

Structure Constants in $\mathcal{N} = 4$ SYM and Separation of Variables

Based on 2210.04923 with C. Bercini, P. Vieira

Alexandre Homrich

IGST 2023

Context

Context

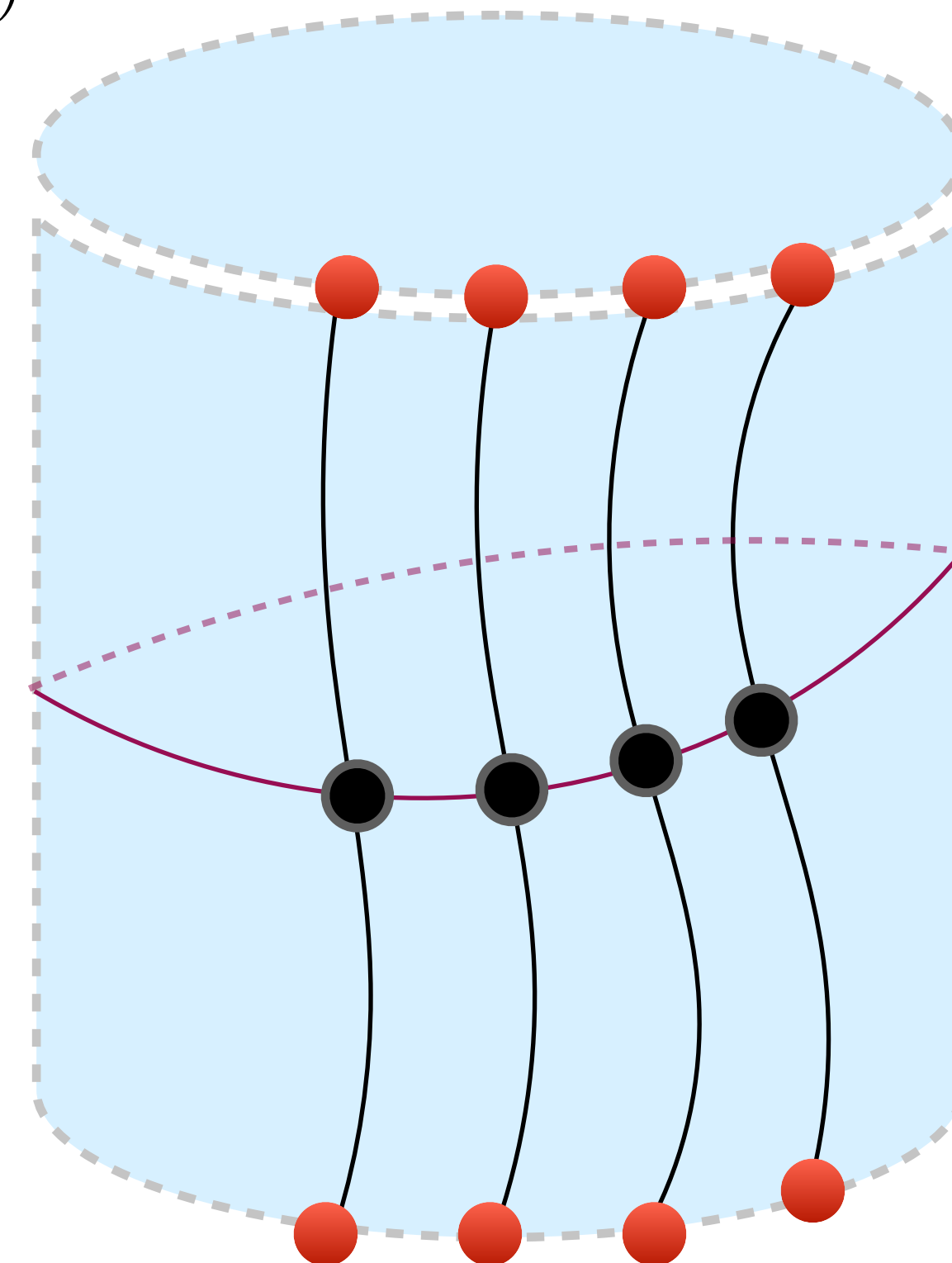
- How to compute observables in QFT from integrability?

Context

- How to compute observables in QFT from integrability?
 - ↳ Reduce computation to generalized 2d scattering problem and bootstrap it.

Context

- How to compute observables in QFT from integrability?
 - ↳ Reduce computation to generalized 2d scattering problem and bootstrap it.
- Classical example in $\mathcal{N} = 4$ SYM: 2-pt functions (spectrum)



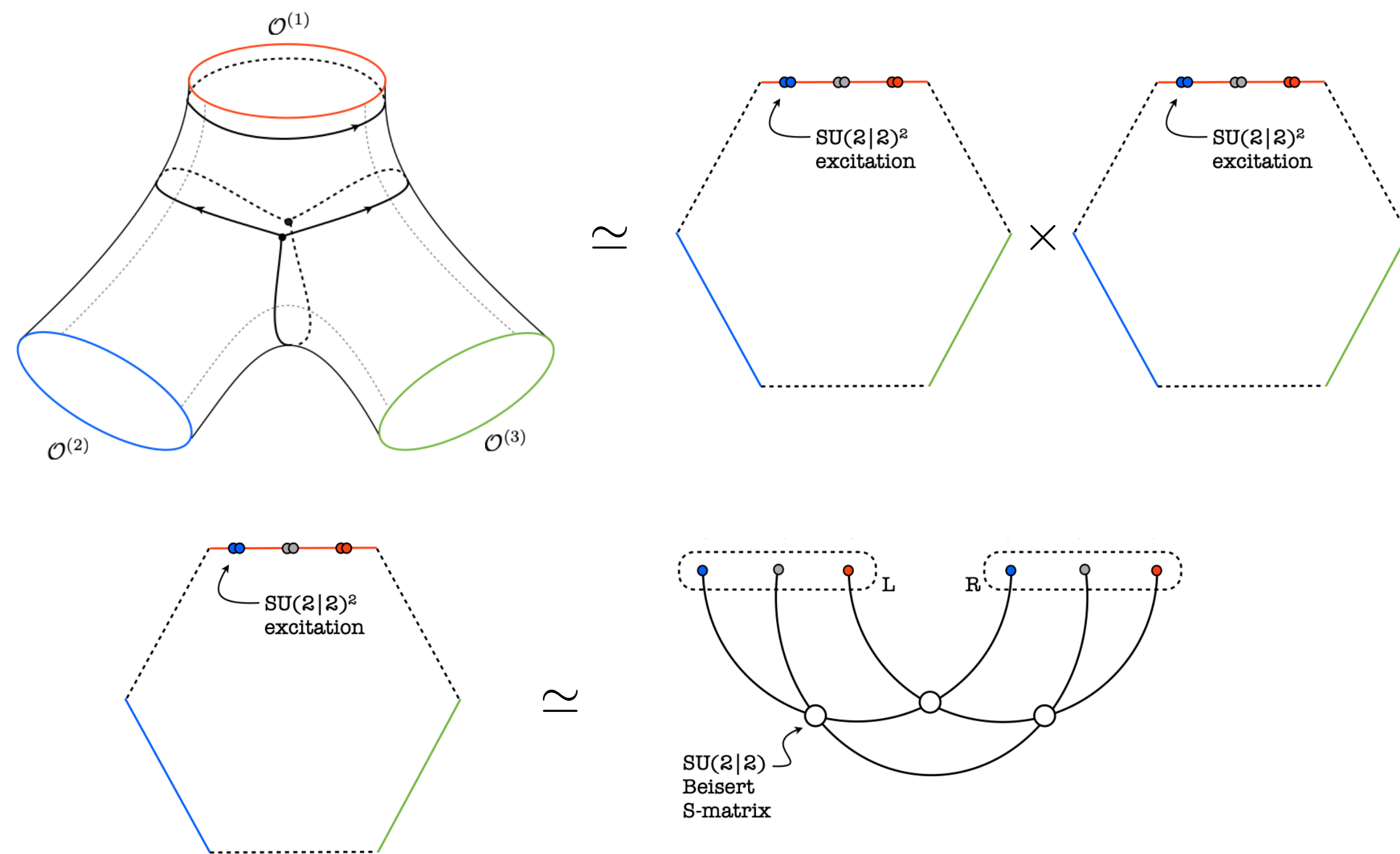
[Minahan, Zarembo], [Beisert], ..., [Gromov et. al - Arutyunov et. al. - Bombardelli et. al.], ...

Context

- How to compute observables in QFT from integrability?
 - ↳ Reduce computation to generalized 2d scattering problem and bootstrap it.

- Classical example in $\mathcal{N} = 4$ SYM: 2-pt functions (spectrum)

- Since then: \rightarrow n-point functions



[Basso, Komatsu, Vieira][Komatsu, Fleury]

Context

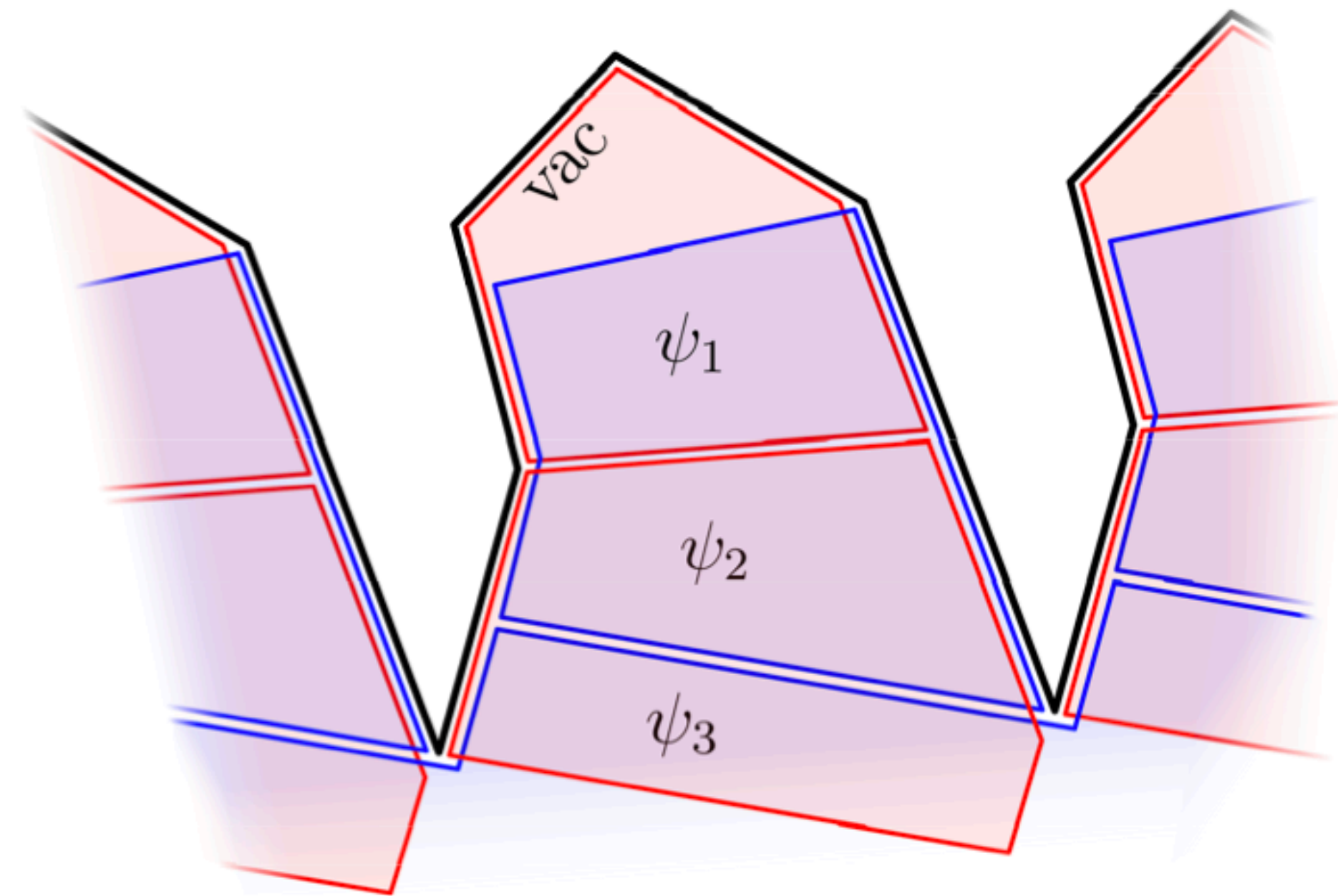
- How to compute observables in QFT from integrability?
 - ↳ Reduce computation to generalized 2d scattering problem and bootstrap it.
- Classical example in $\mathcal{N} = 4$ SYM: 2-pt functions (spectrum)
- Since then: → n-point functions
→ amplitudes

$$\mathcal{W}_{\text{hexagon}} = \text{vacuum} + e^{-E\tau} \text{[1 wavy line]} + e^{-(E_1+E_2)\tau} \text{[2 wavy lines]} + \dots$$

[Basso, Sever, Vieira]

Context

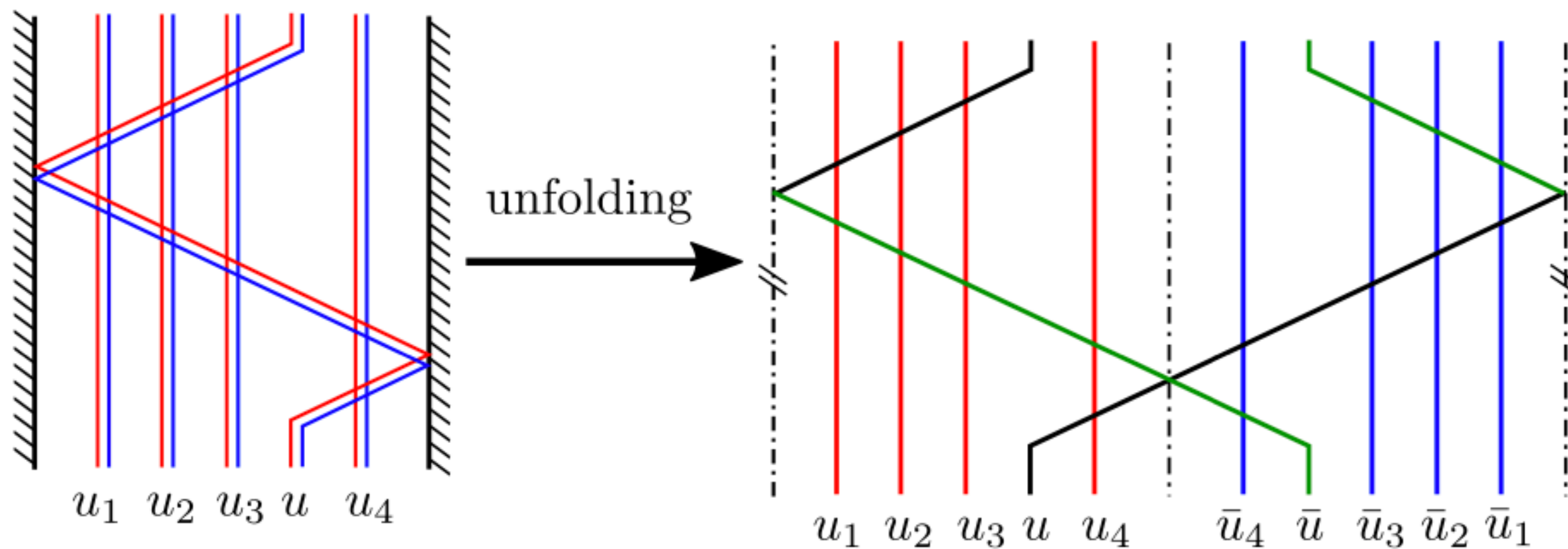
- How to compute observables in QFT from integrability?
 - ↳ Reduce computation to generalized 2d scattering problem and bootstrap it.
- Classical example in $\mathcal{N} = 4$ SYM: 2-pt functions (spectrum)
- Since then: → n-point functions
→ amplitudes
→ form factors



[Sever, Tumanov, Wilhelm]

Context

- How to compute observables in QFT from integrability?
 - ↳ Reduce computation to generalized 2d scattering problem and bootstrap it.
- Classical example in $\mathcal{N} = 4$ SYM: 2-pt functions (spectrum)
- Since then: → n-point functions
→ amplitudes
→ form factors
→ determinants
...



[Jiang, Komatsu, Vescovi]

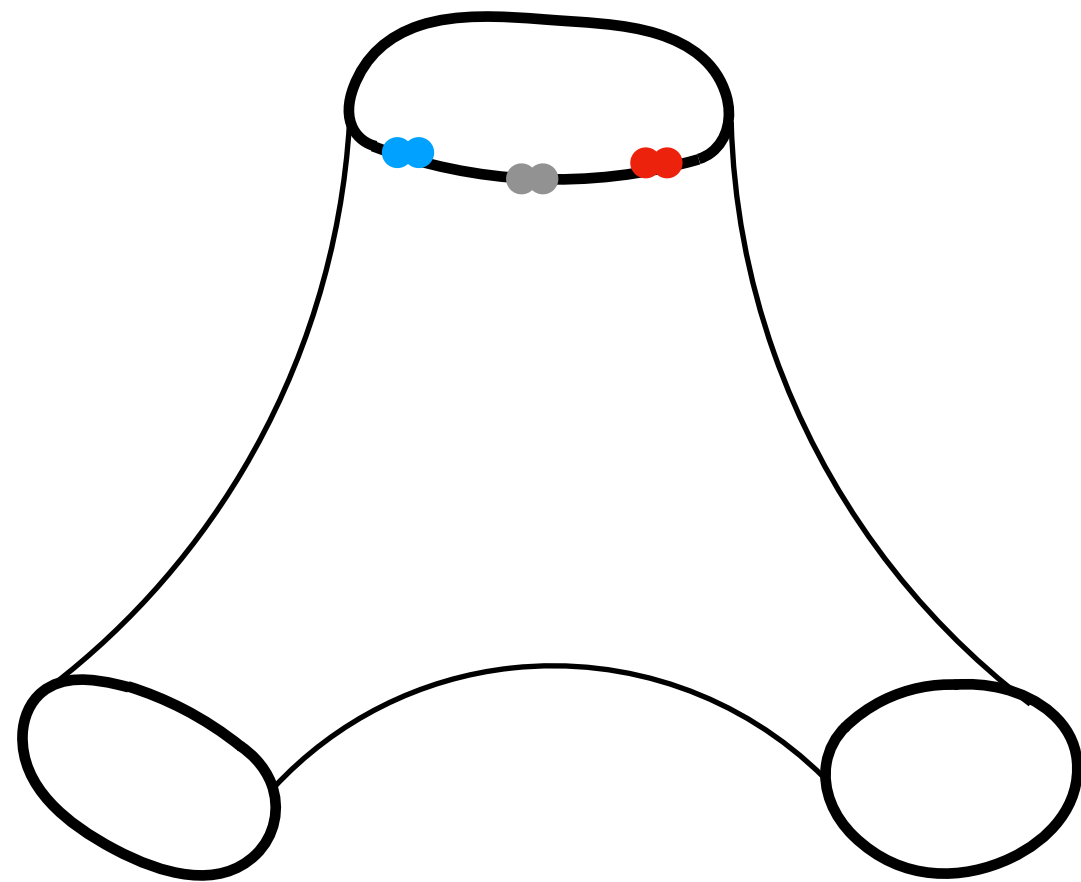
Context

Context

- Overall strategy:

