

Spin- s Rational Q -system

based on: 2303.07640v2



Jue Hou¹, Yunfeng Jiang¹, Rui-Dong Zhu²

¹School of physics & Shing-Tung Yau Center, Southeast University, Nanjing 211189, P. R. China

²Institute for Advanced Study & School of Physical Science and Technology, Soochow University, Suzhou 215006, P. R. China

Abstract

Bethe ansatz equations for spin- s XXX spin chain (XXX _{s} chain) with $s \geq 1$ are significantly more difficult to analyze than the spin- $\frac{1}{2}$ case, due to the presence of repeated roots. As a result, it is challenging to derive extra conditions for the Bethe roots to be physical and study the related completeness problem.

In our paper[1], we propose the rational Q -system for the XXX _{s} chain, which gives all and only physical solutions of the Bethe ansatz equations(BAEs), and the extra conditions for solutions to be physical. Then we prove the completeness of spin- s Bethe ansatz rigorously.

I Introduction

The BAEs for the XXX _{s} chain of length L with M magnons read

$$\left(\frac{u_j + is}{u_j - is}\right)^L = \prod_{k \neq j}^M \frac{u_j - u_k + i}{u_j - u_k - i}, \quad j = 1, \dots, M. \quad (1)$$

There are two types of solutions which need special care. One is the solutions with *repeated roots* and the other are the so-called *singular solutions*. The repeated solutions can take the following form

$$\{\underbrace{u_1, \dots, u_1}_{n_1}, \underbrace{u_2, \dots, u_2}_{n_2}, \dots, \underbrace{u_M, \dots, u_M}_{n_M}\} \quad (2)$$

A generic singular solution can take the following form

$$\{(is)^{m_0}, (i(s-1))^{m_1}, \dots, (-is)^{m_{2s}}, u_1, u_2, \dots, u_{M_r}\} \quad (3)$$

where $\{u_1, \dots, u_{M_r}\}$ are the regular roots. Not all of them are physical, which can be derived by algebraic Bethe ansatz. To get physical solutions, we need other constraints besides BAEs.

The TQ -relation method can eliminate non-physical repeated solutions automatically. Baxter's TQ -relation reads

$$\tau(u)Q(u) = (u + \frac{i}{2})^L Q(u - i) + (u - \frac{i}{2})^L Q(u + i), \quad (4)$$

where $\tau(u)$, $Q(u)$ are polynomials. The zeros of $Q(u)$ are the Bethe roots $\{u_k\}$. The polynomiality of $\tau(u)$ imposes extra constraints for the Bethe roots. The Q -system method can eliminate non-physical singular solutions automatically.

