Thank you for listening!

Backup Slides

Simplified Representations

- The polylogarithms with **upper indices** admit **simplified representations** in terms of the **iterated abelian integrals**, their **complex conjugates** and **contractions with Y^J** .
- For words with a **single letter b_I** we have:

$$\Gamma^I(x, y; p) = \pi Y^J (\Gamma_J(x, y; p) - \overline{\Gamma_J(x, y; p)})$$

- For **two-letter words that contain at least one b_I** , we have:

$$\Gamma_I^J(x, y; p) = \pi Y^{JK} \Gamma_{IK}(x, y; p) + \int_y^x dt \left(-\partial_t \Phi^J_I(t) + \delta^J_I \partial_t \mathcal{G}(t, p) - \pi \omega_I(t) Y^{JK} \overline{\Gamma_K(t, y; p)} \right)$$

$$\begin{aligned} \Gamma^I_J(x, y; p) &= \pi Y^{IK} (\Gamma_{KJ}(x, y; p) - \Gamma_J(x, y; p) \overline{\Gamma_K(x, y; p)}) \\ &\quad + \int_y^x dt \left(\partial_t \Phi^I_J(t) - \delta^I_J \partial_t \mathcal{G}(t, p) + \pi \omega_J(t) Y^{IK} \overline{\Gamma_K(t, y; p)} \right) \end{aligned}$$

$$\begin{aligned} \Gamma^J(x, y; p) &= \pi^2 Y^{IK} Y^{JL} \left(\Gamma_{KL}(x, y; p) + \overline{\Gamma_{KL}(x, y; p)} - \overline{\Gamma_K(x, y; p)} \Gamma_L(x, y; p) \right) \\ &\quad + \pi \int_y^x dt \left(\partial_t \Phi^I_K(t) Y^{KJ} - \partial_t \Phi^J_K(t) Y^{KI} \right. \\ &\quad \left. + \pi \omega^J(t) Y^{IK} \overline{\Gamma_K(t, y; p)} - \pi \omega^I(t) Y^{JK} \overline{\Gamma_K(t, y; p)} \right) \end{aligned}$$