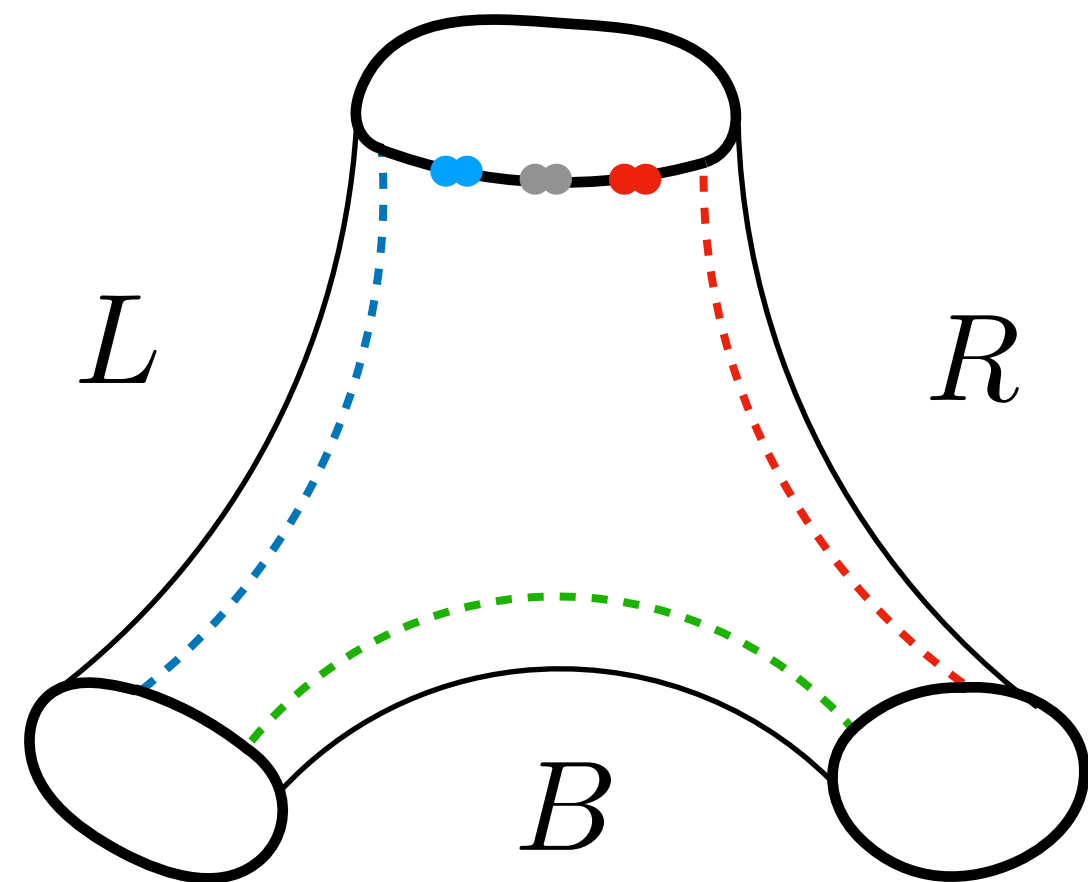
Context

- Overall strategy:
 - ↳ Start with stringy picture of observable



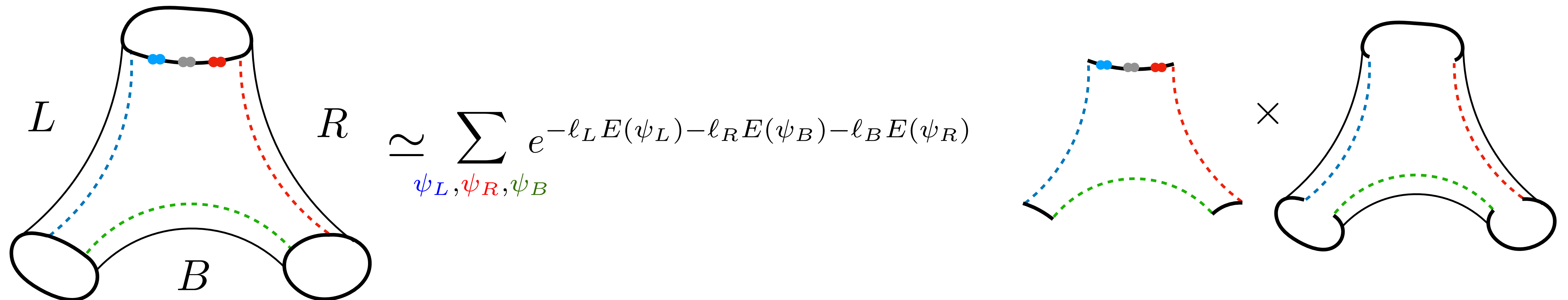
Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion



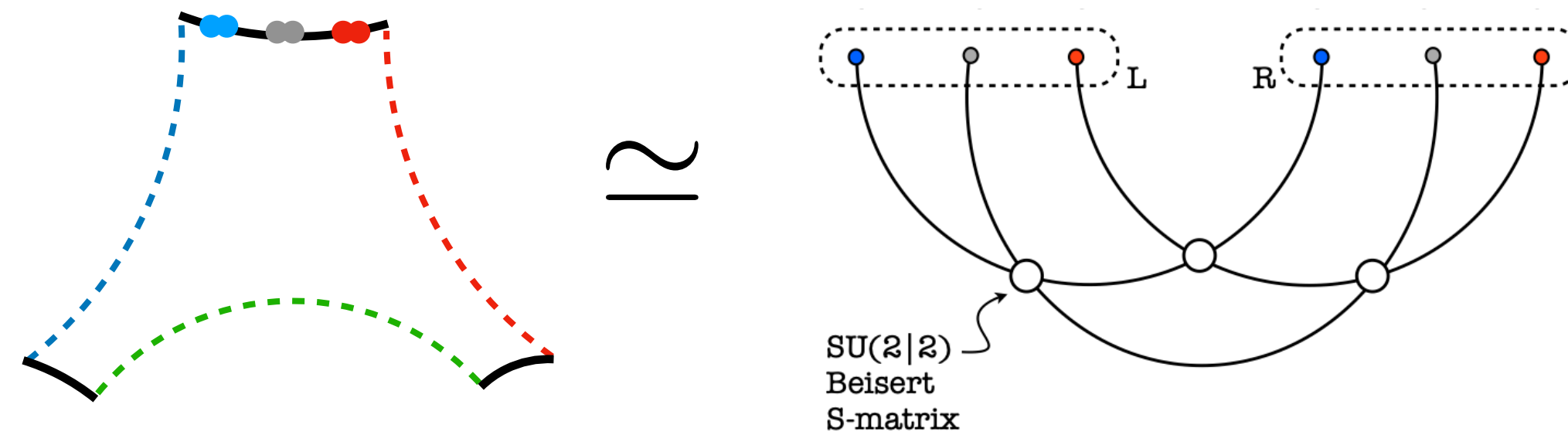
Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion



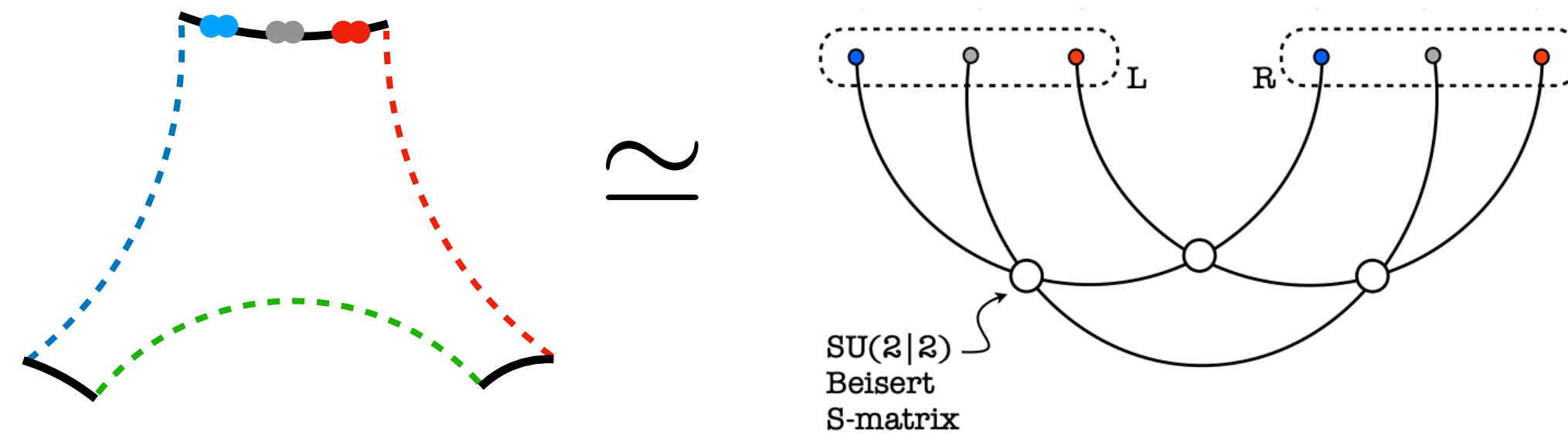
Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion
 - ↳ Each open patch is a 2D scattering problem solved by bootstrap



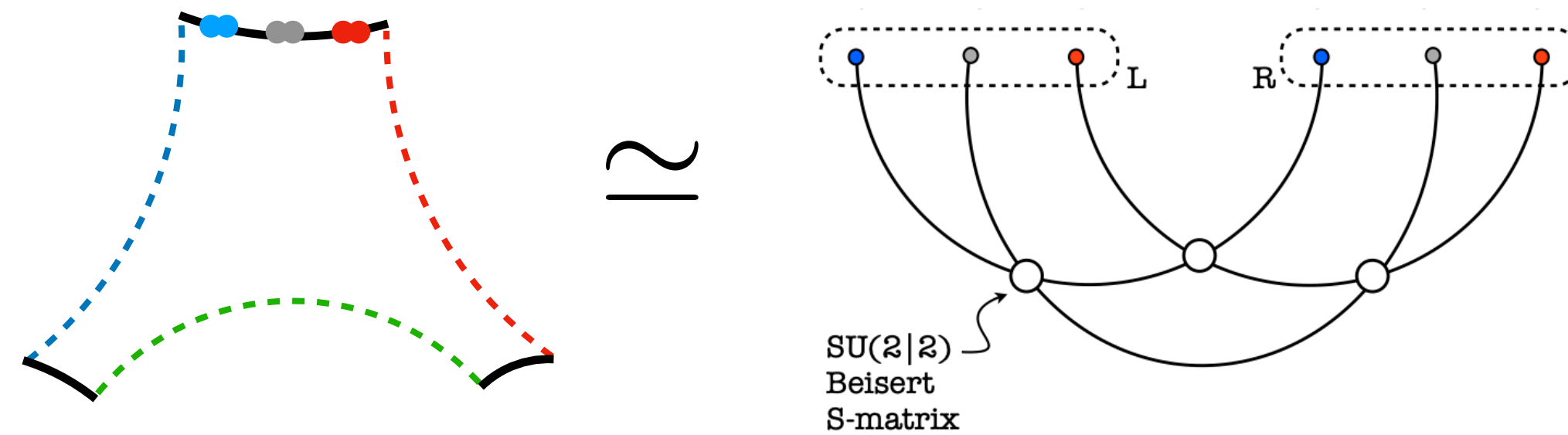
Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion
 - ↳ Each open patch is a 2D scattering problem solved by bootstrap
- Very physical framework.



Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion
 - ↳ Each open patch is a 2D scattering problem solved by bootstrap
- Very physical framework.
 - ↳ Connects large N gauge theory to the worldsheet cartoon



Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion
 - ↳ Each open patch is a 2D scattering problem solved by bootstrap
- Very physical framework.
 - ↳ Connects large N gauge theory to the worldsheet cartoon

Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion
 - ↳ Each open patch is a 2D scattering problem solved by bootstrap
- Very physical framework.
 - ↳ Connects large N gauge theory to the worldsheet cartoon
- Evaluating and resumming large geometry expansion is too hard in practice

Context

- Overall strategy:
 - ↳ Start with stringy picture of observable
 - ↳ Cut worldsheet open \Leftrightarrow Large geometry/low temperature expansion
 - ↳ Each open patch is a 2D scattering problem solved by bootstrap
- Very physical framework.
 - ↳ Connects large N gauge theory to the worldsheet cartoon
- Evaluating and resumming large geometry expansion is too hard in practice
 - ↳ Need truncating limits: weak coupling, large charge, collinear kinematics...

Context

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
[Gromov, Kazakov, Leurent, Volin]

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables
 - ↳ Evidence that this is promising has been piling up recently
[Cavaglià, Gromov, Levkovich-Maslyuk][Caetano, Komatsu][Giombi, Komatsu][Giombi, Komatsu, Offertaler]...

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables
 - ↳ Evidence that this is promising has been piling up recently
[Cavaglià, Gromov, Levkovich-Maslyuk][Caetano, Komatsu][Giombi, Komatsu][Giombi, Komatsu, Offertaler]...
 - ↳ Number of technical developments serve as a backbone to these explorations
[Cavaglià, Gromov, Levkovich-Maslyuk]²[Gromov, Levkovich-Maslyuk, Ryan][Gromov, Levkovich-Maslyuk, Ryan, Volin]

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables
 - ↳ Evidence that this is promising has been piling up recently
[Cavaglià, Gromov, Levkovich-Maslyuk][Caetano, Komatsu][Giombi, Komatsu][Giombi, Komatsu, Offertaler]...
 - ↳ Number of technical developments serve as a backbone to these explorations
[Cavaglià, Gromov, Levkovich-Maslyuk]²[Gromov, Levkovich-Maslyuk, Ryan][Gromov, Levkovich-Maslyuk, Ryan, Volin]
- In this talk I will present:

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables
 - ↳ Evidence that this is promising has been piling up recently
[Cavaglià, Gromov, Levkovich-Maslyuk][Caetano, Komatsu][Giombi, Komatsu][Giombi, Komatsu, Offertaler]...
 - ↳ Number of technical developments serve as a backbone to these explorations
[Cavaglià, Gromov, Levkovich-Maslyuk]²[Gromov, Levkovich-Maslyuk, Ryan][Gromov, Levkovich-Maslyuk, Ryan, Volin]
- In this talk I will present:
 - ↳ Structure constants formulas at weak coupling (up to N^2LO) based on Q-functions and the “SoV” formalism.

Context

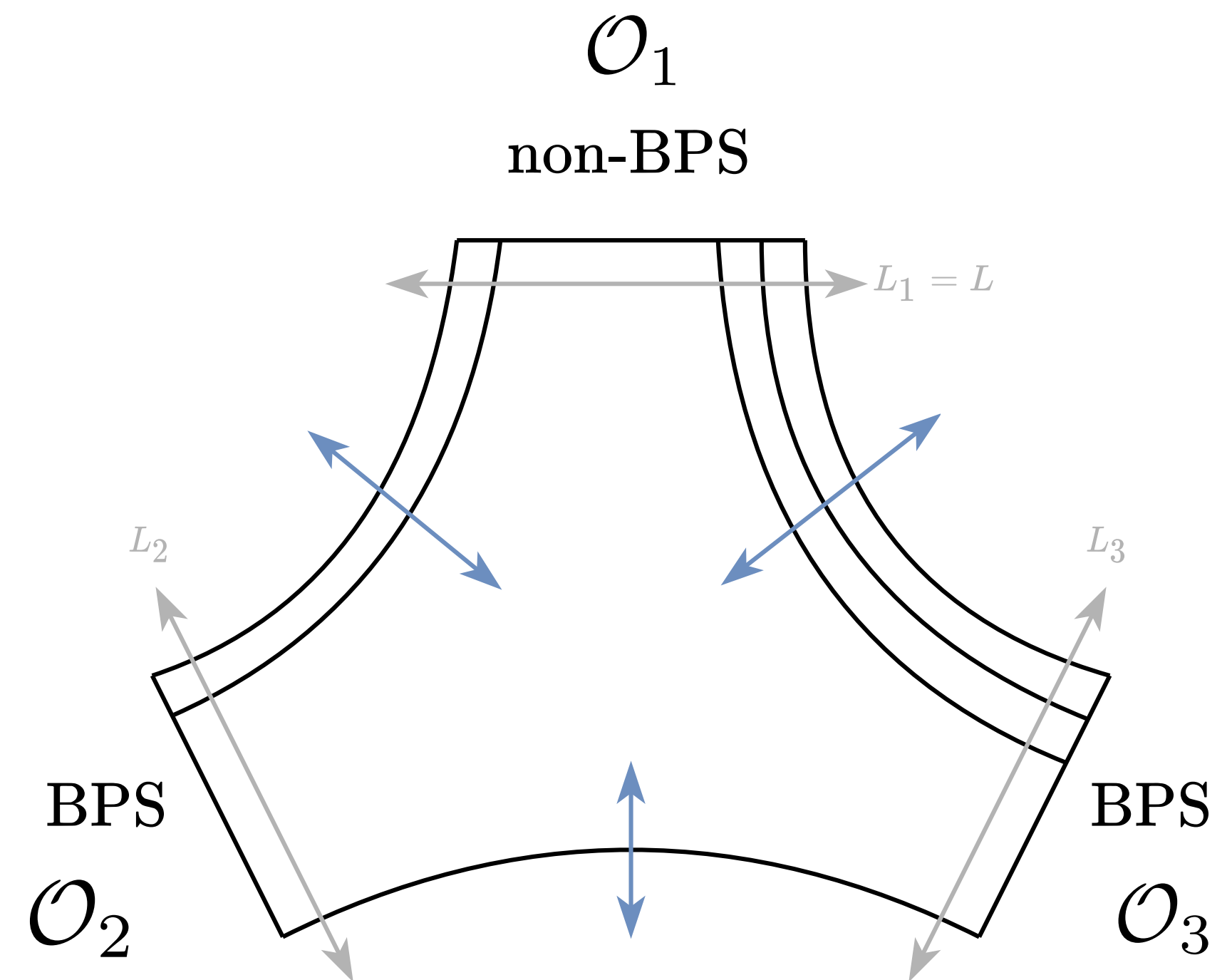
- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables
 - ↳ Evidence that this is promising has been piling up recently
[Cavaglià, Gromov, Levkovich-Maslyuk][Caetano, Komatsu][Giombi, Komatsu][Giombi, Komatsu, Offertaler]...
 - ↳ Number of technical developments serve as a backbone to these explorations
[Cavaglià, Gromov, Levkovich-Maslyuk]²[Gromov, Levkovich-Maslyuk, Ryan][Gromov, Levkovich-Maslyuk, Ryan, Volin]
- In this talk I will present:
 - ↳ Structure constants formulas at weak coupling (up to N^2LO) based on Q-functions and the “SoV” formalism.
 - ↳ Consider only **rank-1** NBPS - BPS - BPS correlators

Context

- The Quantum Spectral Curve resolves this issue in the case of 2-pt functions.
 - ↳ Elegant and efficient framework. [Gromov, Kazakov, Leurent, Volin]
 - ↳ Makes manifest hidden simplicity of the intricate low energy expansion.
- Only natural to look for a Q-function (“SoV”) approach to other observables
 - ↳ Evidence that this is promising has been piling up recently
[Cavaglià, Gromov, Levkovich-Maslyuk][Caetano, Komatsu][Giombi, Komatsu][Giombi, Komatsu, Offertaler]...
 - ↳ Number of technical developments serve as a backbone to these explorations
[Cavaglià, Gromov, Levkovich-Maslyuk]²[Gromov, Levkovich-Maslyuk, Ryan][Gromov, Levkovich-Maslyuk, Ryan, Volin]
- In this talk I will present:
 - ↳ Structure constants formulas at weak coupling (up to N^2LO) based on Q-functions and the “SoV” formalism.
 - ↳ Consider only **rank-1** NBPS - BPS - BPS correlators
 - ↳ Match known CFT data which include finite size effects both in the **adjacent** and **bottom** channels.