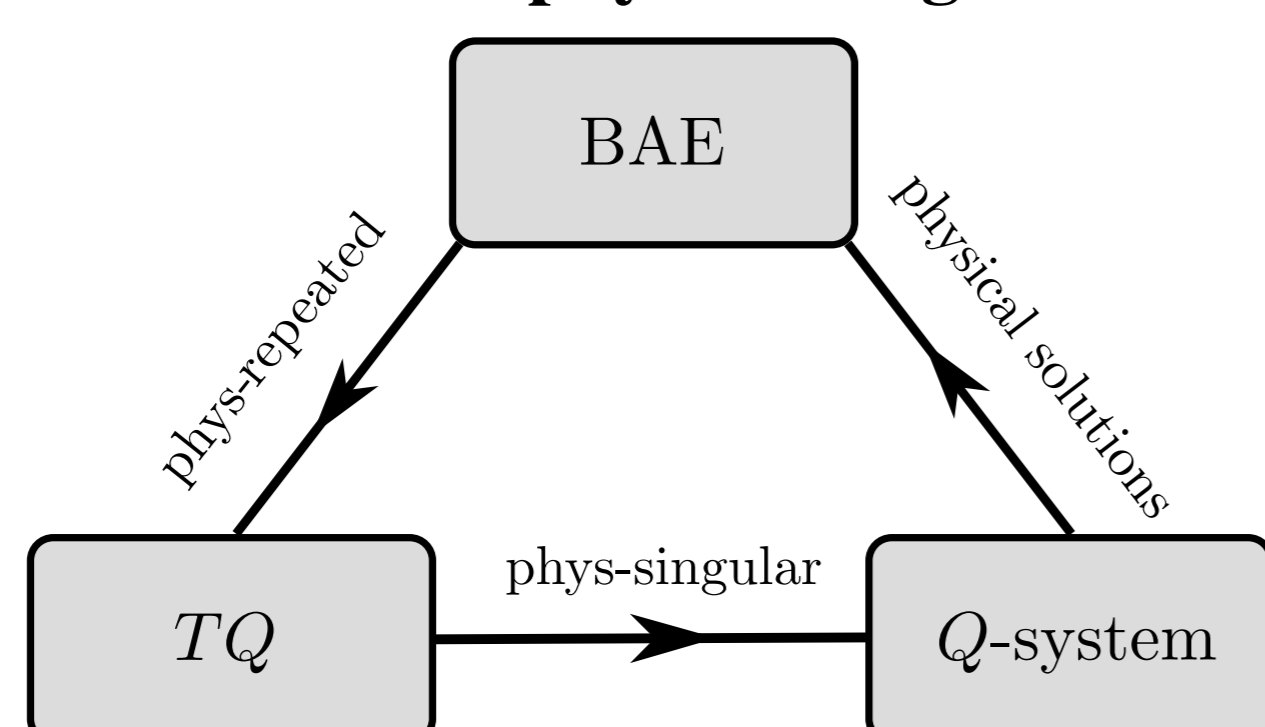


Figure 1: Relations between BAE, TQ and rational Q -system.

II Spin- s rational Q -system

To construct the rational Q -system, we need the following ingredients

1. A Young diagram on which we define Q -functions;
2. QQ -relations which relates the four Q -functions defined on each box;
3. Boundary conditions which specify the Q -functions partially or completely on the left and upper boundary of the Young diagram.

Young diagram The Young diagram has two rows and is given by $(2sL - M, M)$. On each node, we define a function $Q_{a,b}(u)$,

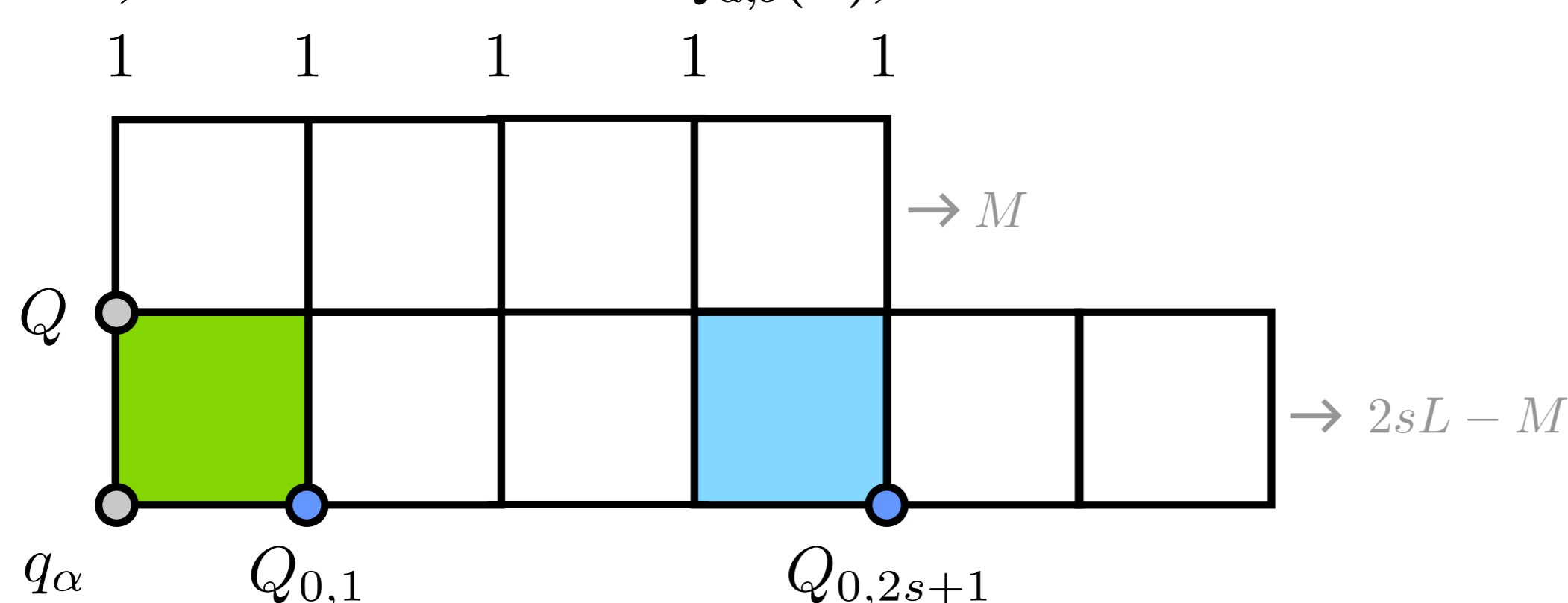


Figure 2: Rational Q -system for XXX _{s} spin chain.