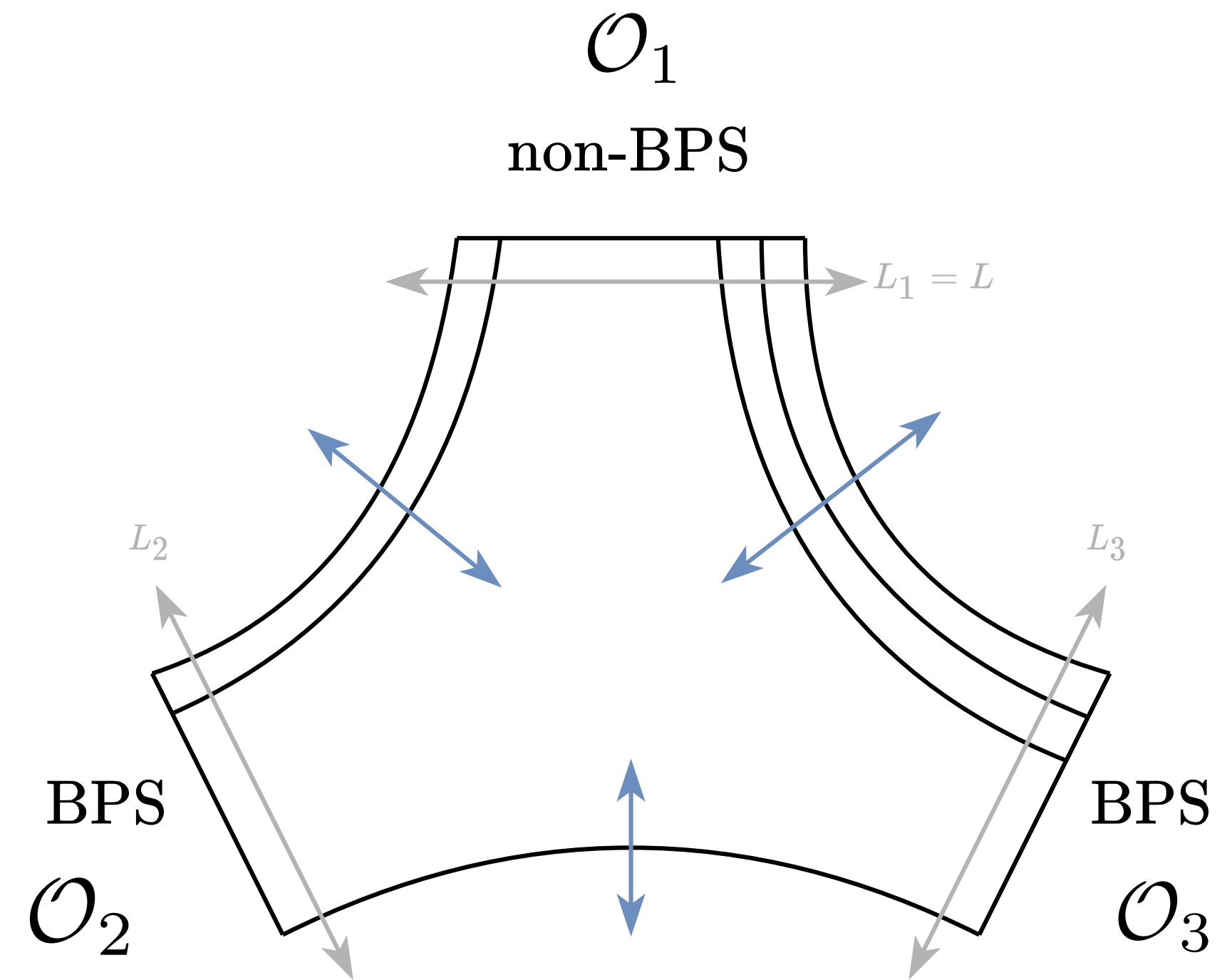
Setup

Setup



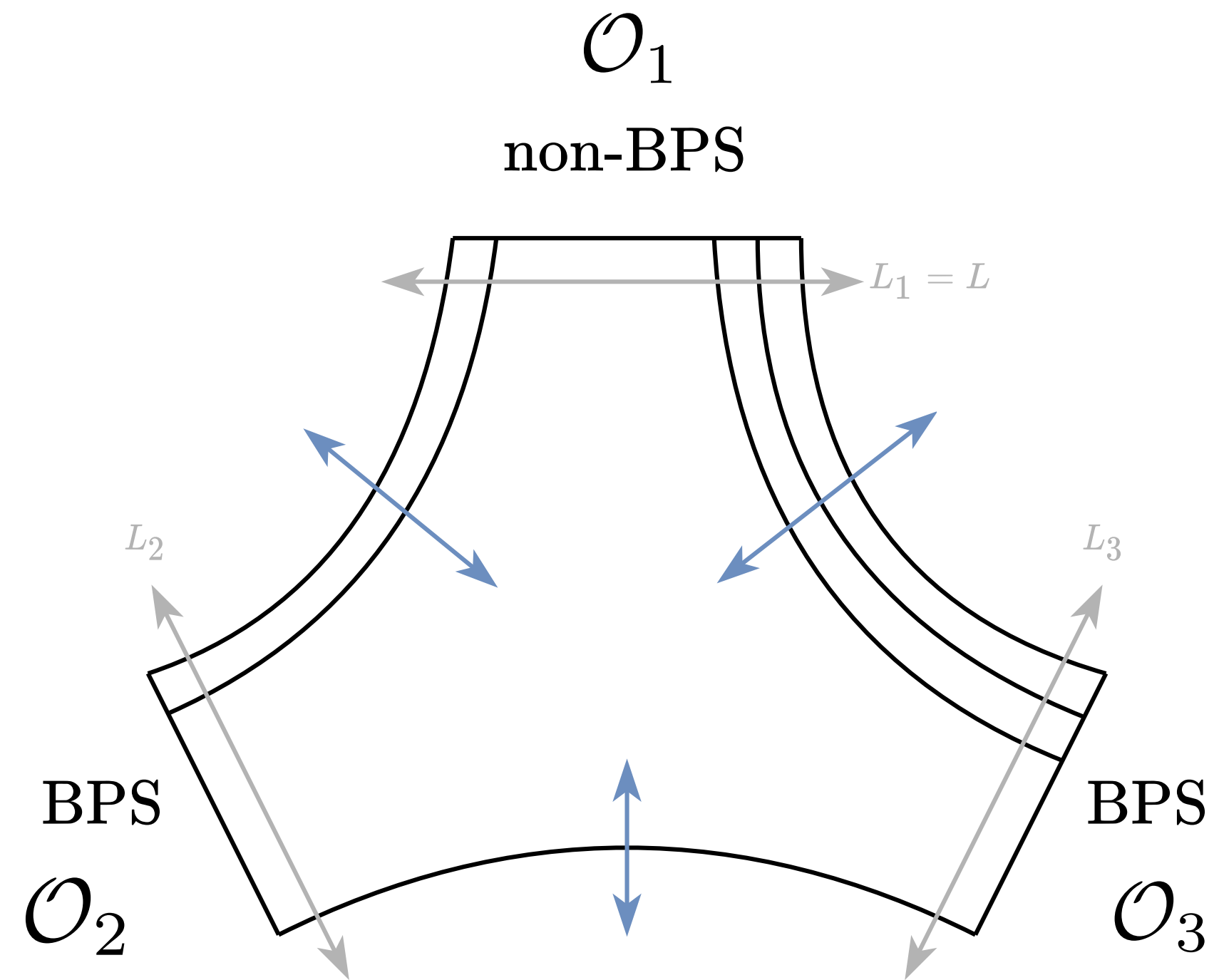
Setup

\mathcal{O}_1 :



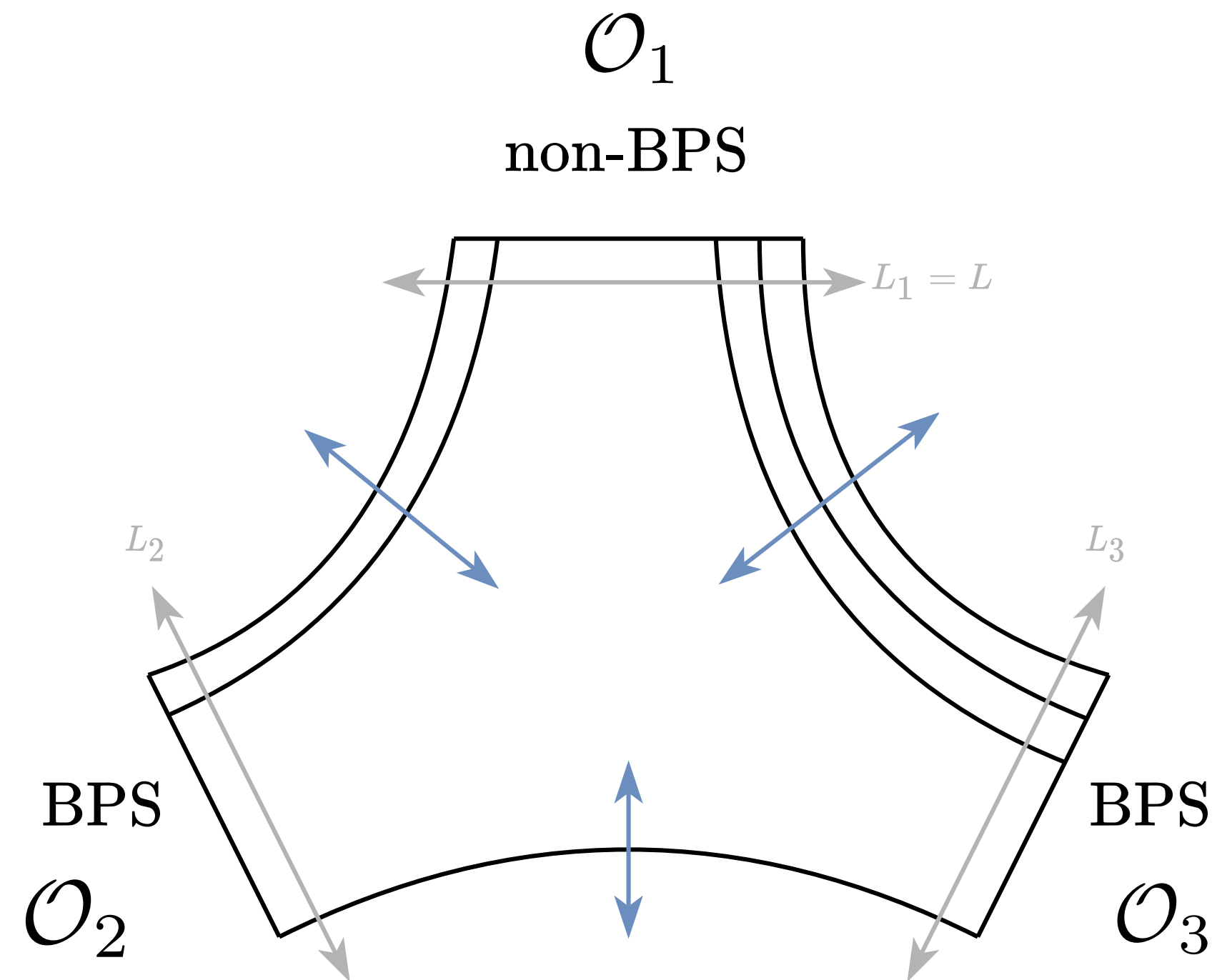
Setup

\mathcal{O}_1 : $SL(2)$: $\text{Tr}(D_+^J Z^L)$ + permutations



Setup

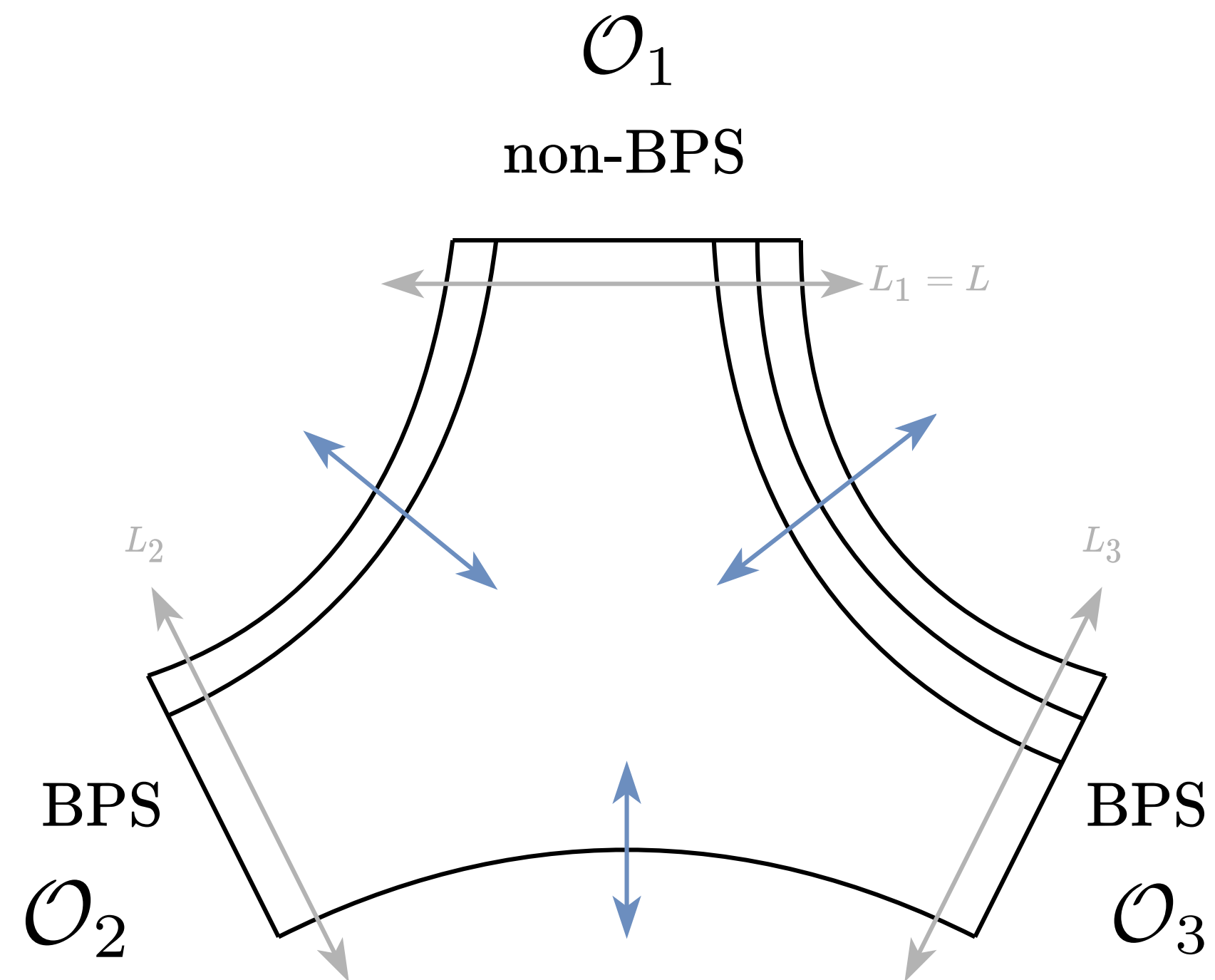
$$\mathcal{O}_1 : \begin{cases} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{cases}$$



Setup

$$\mathcal{O}_1 : \begin{cases} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{cases}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

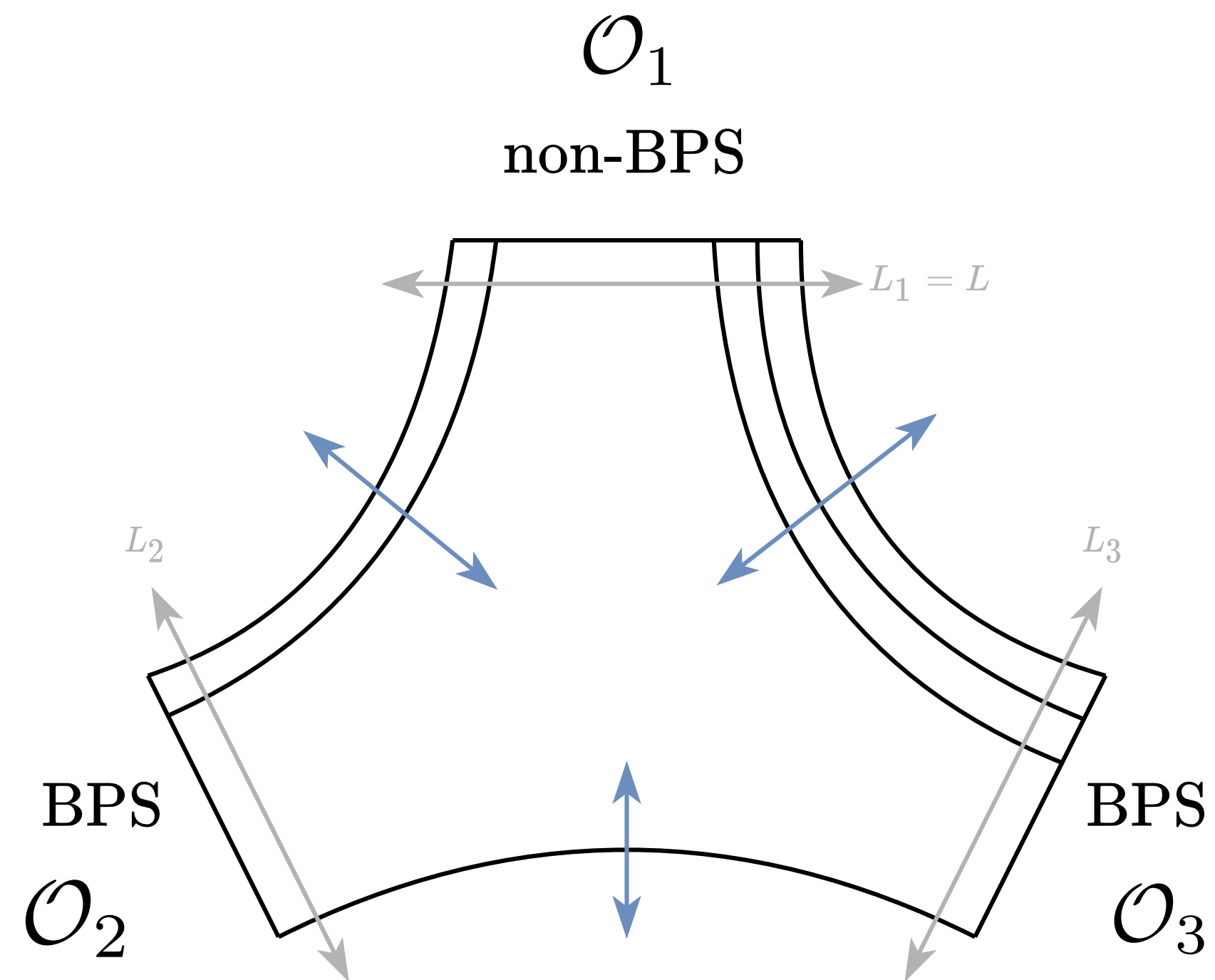


Setup

$$\mathcal{O}_1 : \begin{array}{l} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{array}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

$$\mathcal{O}_3 : \text{Tr}(y_3 \cdot \phi)^{L_3}$$

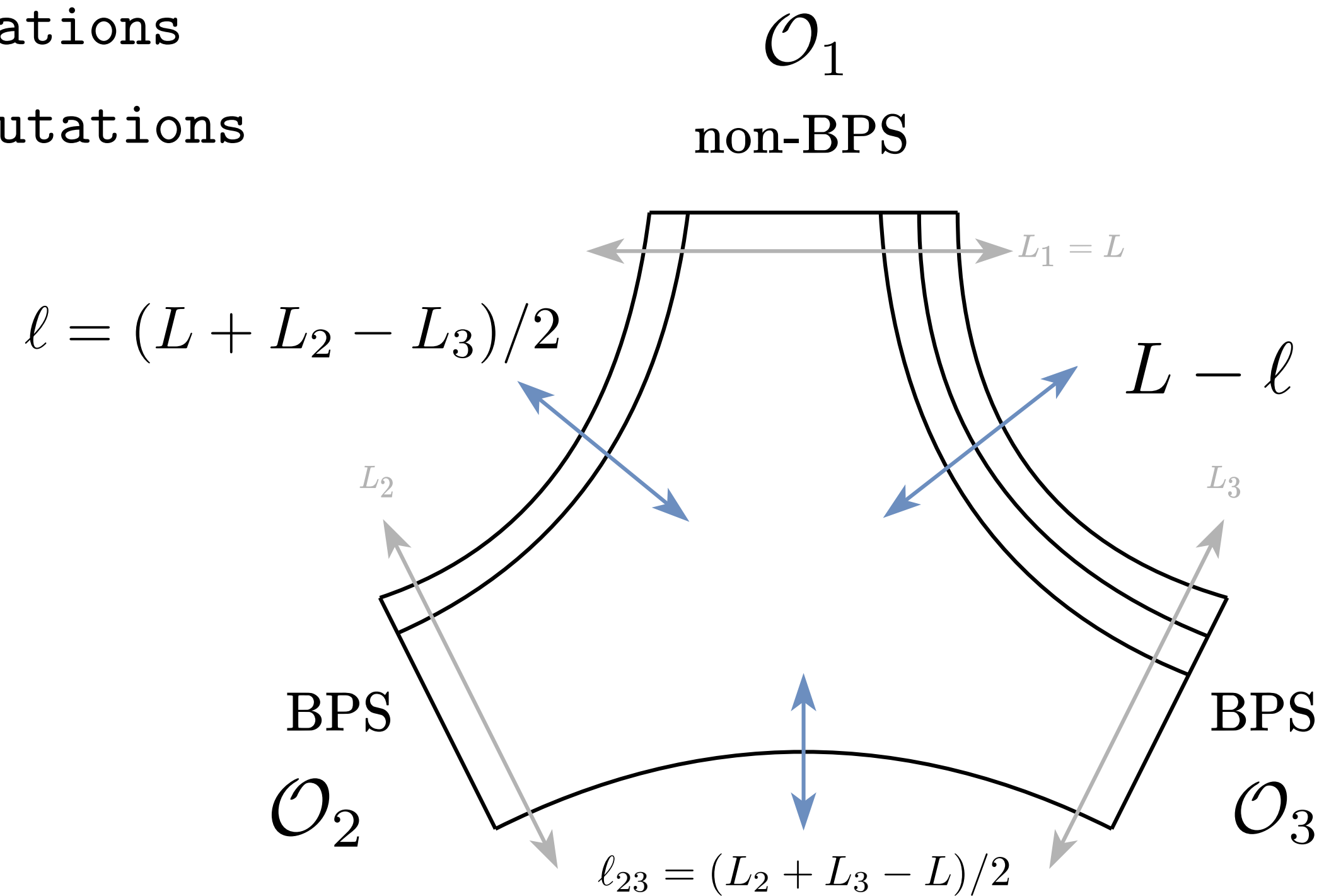


Setup

$$\mathcal{O}_1 : \begin{array}{l} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{array}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