QQ -relations

$$Q_{a,b+1}Q_{a+1,b} = Q_{a,b}^+ Q_{a+1,b+1}^- - Q_{a,b}^- Q_{a+1,b+1}^+. \quad (5)$$

Boundary conditions

$$\begin{aligned} Q_{2,b} &= 1, & b &= 0, 1, \dots, M, \\ Q_{1,0} &= Q(u), \\ Q_{0,0} &= q_\alpha(u). \end{aligned} \quad (6)$$

Solve the Q -system We make the usual ansatz

$$Q(u) = u^M + \sum_{k=0}^{M-1} c_k u^k. \quad (7)$$

We then require $Q_{0,1}/\alpha(u)$ and $Q_{0,2s+1}$ to be polynomials.

Here, $F^\pm(u) \equiv F(u \pm \frac{i}{2})$, $q_\alpha(u) \equiv \prod_{k=0}^{2s-1} (u - i(s-k-\frac{1}{2}))^L$, and $\alpha(u) \equiv \prod_{k=1}^{2s-1} (u - i(s-k))^L$. c_k s are undetermined constants and we will obtain them by solving the Q -system.

III Completeness of spin- s Bethe ansatz

All physical Bethe vectors span the Hilbert space. The expected number of physical solutions for a spin- s L -site chain with M magnons is given by[2]

$$\mathcal{N}_s(L, M) = c_s(L, M) - c_s(L, M - 1), \quad (8)$$

with

$$c_s(L; M) = \sum_{j=0}^{\lfloor \frac{L+M-1}{2s+1} \rfloor} (-1)^j \binom{L}{j} \binom{L+M-1-(2s+1)j}{L-1}. \quad (9)$$

Numerical Checks.

	$\mathcal{N}_{\text{phys}}$		$\mathcal{N}_{\text{phys}}$
$L = 3, M = 1, s = 1$	2	$L = 8, M = 4, s = \frac{3}{2}$	202
$L = 3, M = 2, s = 1$	3	$L = 8, M = 6, s = \frac{3}{2}$	700
$L = 3, M = 3, s = 1$	1	$L = 9, M = 4, s = \frac{3}{2}$	321
$L = 6, M = 4, s = 1$	40	$L = 10, M = 4, s = \frac{3}{2}$	485
$L = 6, M = 5, s = 1$	36	$L = 6, M = 5, s = 2$	120
$L = 6, M = 6, s = 1$	15	$L = 6, M = 6, s = 2$	180
$L = 9, M = 1, s = 1$	8	$L = 7, M = 5, s = 2$	245
$L = 9, M = 2, s = 1$	36	$L = 7, M = 6, s = 2$	420
$L = 9, M = 3, s = 1$	111		
$L = 9, M = 4, s = 1$	258		
$L = 12, M = 3, s = 1$	274		
$L = 12, M = 4, s = 1$	869		

Table 1: $\mathcal{N}_{\text{phys}}$ denotes the number of all physical solutions obtained from the Q -system of a spin- s L -site chain with M magnons.

Rigorously Proof.

In [3], authors prove if the difference equation (4) about $Q(u)$ has two polynomial solutions $Q(u)$, $P(u)$ such that $\deg(Q) < \deg(P)$, then the completeness of spin- s Bethe ansatz is valid. In our paper [1], we prove for XXX _{s} chain, the above premise is valid.

Please read our paper for more details.

IV Conclusions

1. We propose the rational Q -system for spin- s XXX spin chain, which gives all and only physical solutions. Furthermore, the additional constraints for singular and repeated solutions are derived (see [1]).
2. The completeness of spin- s Bethe ansatz is tested numerically and proved rigorously.

References

- [1] J. Hou, Y. Jiang and R. D. Zhu, Spin- s Rational Q -system, [arXiv:2303.07640].
- [2] A. N. Kirillov, Combinatorial identities, and completeness of eigenstates of the Heisenberg magnet, *Journal of Soviet Mathematics* 30 no. 4, (1985) 2298–2310.
- [3] V. Tarasov, Completeness of the Bethe ansatz for the periodic isotropic Heisenberg model, *Rev. Math. Phys.* 30 (2018) 1840018.