$$\mathcal{O}_3 : \text{Tr}(y_3 \cdot \phi)^{L_3}$$



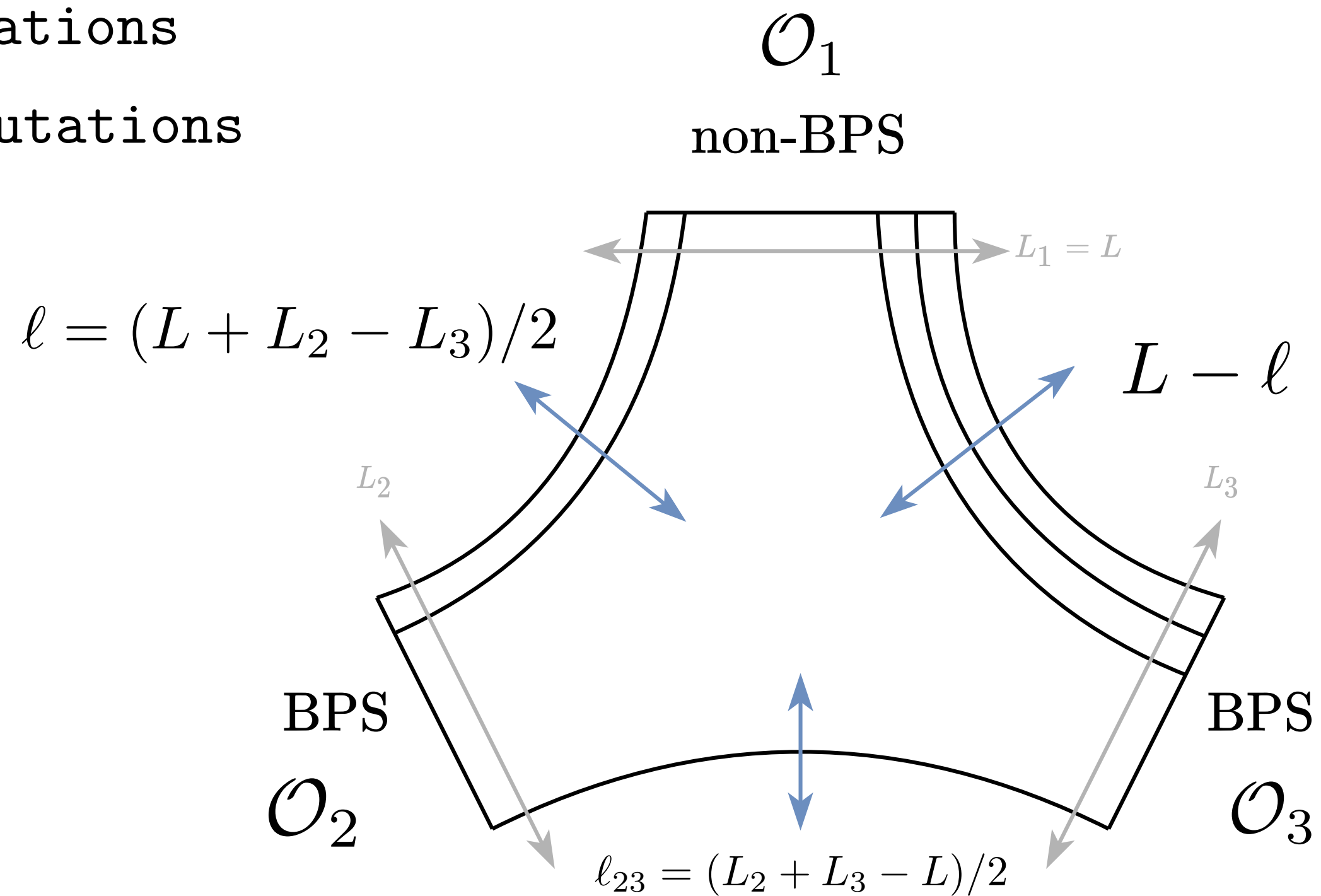
Setup

$$\mathcal{O}_1 : \begin{cases} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{cases}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

$$\mathcal{O}_3 : \text{Tr}(y_3 \cdot \phi)^{L_3}$$

- Step 0: Leading Order (LO)



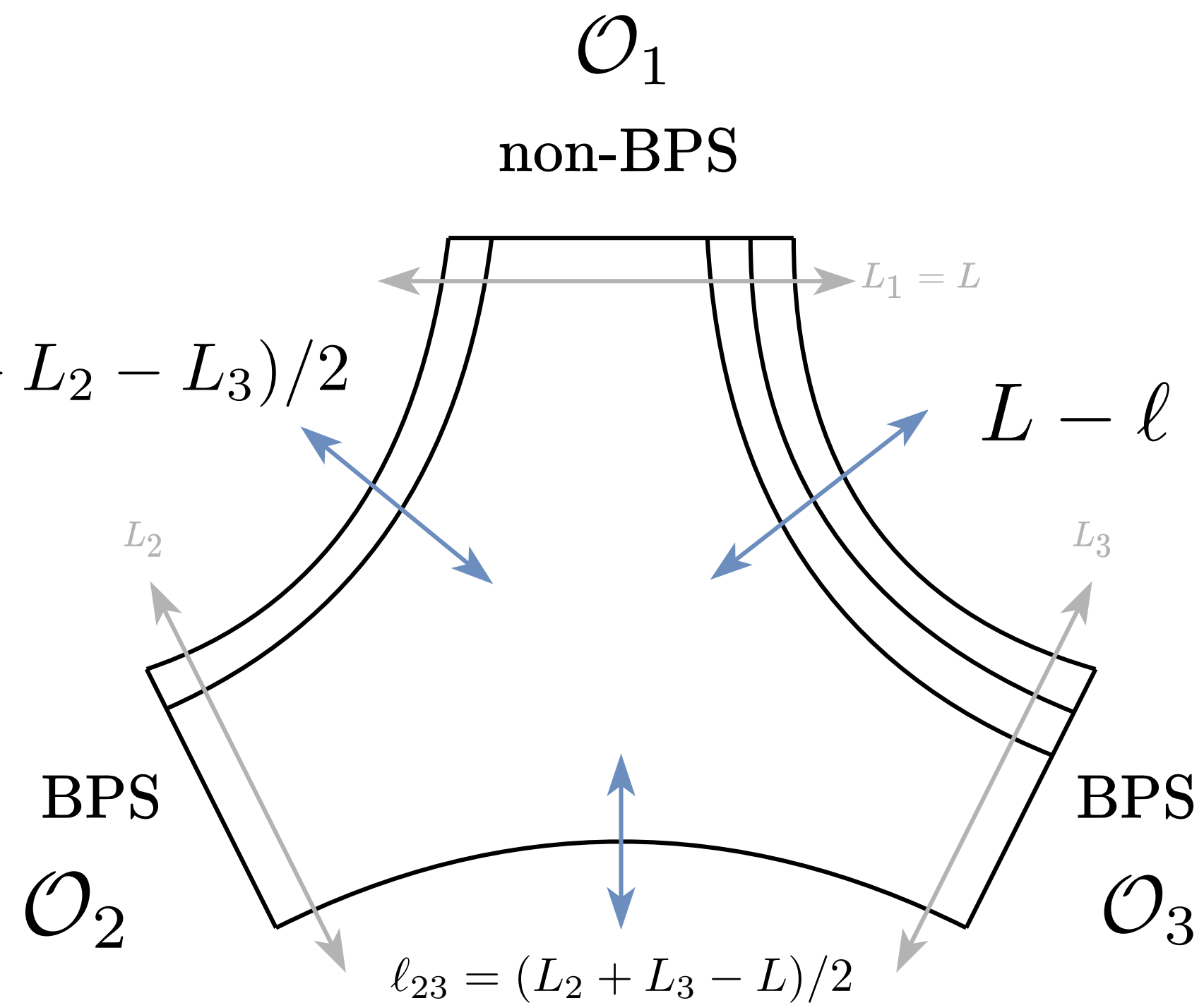
Setup

$$\mathcal{O}_1 : \begin{array}{l} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{array}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

$$\mathcal{O}_3 : \text{Tr}(y_3 \cdot \phi)^{L_3}$$

$$\ell = (L + L_2 - L_3)/2$$



- Step 0: Leading Order (LO)

↳ NBPS wave-function diagonalize 1-loop H

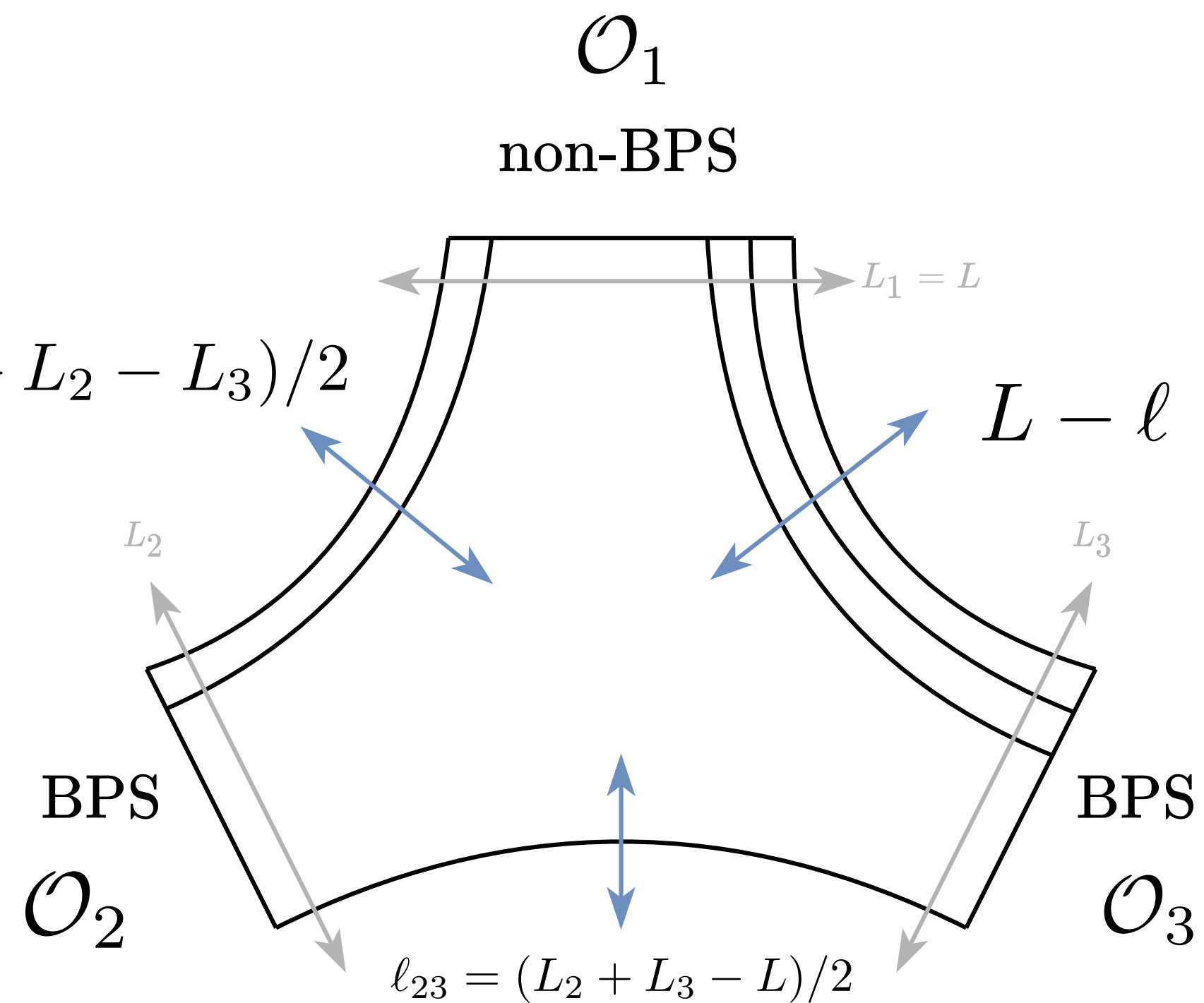
Setup

$$\mathcal{O}_1 : \begin{array}{l} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{array}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

$$\mathcal{O}_3 : \text{Tr}(y_3 \cdot \phi)^{L_3}$$

$$\ell = (L + L_2 - L_3)/2$$



- Step 0: Leading Order (LO)

- ↳ NBPS wave-function diagonalize 1-loop H
- ↳ Structure constant by Wick contractions

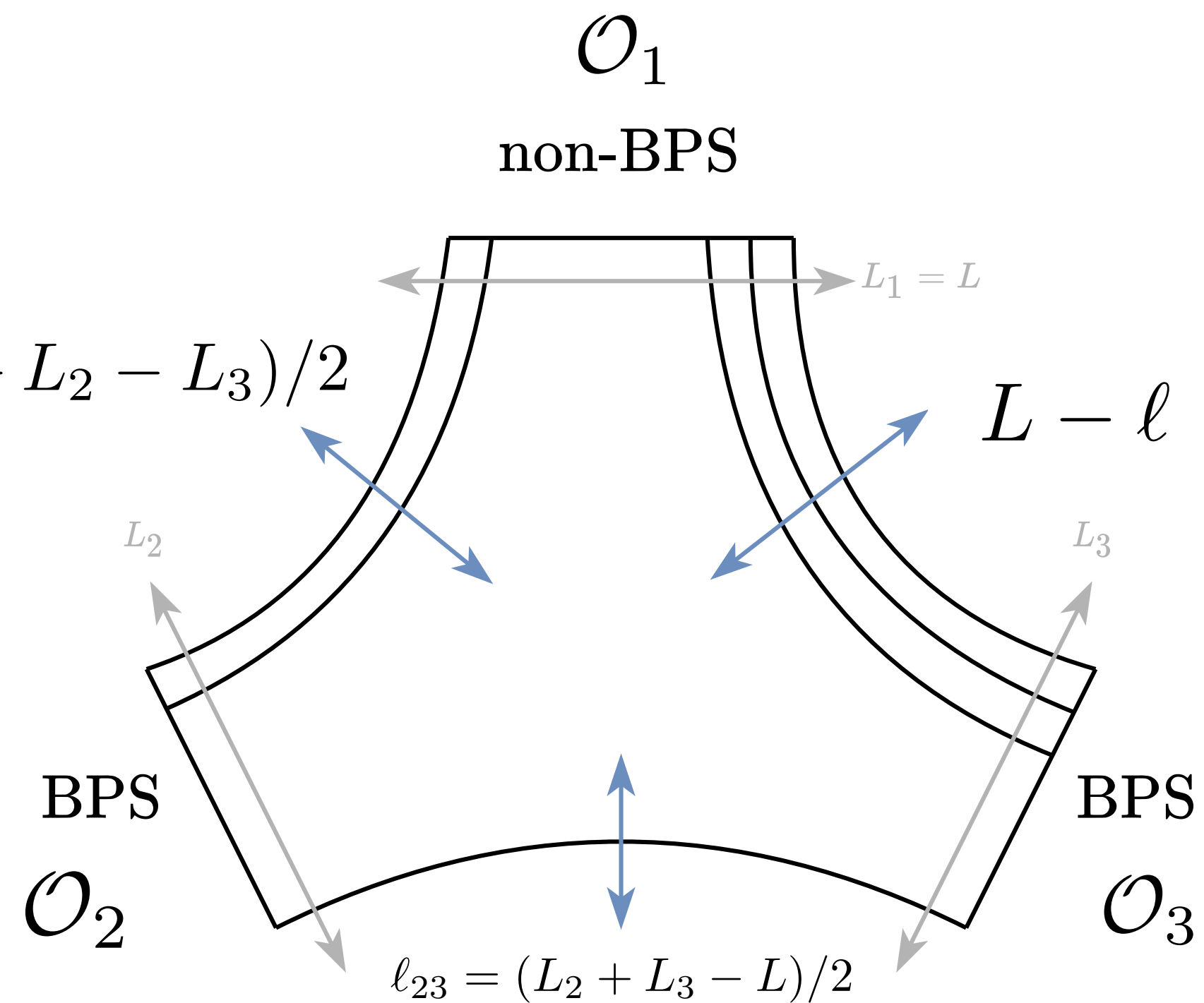
Setup

$$\mathcal{O}_1 : \begin{array}{l} SL(2) : \text{Tr}(D_+^J Z^L) + \text{permutations} \\ SU(2) : \text{Tr}(X^J Z^{L-J}) + \text{permutations} \end{array}$$

$$\mathcal{O}_2 : \text{Tr}(y_2 \cdot \phi)^{L_2}$$

$$\mathcal{O}_3 : \text{Tr}(y_3 \cdot \phi)^{L_3}$$

$$\ell = (L + L_2 - L_3)/2$$



- Step 0: Leading Order (LO)

- ↳ NBPS wave-function diagonalize 1-loop H

- ↳ Structure constant by Wick contractions

- ↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_{\ell^2}}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A{}^2}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A{}^2}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}^2$$

→ off-shell overlap in spin-chain of length ℓ
→ on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

$$Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$$

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ
 on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

$$Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$$

$$\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

↘ off-shell overlap in spin-chain of length ℓ
→ on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

$$Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$$

$$\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ
 on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

$$Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$$

$$\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

SL(2):

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

→ off-shell overlap in spin-chain of length ℓ
→ on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

$$Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$$

$$\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

SL(2):

$$\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$$

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | |
|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ |
|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

→ off-shell overlap in spin-chain of length ℓ
→ on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | |
|--|---|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_{\Gamma} d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1}$ |
|--|---|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | |
|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ |
|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | |
|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_{\Gamma} d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <div style="display: flex; justify-content: space-between;"> <div style="width: 45%;"> <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ </div> <div style="width: 45%;"> <p>SU(2):</p> </div> </div> $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ |
|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | | |
|--|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ | <p>SU(2):</p> $\mu_1(u) = \frac{\sinh(2\pi u)}{(u^2 + 1/4)^2}$ |
|--|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | | |
|--|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ | <p>SU(2):</p> $\mu_1(u) = \frac{\sinh(2\pi u)}{(u^2 + 1/4)^2}$ $\mu_2(u, v) = \frac{\sinh(2\pi(u - v))(u - v)}{2(u^2 + 1/4)(v^2 + 1/4)}$ |
|--|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | | |
|--|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ | <p>SU(2):</p> $\mu_1(u) = \frac{\sinh(2\pi u)}{(u^2 + 1/4)^2}$ $\mu_2(u, v) = \frac{\sinh(2\pi(u - v))(u - v)}{2(u^2 + 1/4)(v^2 + 1/4)}$ $\Gamma = S^{1^{\ell-1}}$ |
|--|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | | |
|--|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ | <p>SU(2):</p> $\mu_1(u) = \frac{\sinh(2\pi u)}{(u^2 + 1/4)^2}$ $\mu_2(u, v) = \frac{\sinh(2\pi(u - v))(u - v)}{2(u^2 + 1/4)(v^2 + 1/4)}$ $\Gamma = S^{1^{\ell-1}} \quad \mathcal{N} = \binom{\ell}{J_A + J_B}$ |
|--|--|--|

Setup

- Step 0: Leading Order (LO)

↳ Combinatorics of contractions: rational spin-chain scalar products [Escobedo, Gromov, Sever, Vieira]

$$(C^{\bullet\circ\circ})^2 = \frac{J!}{2J!} \frac{\overbrace{\langle \Omega | u_1, \dots, u_J \rangle_\ell}^A{}^2}{\underbrace{\langle u_1, \dots, u_J | u_1, \dots, u_J \rangle_L}_B}$$

off-shell overlap in spin-chain of length ℓ

on-shell norm in spin-chain of length L

$$\mathcal{A} = \langle Q_J, 1 \rangle_\ell / \langle 1, 1 \rangle_\ell$$

$$\mathcal{B} = \langle Q_J, Q_J \rangle_L / \langle 1, 1 \rangle_L$$

- Rational scalar products admit known integral representations from SoV! [Derkachov, Korchemsky, Manashov][Kazama, Komatsu, Nishimura]

| | | |
|--|--|--|
| $Q_A(u) = \prod_{u_i \in A} \frac{u - u_i}{\sqrt{u_i^2 + 1/4}}$ $\langle Q_A, Q_B \rangle_\ell = \mathcal{N} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$ $d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$ | <p>SL(2):</p> $\mu_1(u) = \frac{\pi}{2 \cosh(\pi u)^2}$ $\mu_2(u, v) = \frac{\pi(u - v) \sinh(\pi(u - v))}{\cosh(\pi u) \cosh(\pi v)}$ $\Gamma = \mathbb{R}^{\ell-1} \quad \mathcal{N} = \binom{J_A + J_B + \ell - 1}{\ell - 1}$ | <p>SU(2):</p> $\mu_1(u) = \frac{\sinh(2\pi u)}{(u^2 + 1/4)^2}$ $\mu_2(u, v) = \frac{\sinh(2\pi(u - v))(u - v)}{2(u^2 + 1/4)(v^2 + 1/4)}$ $\Gamma = S^{1^{\ell-1}} \quad \mathcal{N} = \binom{\ell}{J_A + J_B}$ |
|--|--|--|

Plan

Plan

- Revisit LO results from perspective of functional SoV

Plan

- Revisit LO results from perspective of functional SoV
- Generalize and lift $SL(2)$ results to the quantum level (NLO)

Plan

- Revisit LO results from perspective of functional SoV
- Generalize and lift $SL(2)$ results to the quantum level (NLO)
- Alternative derivation up to N^2LO for $SU(2)$, including adjacent finite volume effects.

Plan

- Revisit LO results from perspective of functional SoV
- Generalize and lift $SL(2)$ results to the quantum level (NLO)
- Alternative derivation up to N^2LO for $SU(2)$, including adjacent finite volume effects.
- Conclusion: features, issues, puzzles and prospects

Plan

- Revisit LO results from perspective of functional SoV
- Generalize and lift $SL(2)$ results to the quantum level (NLO)
- Alternative derivation up to N^2LO for $SU(2)$, including adjacent finite volume effects.
- Conclusion: features, issues, puzzles and prospects
- Extra: some $SL(2)$ N^2LO results, matching of bottom finite volume effects.

Plan

- Revisit LO results from perspective of functional SoV
- Generalize and lift $SL(2)$ results to the quantum level (NLO)
- Alternative derivation up to N^2LO for $SU(2)$, including adjacent finite volume effects.
- Conclusion: features, issues, puzzles and prospects
- Extra: some $SL(2)$ N^2LO results, matching of bottom finite volume effects.

Questions about context, setup and plan?

Functional SoV

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2 SL(2)$. Want to determine (μ, Γ) .

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2 SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.
 - ↳ Consider Baxter operator

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator

$$\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$$

$$\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator

$$\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$$

$$\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator

$$\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$$

$$\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A))$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left((u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial} \right) Q_B$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left(\overset{u \rightarrow u - i}{\cancel{(u + i/2)^2}} e^{i\partial} + (u - i/2)^2 e^{-i\partial} \right) Q_B$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left(\begin{array}{c} u \rightarrow u - i \\ \nearrow \\ (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial} \\ \searrow \\ u \rightarrow u + i \end{array} \right) Q_B$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left((u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial} \right) Q_B$$

$u \rightarrow u - i$ (pointing to $(u + i/2)^2$)
 $u \rightarrow u + i$ (pointing to $(u - i/2)^2$)

$$= Q_B (\mu(u - i)(u - i/2)^2 e^{-i\partial} + \mu(u + i)(u + i/2)^2 e^{+i\partial}) Q_A$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$

- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left(\overset{u \rightarrow u - i}{\cancel{(u + i/2)^2 e^{i\partial}}} + \overset{u \rightarrow u + i}{\cancel{(u - i/2)^2 e^{-i\partial}}} \right) Q_B$$

μ is i-periodic

$$= Q_B (\mu(u - i)(u - i/2)^2 e^{-i\partial} + \mu(u + i)(u + i/2)^2 e^{+i\partial}) Q_A$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$
- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left((u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial} \right) Q_B$$

$u \rightarrow u - i$ (arrow pointing to $(u + i/2)^2 e^{i\partial}$)
 $u \rightarrow u + i$ (arrow pointing to $(u - i/2)^2 e^{-i\partial}$)

μ is i -periodic
No poles when deforming contour

$$= Q_B (\mu(u - i)(u - i/2)^2 e^{-i\partial} + \mu(u + i)(u + i/2)^2 e^{+i\partial}) Q_A$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$

- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left(\begin{array}{c} u \rightarrow u - i \\ \nearrow \\ (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial} \\ \searrow \\ u \rightarrow u + i \end{array} \right) Q_B$$

μ is i -periodic
No poles when deforming contour $\rightarrow \mathbb{B}$ is self-adjoint

$$= Q_B (\mu(u - i)(u - i/2)^2 e^{-i\partial} + \mu(u + i)(u + i/2)^2 e^{+i\partial}) Q_A$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$

- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left((u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial} \right) Q_B$$

$u \rightarrow u - i$ (arrow pointing to $(u + i/2)^2 e^{i\partial}$)
 $u \rightarrow u + i$ (arrow pointing to $(u - i/2)^2 e^{-i\partial}$)

μ is i -periodic
 No poles when deforming contour $\rightarrow \mathbb{B}$ is self-adjoint

$$\mu(u) = \frac{\pi}{2 \cosh(\pi u)^2} \quad \Gamma = \mathbb{R}$$

$$= Q_B (\mu(u - i)(u - i/2)^2 e^{-i\partial} + \mu(u + i)(u + i/2)^2 e^{+i\partial}) Q_A$$

Functional SoV

[Cavaglià, Gromov, Levkovich-Maslyuk]

- Consider simplest case of $L = 2$ $SL(2)$. Want to determine (μ, Γ) . $\int_{\Gamma} du \mu Q_A Q_B$

- Look for hermitian scalar product: distinct Bethe states/Baxter functions must be orthogonal.

↳ Consider Baxter operator $\mathbb{B} = (u + i/2)^2 e^{i\partial} + (u - i/2)^2 e^{-i\partial}$ $\mathbb{B}Q_A = \underbrace{(2u^2 + c_A)}_{\tau_A} Q_A$

- Suppose (μ, Γ) make \mathbb{B} self-adjoint.

$$0 = \int_{\Gamma} du \mu (Q_A(\mathbb{B}Q_B) - Q_B(\mathbb{B}Q_A)) = (c_A - c_B) \langle Q_A, Q_B \rangle$$

$$\mu(u) Q_A \left(\overset{u \rightarrow u - i}{\cancel{(u + i/2)^2 e^{i\partial}}} + (u - i/2)^2 e^{-i\partial} \right) Q_B$$

$$\mu(u) Q_A \left((u + i/2)^2 e^{i\partial} + \underset{u \rightarrow u + i}{\cancel{(u - i/2)^2 e^{-i\partial}}} \right) Q_B$$

μ is i -periodic
No poles when deforming contour $\rightarrow \mathbb{B}$ is self-adjoint

$$\mu(u) = \frac{\pi}{2 \cosh(\pi u)^2} \quad \Gamma = \mathbb{R}$$

$$= Q_B (\mu(u - i)(u - i/2)^2 e^{-i\partial} + \mu(u + i)(u + i/2)^2 e^{+i\partial}) Q_A$$

Poles of μ cancel with zeroes of \mathbb{B}

Quantum Lift

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$
- Inconvenient: Baxter is state dependent
 - ↳ Define “dressed Q-function”

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$

$$H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$

$$H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial}$$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$

$$H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial} \quad \tilde{\mathbb{B}} Q_A = \tau_A Q_A$$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$

$$H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial} \quad \tilde{\mathbb{B}} Q_A = \tau_A Q_A$$

Repeat strategy to fix μ

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$

$$H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial} \quad \tilde{\mathbb{B}} Q_A = \tau_A Q_A$$

Repeat strategy to fix μ

All poles cancel provided

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$ $H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial} \quad \tilde{\mathbb{B}} Q_A = \tau_A Q_A$$

Repeat strategy to fix μ

All poles cancel provided $\mu(u) = \frac{\pi}{2 \cosh(\pi u)^2} (1 + g^2 \pi (3 \tanh(\pi u)^2 - 1))$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$ $H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial} \quad \tilde{\mathbb{B}} Q_A = \tau_A Q_A$$

Repeat strategy to fix μ

All poles cancel provided

$$\mu(u) = \frac{\pi}{2 \cosh(\pi u)^2} (1 + g^2 \pi (3 \tanh(\pi u)^2 - 1)) = \oint \frac{dv}{2\pi i} \frac{\pi}{2 \cosh^2(\pi(u-v))} \frac{1}{x(v)}$$

Quantum Lift

- $\mathbb{B} = (x^+)^2 \left(1 - \frac{g^2 Q_1^+}{x^-}\right) e^{i\partial} + (x^-)^2 \left(1 - \frac{g^2 Q_1^+}{x^+}\right) e^{-i\partial}$ $Q_1^+ \equiv \sum_i \frac{1}{u_i^2 + 1/4}$

- Inconvenient: Baxter is state dependent

↳ Define “dressed Q-function”

- $Q(u) = \left(\prod_i \frac{u - u_i}{\sqrt{x_i^+ x_i^-}} \right) e^{g^2 Q_1^+ H_1^+(u)/2}$ $H_1^+ \equiv H_1(-1/2 + iu) + H_1(-1/2 - iu)$

$$\tilde{\mathbb{B}} = (x^+)^2 e^{i\partial} + (x^-)^2 e^{-i\partial} \quad \tilde{\mathbb{B}} Q_A = \tau_A Q_A$$

Repeat strategy to fix μ

“Zhukowsky-fy tree level”

All poles cancel provided

$$\mu(u) = \frac{\pi}{2 \cosh(\pi u)^2} (1 + g^2 \pi (3 \tanh(\pi u)^2 - 1)) = \oint \frac{dv}{2\pi i} \frac{\pi}{2 \cosh^2(\pi(u-v))} \frac{1}{x(v)}$$

Quantum Lift

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle \mathcal{Q}_A \mathcal{Q}_B \rangle_2 \propto \delta_{AB}$

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle \mathcal{Q}_A \mathcal{Q}_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

$$(C^{\bullet\circ\circ})^2 = \frac{\mathbb{J}!^2}{(2\mathbb{J})!} \frac{\langle \mathbf{1}, Q_J \rangle_\ell^2}{\langle Q_J, Q_J \rangle_L}$$

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

$$(C^{\bullet\circ\circ})^2 = \frac{\mathbb{J}!^2}{(2\mathbb{J})!} \frac{\langle \mathbf{1}, Q_J \rangle_\ell^2}{\langle Q_J, Q_J \rangle_L}$$

Lessons:

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

$$(C^{\bullet\circ\circ})^2 = \frac{\mathbb{J}!^2}{(2\mathbb{J})!} \frac{\langle \mathbf{1}, Q_J \rangle_\ell^2}{\langle Q_J, Q_J \rangle_L}$$

Lessons:

Q gets dressed (no longer polynomial)

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

$$(C^{\bullet\circ\circ})^2 = \frac{\mathbb{J}!^2}{(2\mathbb{J})!} \frac{\langle \mathbf{1}, Q_J \rangle_\ell^2}{\langle Q_J, Q_J \rangle_L}$$

Lessons:

Q gets dressed (no longer polynomial)

Measure and norm gets corrected

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

$$(C^{\bullet\circ\circ})^2 = \frac{\mathbb{J}!^2}{(2\mathbb{J})!} \frac{\langle \mathbf{1}, Q_J \rangle_\ell^2}{\langle Q_J, Q_J \rangle_L}$$

Lessons:

Q gets dressed (no longer polynomial)

Measure and norm gets corrected

Otherwise classical structure survives

Quantum Lift

- This provides an orthogonal measure for twist 2. $\langle Q_A Q_B \rangle_2 \propto \delta_{AB}$
- From there it is easy to generalize to arbitrary state in the $SL(2)$ sector:

$$\langle Q_A, Q_B \rangle_\ell = \binom{\mathbb{J}_A + \mathbb{J}_B + \ell - 1}{\ell - 1} \int_\Gamma d\mu \prod_{i=1}^{\ell-1} Q_A(x_i) Q_B(x_i)$$

$$d\mu = \prod_{i=1}^{\ell-1} dx_i \mu_1(x_i) \prod_{j=1}^{\ell-1} \mu_2(x_i, x_j)$$

$$\mu_2(u, v) / \mu_2^{\text{tree}} = (1 + g^2 \pi^2 ((\tanh(\pi u) + \tanh(\pi v))^2 - 4/3))$$

$$\mathbb{J} = J + g^2 Q_1^+$$

$$(C^{\bullet\circ\circ})^2 = \frac{\mathbb{J}!^2}{(2\mathbb{J})!} \frac{\langle \mathbf{1}, Q_J \rangle_\ell^2}{\langle Q_J, Q_J \rangle_L}$$

Lessons:

Q gets dressed (no longer polynomial)

Measure and norm gets corrected

Otherwise classical structure survives

Orthogonality remains key

$SU(2)$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0}$$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\circ\circ})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} \begin{matrix} \xrightarrow{\text{ABA}} \\ + O(g^4) \\ \xrightarrow{\text{ABA}} \end{matrix}$$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} \begin{matrix} \xrightarrow{\text{ABA}} \\ + O(g^4) \\ \xrightarrow{\text{ABA}} \end{matrix}$$

$$\mathcal{M} = \exp \left[g^2 \left(\sum_{i=1}^L (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 \right) - ig^2 Q_1^+ (\partial_{\theta_1} - \partial_{\theta_L}) \right]$$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} \begin{matrix} \xrightarrow{\text{SOV}} \\ + O(g^4) \\ \xrightarrow{\text{SOV}} \end{matrix}$$

$$\mathcal{M} = \exp \left[g^2 \left(\sum_{i=1}^L (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 \right) - ig^2 Q_1^+ (\partial_{\theta_1} - \partial_{\theta_L}) \right]$$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} \begin{array}{c} \xrightarrow{\text{SOV}} \\ + O(g^4) \\ \xrightarrow{\text{SOV}} \end{array}$$

$$\mathcal{M} = \exp \left[g^2 \left(\sum_{i=1}^L (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 - \frac{1}{4} g^2 (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 (\partial_{\theta_{i+1}} - \partial_{\theta_{i+2}})^2 \right) - ig^2 Q_1^+ (\partial_{\theta_1} - \partial_{\theta_L}) + g^4 \delta \mathcal{M}_{\text{NNLO-b}} \right]$$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} \begin{array}{c} \xrightarrow{\text{SOV}} \\ + O(g^6) \\ \xrightarrow{\text{SOV}} \end{array}$$

$$\mathcal{M} = \exp \left[g^2 \left(\sum_{i=1}^L (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 - \frac{1}{4} g^2 (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 (\partial_{\theta_{i+1}} - \partial_{\theta_{i+2}})^2 \right) - ig^2 Q_1^+ (\partial_{\theta_1} - \partial_{\theta_L}) + g^4 \delta \mathcal{M}_{\text{NNLO-b}} \right]$$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} + O(g^6)$$

SOV
SOV

$$\mathcal{M} = \exp \left[g^2 \left(\sum_{i=1}^L (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 - \frac{1}{4} g^2 (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 (\partial_{\theta_{i+1}} - \partial_{\theta_{i+2}})^2 \right) - ig^2 Q_1^+ (\partial_{\theta_1} - \partial_{\theta_L}) + g^4 \delta \mathcal{M}_{\text{NNLO-b}} \right]$$

e.g. $(Q_1^+)^2 \partial_{\theta_{L-1}}^2$

SU(2)

- Let us consider an alternative approach in the SU(2) sector.
- SU(2) S-matrix does not receive quantum corrections up to N²LO.
 - ↳ Quantum effects only affect propagation.
 - ↳ “Simulate” quantum “friction” through background of impurities θ_i
 - ↳ This was the strategy of [Gromov, Vieira] in “Tailoring IV: Θ -morphism” paper.

$$(C^{\bullet\bullet\bullet})^2 = \left| \frac{(\mathcal{M} \circ \mathcal{A}_\theta)^2}{\Lambda_{\mathcal{B}} \mathcal{M} \circ \mathcal{B}_\theta} \right|_{\theta=0} + O(g^6)$$

SOV
SOV

- \mathcal{A}_θ depend only on $\theta_1, \dots, \theta_\ell$.

$$\mathcal{M} = \exp \left[g^2 \left(\sum_{i=1}^L (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 - \frac{1}{4} g^2 (\partial_{\theta_i} - \partial_{\theta_{i+1}})^2 (\partial_{\theta_{i+1}} - \partial_{\theta_{i+2}})^2 \right) - ig^2 Q_1^+ (\partial_{\theta_1} - \partial_{\theta_L}) + g^4 \delta \mathcal{M}_{\text{NNLO-b}} \right]$$

e.g. $(Q_1^+)^2 \partial_{\theta_{L-1}}^2$

$SU(2)$

SU(2)

- In the end the result is

$$C_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

SU(2)

- In the end the result is

$$C_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

SU(2)

- In the end the result is

$$C_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L} \qquad \langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$
$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$
$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)}$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle \mathbf{1}, Q \rangle\rangle_{\ell,L}}{\langle\langle Q, Q \rangle\rangle_{L,L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell,L} \equiv \langle f, g \rangle_{\ell,L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell,L} \qquad \langle Q_1, Q_2 \rangle_{\ell,L} \equiv \binom{\ell}{J_1 + J_2} \oint_\gamma d\mu_{\ell,L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$d\mu_{\ell,L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell,L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell,L}^2(u_i, u_j)$$

$$\mu_{\ell,L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell,L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell,L}(u)$$

$$\mu_{\ell,L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell,L}(u,v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

As in NLO SL(2):

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \int_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathcal{A}_{\ell, L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset \mathcal{A}_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathcal{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathcal{B}(u)}$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathcal{A}_{\ell, L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathcal{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathcal{B}(u)}$$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

Q is dressed by a function of higher integrable charges. $\frac{g^4(Q_1^+)^2}{x^+ x^-} \subset B(u)$

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle \mathbf{1}, Q \rangle\rangle_{\ell,L}}{\langle\langle Q, Q \rangle\rangle_{L,L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell,L} \equiv \langle f, g \rangle_{\ell,L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell,L}$$

$$d\mu_{\ell,L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell,L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell,L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell,L} \equiv \binom{\ell}{J_1 + J_2} \int_{\gamma} d\mu_{\ell,L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$\mu_{\ell,L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell,L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell,L}(u)$$

$$\mu_{\ell,L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell,L}(u,v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

Q is dressed by a function of higher integrable charges. $\frac{g^4(Q_1^+)^2}{x^+ x^-} \subset B(u)$

New relative to SL(2):

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \int_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

Q is dressed by a function of higher integrable charges. $\frac{g^4(Q_1^+)^2}{x^+ x^-} \subset B(u)$

New relative to SL(2):

Measure in \mathcal{A} is NOT the orthogonal measure for on-shell Q .

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\bullet\bullet}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)} \quad \xrightarrow{\quad} \quad \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

Q is dressed by a function of higher integrable charges. $\frac{g^4(Q_1^+)^2}{x^+ x^-} \subset B(u)$

New relative to SL(2):

Measure in \mathcal{A} is NOT the orthogonal measure for on-shell Q .
(It is for \mathcal{B} .)

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \oint_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)} \xrightarrow{\quad} \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

Q is dressed by a function of higher integrable charges. $\frac{g^4(Q_1^+)^2}{x^+ x^-} \subset B(u)$

New relative to SL(2):

Measure in \mathcal{A} is NOT the orthogonal measure for on-shell Q .
(It is for \mathcal{B} .)

At N²LO the dressing of the state depends on ℓ .

SU(2)

- In the end the result is

$$\mathcal{C}_{\bullet\circ\circ}^2 = \Lambda_\ell \frac{J!^2}{(2J)!} \frac{\langle\langle 1, Q \rangle\rangle_{\ell, L}}{\langle\langle Q, Q \rangle\rangle_{L, L}}$$

$$\langle\langle f, g \rangle\rangle_{\ell, L} \equiv \langle f, g \rangle_{\ell, L} / \langle \mathbf{1}, \mathbf{1} \rangle_{\ell, L}$$

$$d\mu_{\ell, L} = \prod_{i=1}^{\ell-1} du_i \mu_{\ell, L}^1(u_i) \prod_{j \neq i}^{\ell-1} \mu_{\ell, L}^2(u_i, u_j)$$

$$\langle Q_1, Q_2 \rangle_{\ell, L} \equiv \binom{\ell}{J_1 + J_2} \int_{\gamma} d\mu_{\ell, L} \prod_{i=1}^{\ell-1} Q_1(u_i) Q_2(u_i).$$

$$\mu_{\ell, L}^1(u) = \frac{\sinh(2\pi u)}{(x_u^+ x_u^-)^2} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u)} \xrightarrow{\quad} \frac{g^4}{x^+(u)^4} \subset A_{\ell, L}(u)$$

$$\mu_{\ell, L}^2(u, v) = \frac{\sinh(2\pi(u-v))(u-v)}{2x_u^+ x_u^- x_v^+ x_v^-} e^{\delta_{\ell \neq L} \mathbf{A}_{\ell, L}(u, v)}$$

$$Q(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} e^{\delta_{\ell \neq L} \mathbf{B}(u)}$$

As in NLO SL(2):

\mathcal{A} is given by an $\ell - 1$ dimensional integral with factorized measure.

\mathcal{B} is given by an $L - 1$ dimensional integral with factorized measure.

Q is dressed by a function of higher integrable charges. $\frac{g^4(Q_1^+)^2}{x^+ x^-} \subset B(u)$

New relative to SL(2):

Measure in \mathcal{A} is NOT the orthogonal measure for on-shell Q .
(It is for \mathcal{B} .)

At N²LO the dressing of the state depends on ℓ .

$$\delta_{\ell, L-1} \frac{g^4(Q_1^+)^2}{x^{+2}} \subset B(u)$$

Finite Size effect in $SU(2)$

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

Finite Size effect in SU(2)

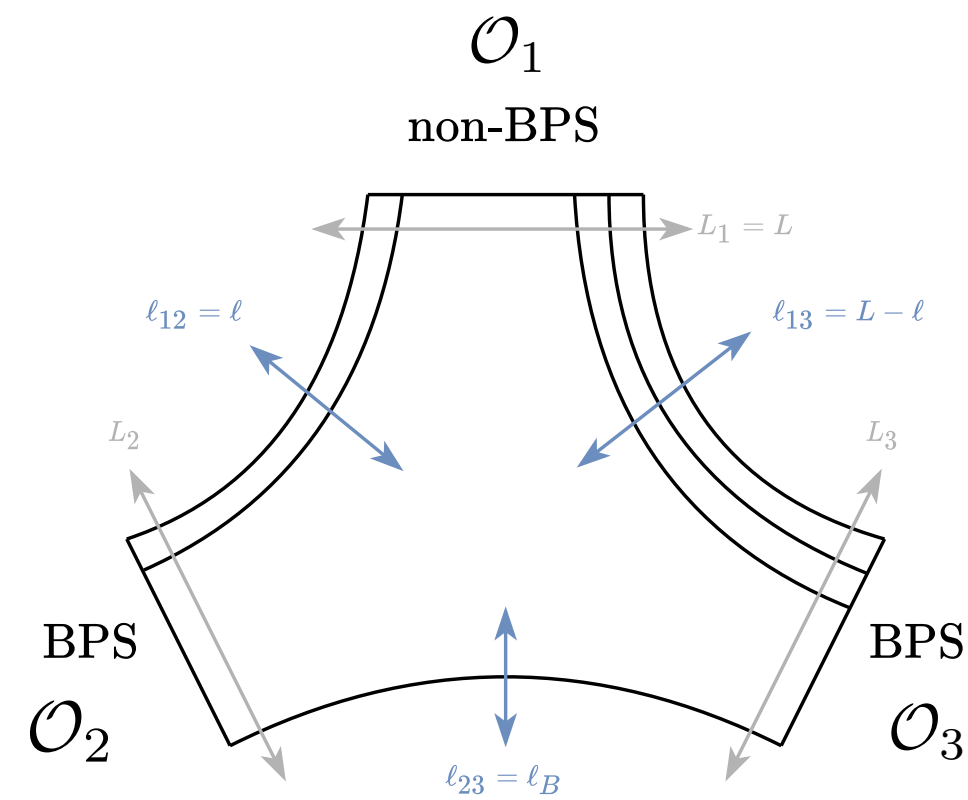
- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.
 - ↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.
 - ↳ C_J vanishes if $\ell < M$ or $L - \ell < M$
- Very easy to understand classically:

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.
 - ↳ C_J vanishes if $\ell < M$ or $L - \ell < M$
- Very easy to understand classically:

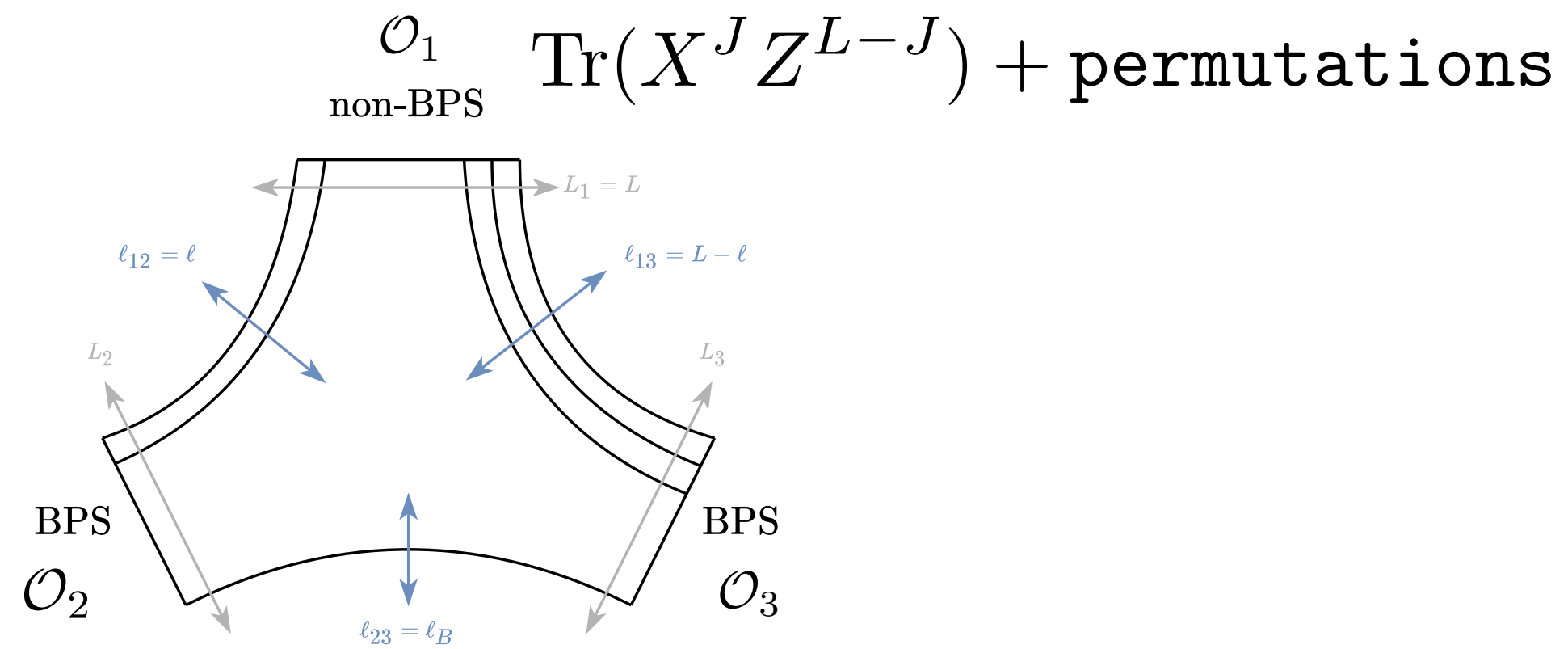


Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:

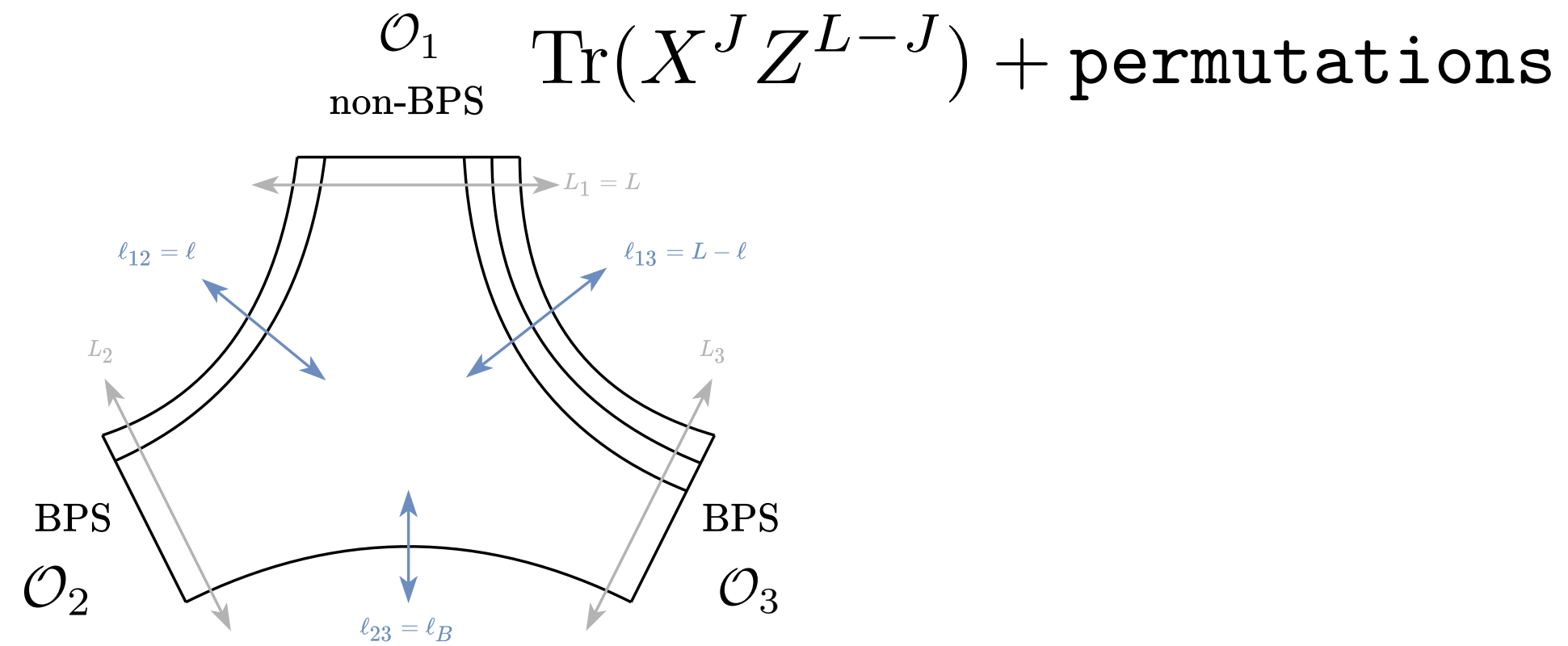


Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:



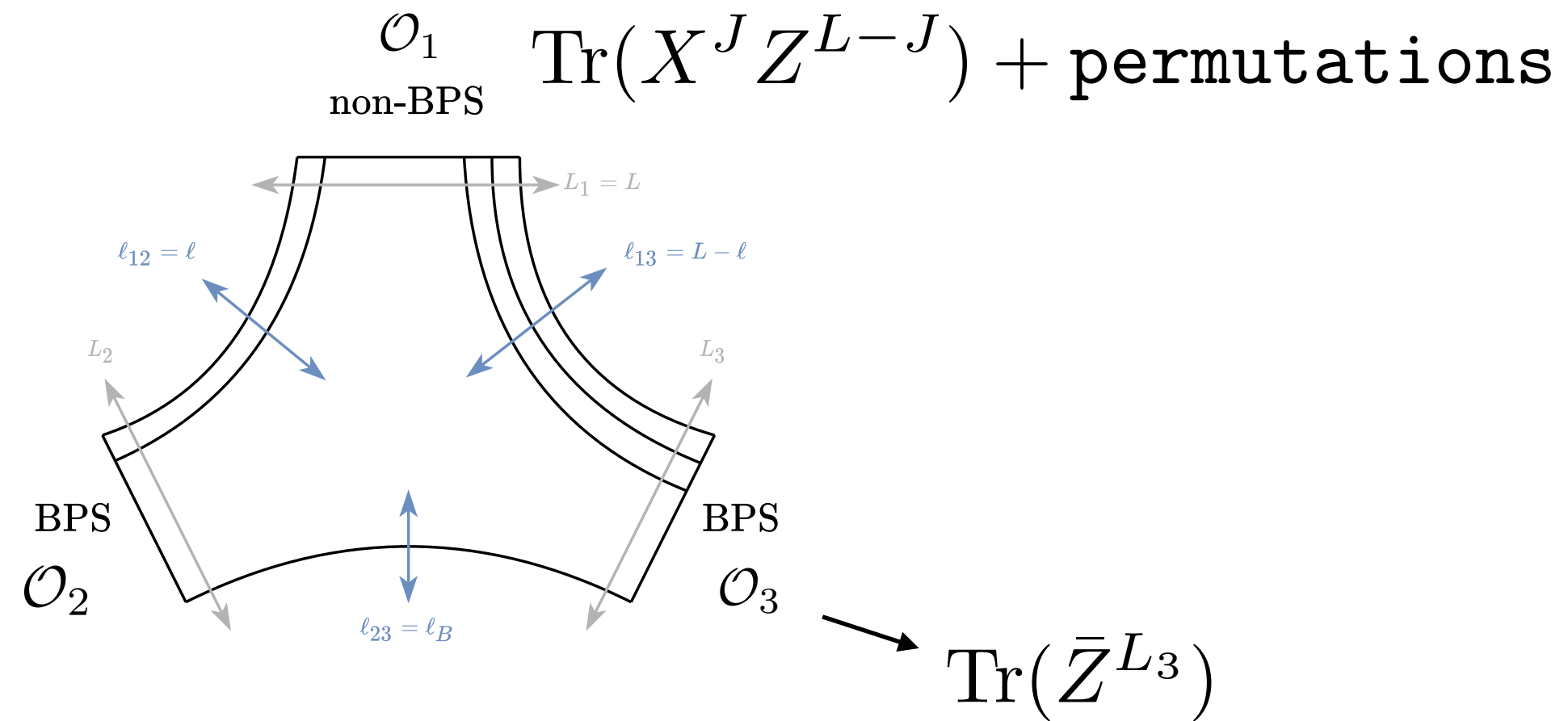
- Choice of super-conformal frame forces all contractions along the same bridge.

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:



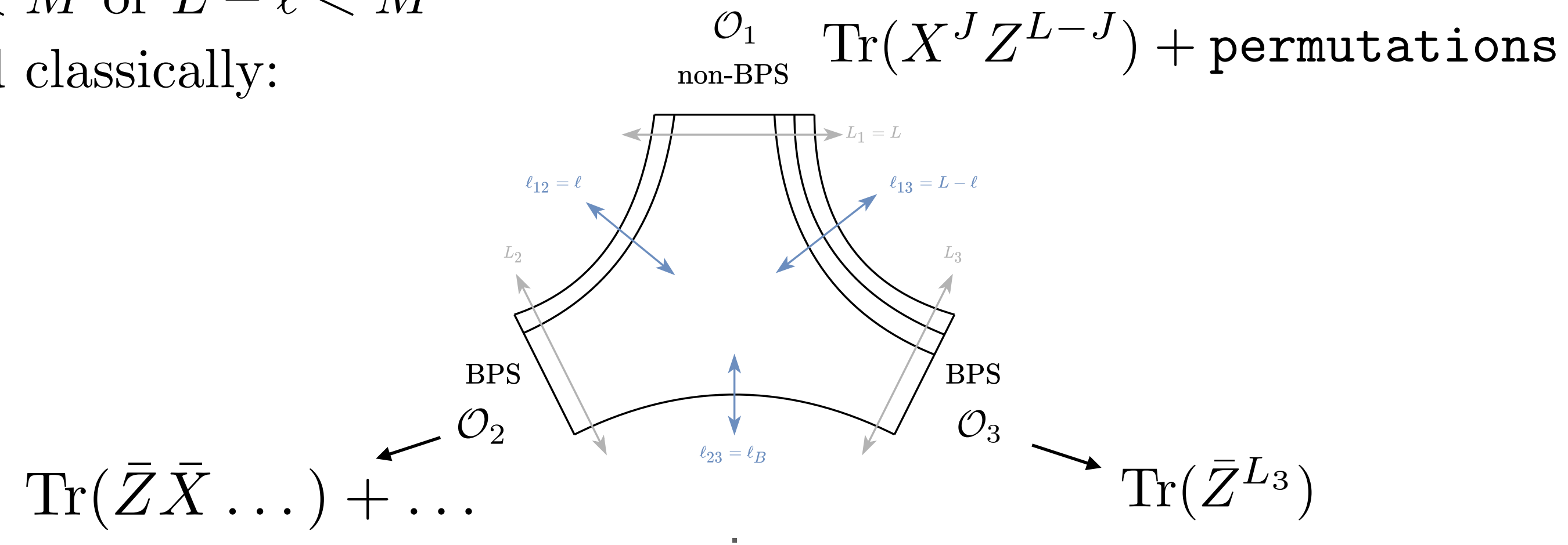
- Choice of super-conformal frame forces all contractions along the same bridge.

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:



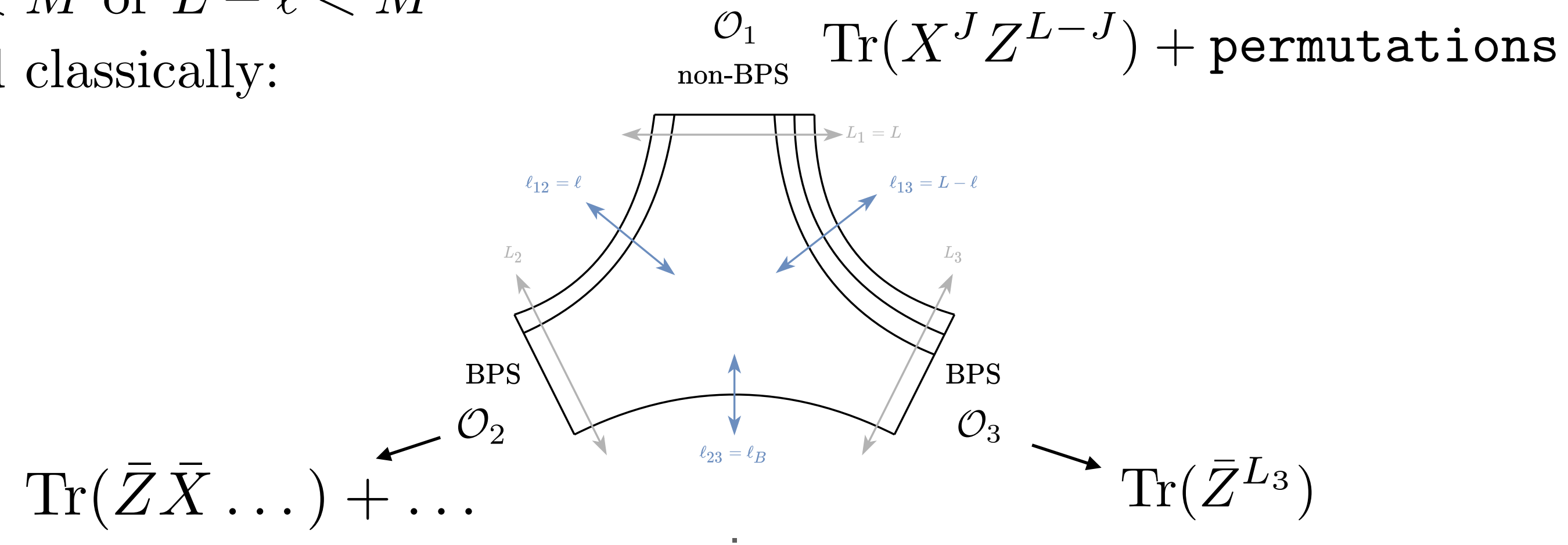
- Choice of super-conformal frame forces all contractions along the same bridge.

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:



- Choice of super-conformal frame forces all contractions along the same bridge.

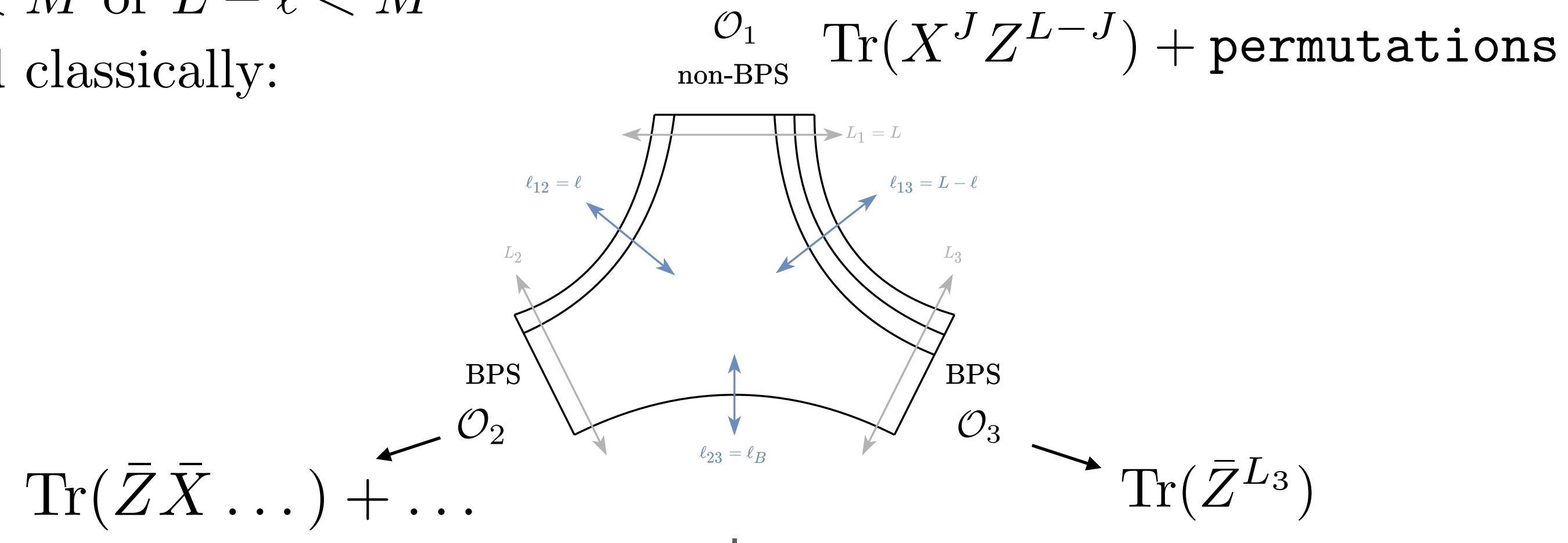
↳ But with ℓ contractions \mathcal{O}_2 can only absorb ℓ of the J excitations.

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:



- Choice of super-conformal frame forces all contractions along the same bridge.

↳ But with ℓ contractions \mathcal{O}_2 can only absorb ℓ of the J excitations.

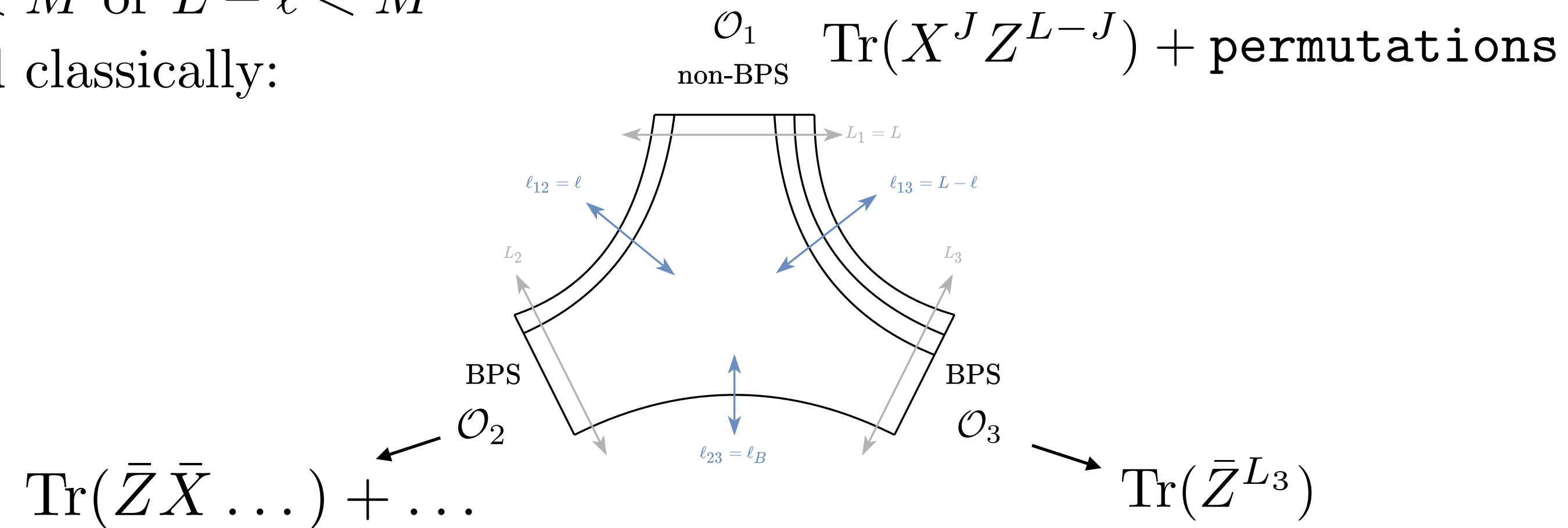
↳ $\ell \leftrightarrow L - \ell$ symmetry by exchanging the role of \mathcal{O}_2 and \mathcal{O}_3 .

Finite Size effect in SU(2)

- Selection rule for 3-pt functions of SU(2) primaries and two $\frac{1}{2}$ BPS operators.

↳ C_J vanishes if $\ell < M$ or $L - \ell < M$

- Very easy to understand classically:



- Choice of super-conformal frame forces all contractions along the same bridge.

↳ But with ℓ contractions \mathcal{O}_2 can only absorb ℓ of the J excitations.

↳ $\ell \leftrightarrow L - \ell$ symmetry by exchanging the role of \mathcal{O}_2 and \mathcal{O}_3 .

- At the quantum level, one can argue group theoretically.

Finite Size effect in $SU(2)$

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

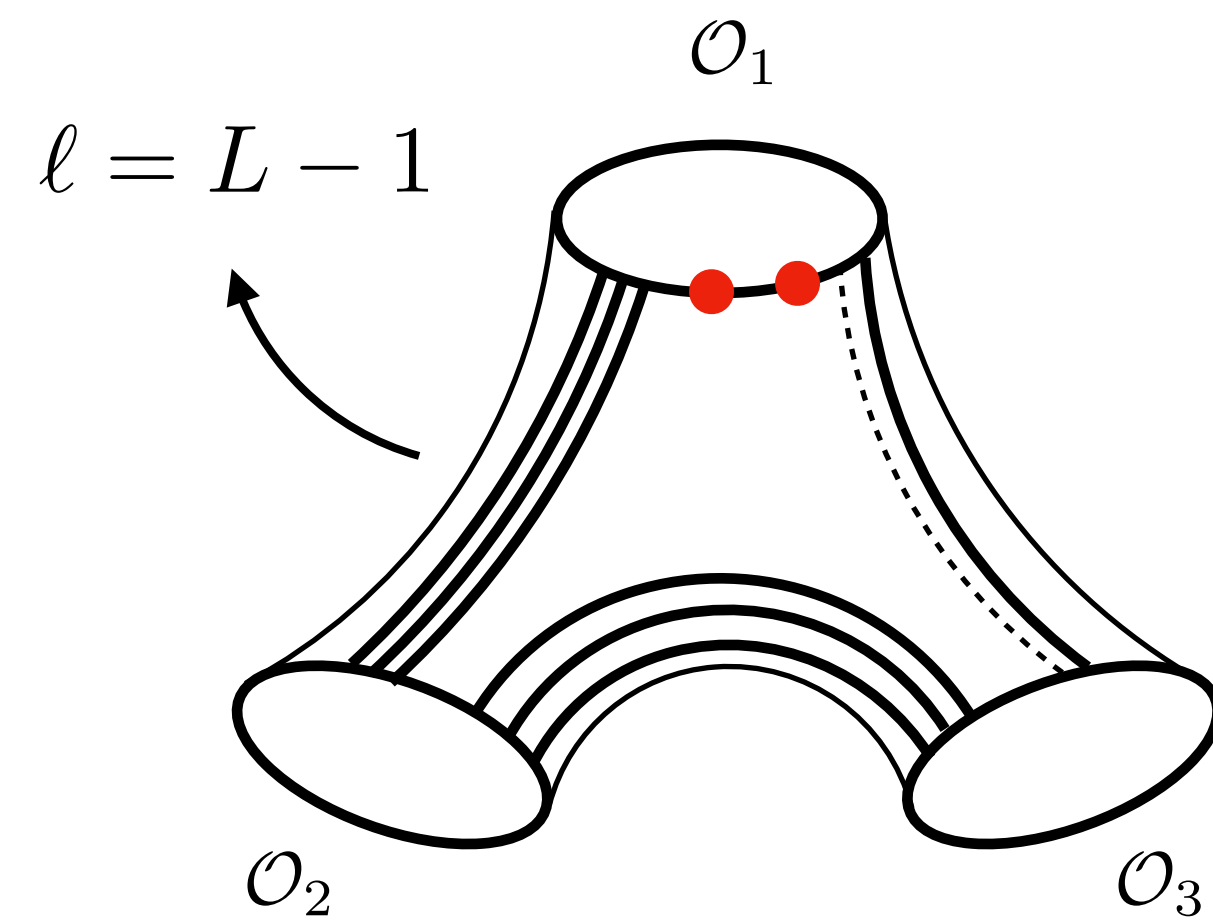
Finite Size effect in SU(2)

- Selection rules are NOT realized term by term in low energy expansion.
 - ↳ At N²LO

Finite Size effect in SU(2)

- Selection rules are NOT realized term by term in low energy expansion.

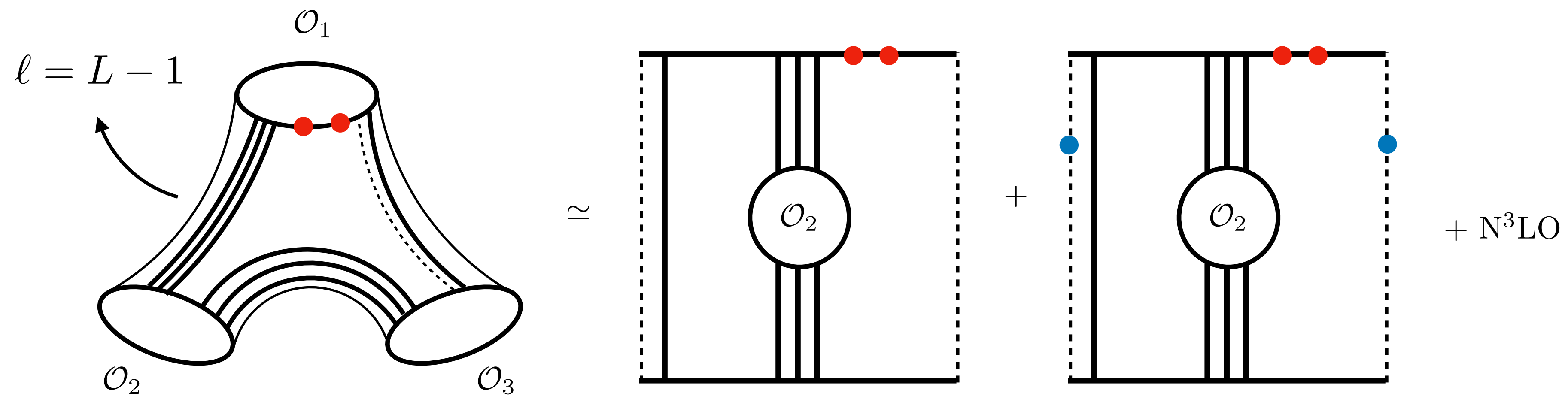
↳ At N²LO



Finite Size effect in SU(2)

- Selection rules are NOT realized term by term in low energy expansion.

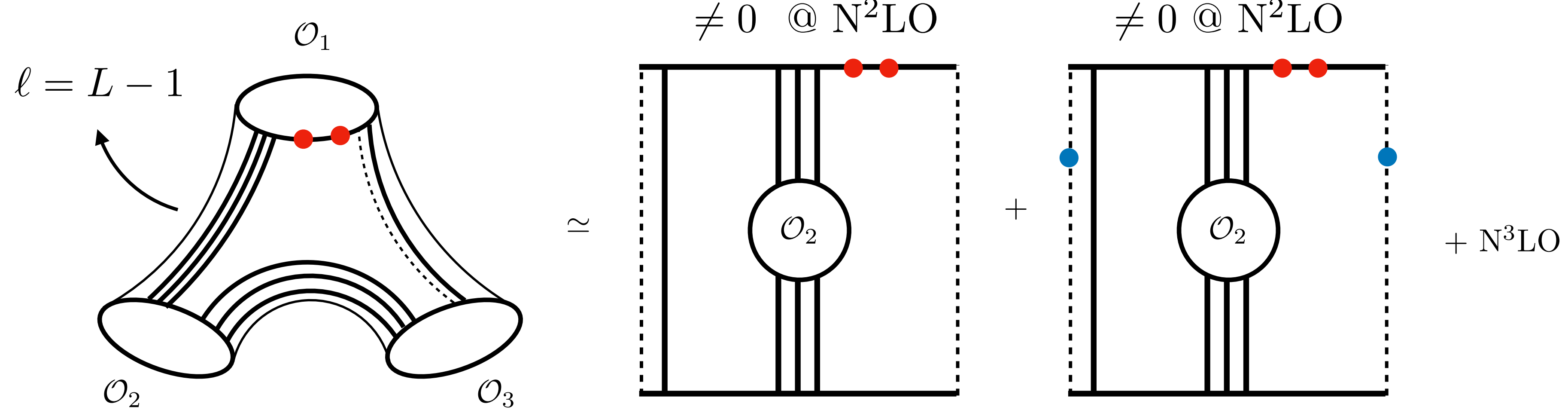
↳ At N²LO



Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

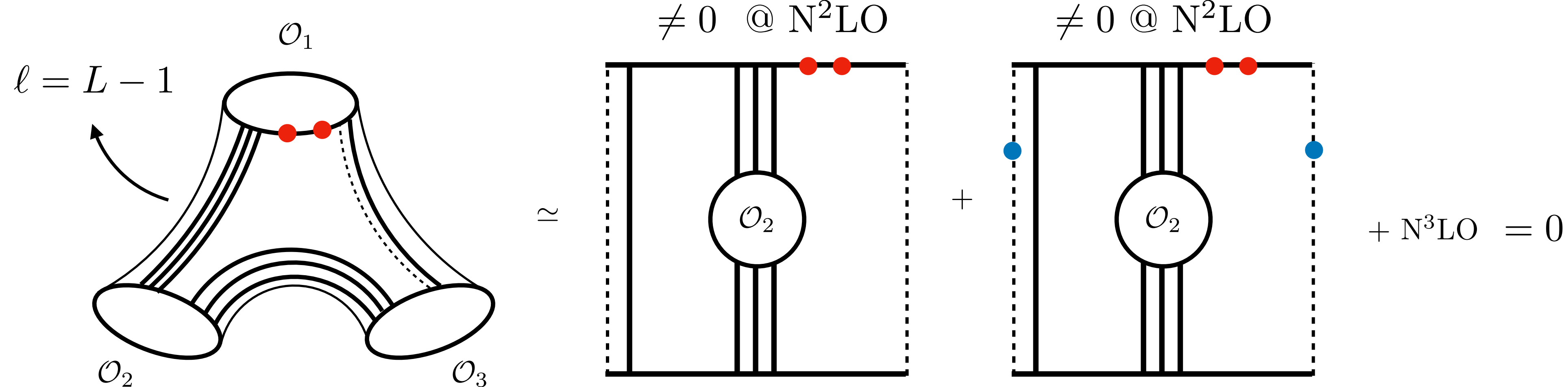
↳ At N^2LO



Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

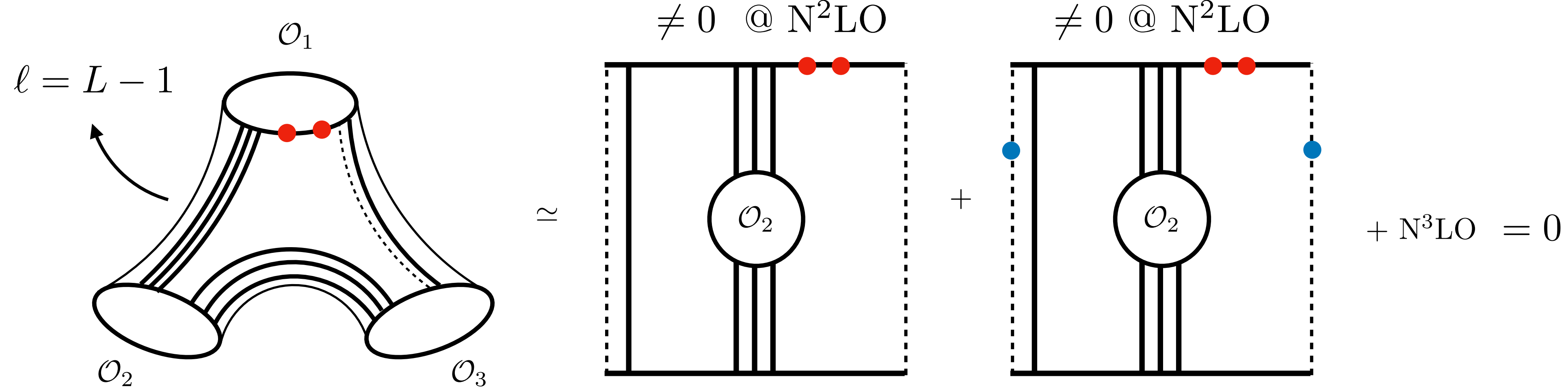
↳ At N^2LO



Finite Size effect in SU(2)

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N²LO

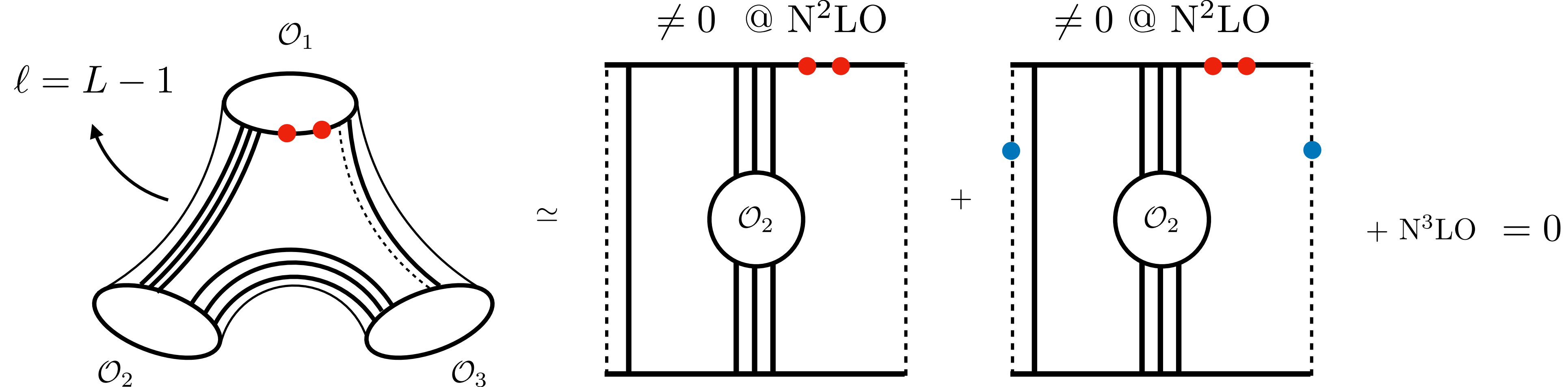


- Only with finite size correction are selection rules satisfied.

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N^2LO

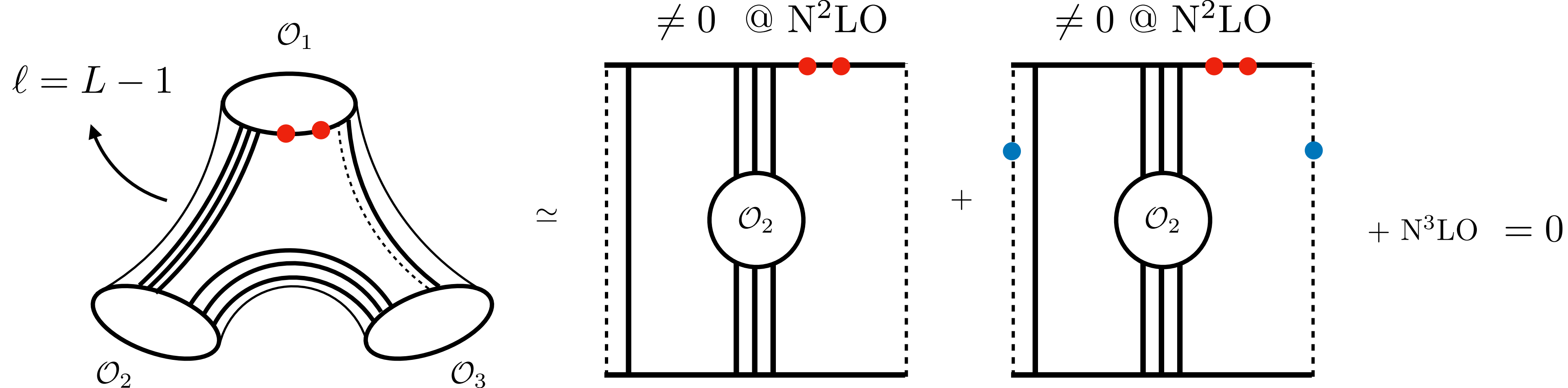


- Only with finite size correction are selection rules satisfied.
- Note this effect is not about wrapping corrections of the NBPS operator.

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N^2LO

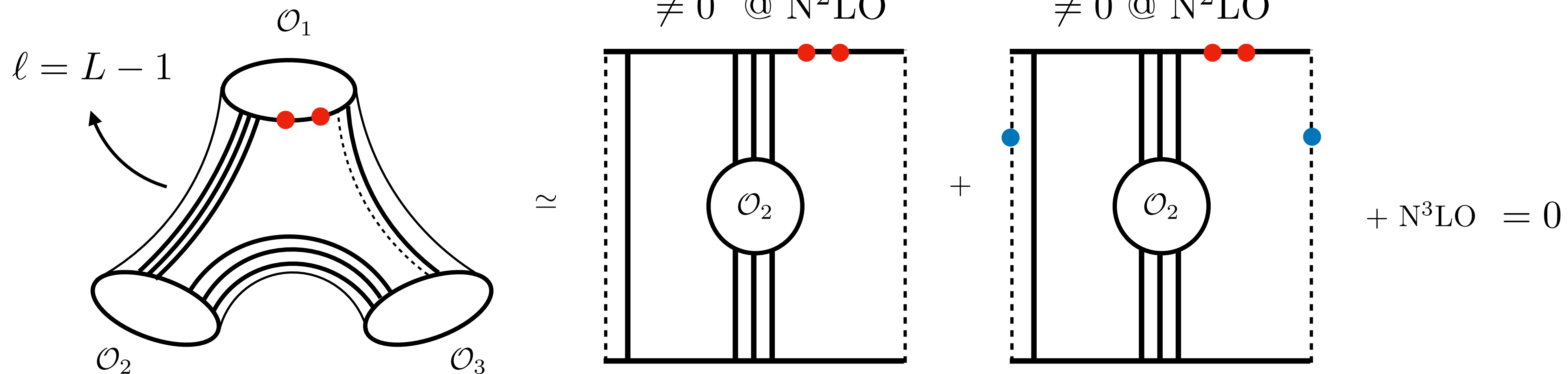


- Only with finite size correction are selection rules satisfied.
- Note this effect is not about wrapping corrections of the NBPS operator.
 ↳ Present even if all operators are very big. True “three point geometry” effect.

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N^2LO

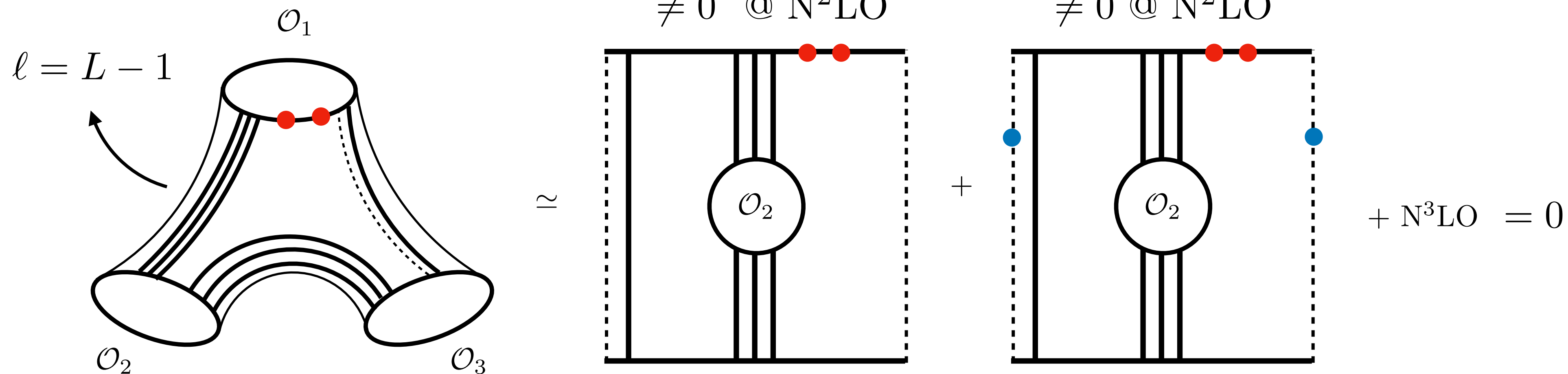


- Only with finite size correction are selection rules satisfied.
- Note this effect is not about wrapping corrections of the NBPS operator.
 ↳ Present even if all operators are very big. True “three point geometry” effect.
- SOV expression satisfy selection rule!

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N^2LO

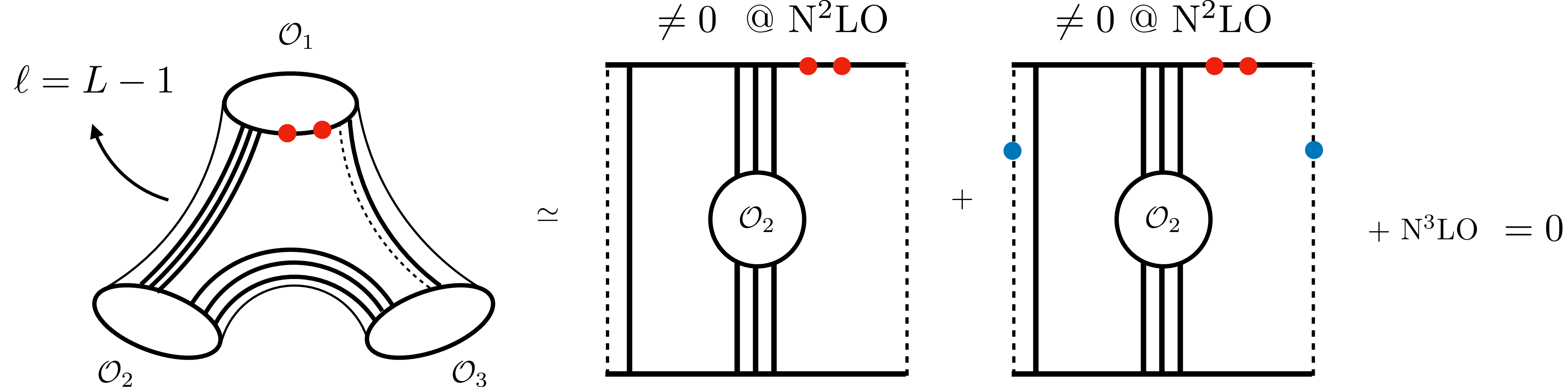


- Only with finite size correction are selection rules satisfied.
- Note this effect is not about wrapping corrections of the NBPS operator.
 ↳ Present even if all operators are very big. True “three point geometry” effect.
- SOV expression satisfy selection rule!
 ↳ For $\ell = L - 1$ new terms in the dressing of \mathcal{Q} .

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N^2LO

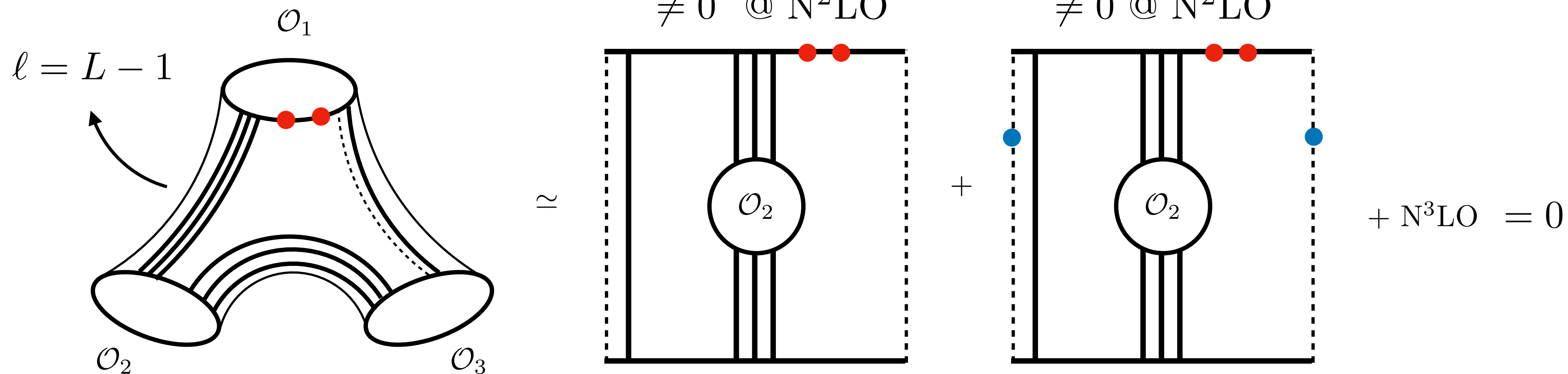


- Only with finite size correction are selection rules satisfied.
- Note this effect is not about wrapping corrections of the NBPS operator.
 ↳ Present even if all operators are very big. True “three point geometry” effect.
- SOV expression satisfy selection rule!
 ↳ For $\ell = L - 1$ new terms in the dressing of \mathcal{Q} . Extra terms are one-to-one with mirror correction.

Finite Size effect in $SU(2)$

- Selection rules are NOT realized term by term in low energy expansion.

↳ At N^2LO



- Only with finite size correction are selection rules satisfied.
- Note this effect is not about wrapping corrections of the NBPS operator.
 - ↳ Present even if all operators are very big. True “three point geometry” effect.
- SOV expression satisfy selection rule!
 - ↳ For $\ell = L - 1$ new terms in the dressing of \mathcal{Q} . Extra terms are one-to-one with mirror correction.
 - ↳ “Finite geometry dress \mathcal{Q} function”.

Conclusions

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?
- Can we derive the perturbative SoV formulas from the hexagon formalism?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?
- Can we derive the perturbative SoV formulas from the hexagon formalism? [Jiang, Komatsu, Kostov, Serban]

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?
- Can we derive the perturbative SoV formulas from the hexagon formalism? [Jiang, Komatsu, Kostov, Serban]
- Can we access a new class of observables in $\mathcal{N} = 4$ SYM through SOV?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?
- Can we derive the perturbative SoV formulas from the hexagon formalism? [Jiang, Komatsu, Kostov, Serban]
- Can we access a new class of observables in $\mathcal{N} = 4$ SYM through SOV?
 - ↳ Baxter/SoV is better suited to continuation in the charges than ABA

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?
- Can we derive the perturbative SoV formulas from the hexagon formalism? [Jiang, Komatsu, Kostov, Serban]
- Can we access a new class of observables in $\mathcal{N} = 4$ SYM through SOV?
 - ↳ Baxter/SoV is better suited to continuation in the charges than ABA
 - ↳ Correlation functions of light-ray operators from integrability?

Conclusions

- We presented N^2LO formulas for C_J in terms of Baxter's Q -functions.
- We match CFT data including adjacent and bottom finite size corrections.
- We fail to identify overarching principle.
 - ↳ Is there a set of axioms governing all formulas presented?
 - ↳ Can we bootstrap them?
 - ↳ Is there a systematic approach to construct these formulas beyond N^2LO ?
- Can we obtain a covariant formulation of the result?
 - ↳ Need to generalize SoV to the higher-rank supersymmetric spin chains?
- What is the connection between our dressed Q 's and the QSC variables?
- Can we derive the perturbative SoV formulas from the hexagon formalism? [Jiang, Komatsu, Kostov, Serban]
- Can we access a new class of observables in $\mathcal{N} = 4$ SYM through SOV?
 - ↳ Baxter/SoV is better suited to continuation in the charges than ABA
 - ↳ Correlation functions of light-ray operators from integrability? [WIP with Simmons-Duffin, Vieira]

SL(2) N²LO

SL(2) N²LO

- Before concluding, a brief result:

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.
- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 +$

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.
- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + g^4 \int \nu(u) \frac{Q(u)}{Q(i/2)} + \text{contact-terms}(Q_n)$

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + g^4 \int \nu(u) \frac{\mathbb{Q}(u)}{\mathbb{Q}(i/2)} + \text{contact-terms}(Q_n)$

$$\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$$

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + g^4 \int \nu(u) \frac{\mathbb{Q}(u)}{\mathbb{Q}(i/2)} + \text{contact-terms}(Q_n)$

$$\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$$

- Lesson: in general should expect more integrals at higher loops.

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + g^4 \int \nu(u) \frac{\mathbb{Q}(u)}{\mathbb{Q}(i/2)} + \text{contact-terms}(Q_n)$ $\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$

- Lesson: in general should expect more integrals at higher loops.

Finite d.o.f. spin-chain \longrightarrow ∞ d.o.f gauge-theory/strings
increase coupling

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + g^4 \int \nu(u) \frac{\mathbb{Q}(u)}{\mathbb{Q}(i/2)} + \text{contact-terms}(Q_n)$ $\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$

- Lesson: in general should expect more integrals at higher loops.

Finite d.o.f. spin-chain \longrightarrow ∞ d.o.f gauge-theory/strings
increase coupling

- So far no dependence on ℓ_{23} anywhere in the talk.

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + g^4 \int \nu(u) \frac{Q(u)}{Q(i/2)} + \text{contact-terms}(Q_n)$ $\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$

- Lesson: in general should expect more integrals at higher loops.

Finite d.o.f. spin-chain \longrightarrow ∞ d.o.f gauge-theory/strings
increase coupling

- So far no dependence on ℓ_{23} anywhere in the talk.
↳ At N²LO there are finite size effects for $\ell_{23} = 1$. \mathcal{A} must depend on ℓ_{23} .

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + 2g^4 \int \nu(u) \frac{Q(u)}{Q(i/2)} + \text{contact-terms}(Q_n)$ $\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$

- Lesson: in general should expect more integrals at higher loops.

Finite d.o.f. spin-chain \longrightarrow ∞ d.o.f gauge-theory/strings
increase coupling

- So far no dependence on ℓ_{23} anywhere in the talk.
↳ At N²LO there are finite size effects for $\ell_{23} = 1$. \mathcal{A} must depend on ℓ_{23} .

SL(2) N²LO

- Before concluding, a brief result:
- Recall: up to NLO, \mathcal{A} is an $\ell - 1$ dimensional integral. No integral for $\ell = 1$!
↳ Actually $\mathcal{A}_{\ell=1} = 1$ in our normalization.

- For $L = 2$, $\mathcal{A}_{\ell=1} = 1 + 2g^4 \int \nu(u) \frac{Q(u)}{Q(i/2)} + \text{contact-terms}(Q_n)$ $\nu(u) = \frac{i\pi - i(u+i/2)\pi^2 \tanh(\pi u)}{2(u+i/2)^3 \cosh^2(\pi u)}$

- Lesson: in general should expect more integrals at higher loops.

Finite d.o.f. spin-chain \longrightarrow ∞ d.o.f gauge-theory/strings
increase coupling

- So far no dependence on ℓ_{23} anywhere in the talk.
↳ At N²LO there are finite size effects for $\ell_{23} = 1$. \mathcal{A} must depend on ℓ_{23} .

Formulas

$$\mathcal{Q}(u) \equiv \left(\prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} \right) e^{\frac{g^2}{2} Q_1^+ H_1^+(u) + \frac{g^2}{2} Q_1^- H_1^-(u) + O(g^4)}$$

$$\mathcal{Q}(u) \equiv \prod_{k=1}^J \frac{u - v_k}{\sqrt{x_k^+ x_k^-}} \times e^{\frac{1}{2} g^2 Q_1^+ H_1^+ + \frac{1}{8} g^4 Q_1^+ Q_2^+ H_1^+ - \frac{1}{2} g^4 Q_1^+ H_3^+}$$

and

$$\begin{aligned} \mu_1(u) &= \frac{\pi}{2c_u^2} \left(1 + \pi^2 g^2 (3t_u^2 - 1) + \pi^4 g^4 \left(\frac{5}{6} - 7t_u^2 + \frac{11}{2} t_u^4 \right) \right. \\ &\quad \left. - \frac{g^4}{8} H_1^+ (Q_1^+(v_1) Q_2^+(v_2) + Q_1^+(v_2) Q_2^+(v_1)) \right), \end{aligned} \quad (14)$$

$$\begin{aligned} \mathbf{A}_{\ell,L}(u) &= g^2 (q_1^-)^2 + g^4 \left(\frac{1}{2} (q_2^+)^2 + 4\alpha (q_1^-)^2 - 6\pi^2 q_2^+ \delta_{\ell=2} \right), \\ \mathbf{A}_{\ell,L}(u, v) &= g^2 q_1^- \tilde{q}_1^- + g^4 \left(\frac{1}{2} q_2^+ \tilde{q}_2^+ + 2\alpha q_2^+ + 4\alpha q_1^- \tilde{q}_1^- \right), \\ \mathbf{B}(u) &= g^2 q_1^- Q_1^- \\ &\quad + g^4 \left(\frac{1}{2} q_2^+ Q_2^+ + \alpha Q_2^+ + 4\alpha q_1^- Q_1^- - \pi^2 q_1^- Q_1^- \delta_{\ell=2} + \right. \\ &\quad \left. + \left(\left(\frac{1}{8} (Q_1^+)^2 - \frac{1}{4} Q_2^+ - Q_1^+ + \frac{3}{8} (Q_1^-)^2 \right) q_2^+ - \frac{1}{2} q_3^- Q_1^- \right) \delta_{\ell,L-1} \right), \end{aligned} \quad (25)$$

$$\begin{aligned} q_k^+(u) &= i^{k+0} (x^+(u))^{-k} + (-i)^{k+0} (x^-(u))^{-k}, \\ q_k^-(u) &= i^{k-1} (x^+(u))^{-k} + (-i)^{k-1} (x^-(u))^{-k}, \end{aligned}$$

$$Q_1^+ = \oint \frac{u du}{4\pi g^2 \sqrt{u^2 - 4g^2}} \log \frac{\mathcal{Q}(u + i/2)}{\mathcal{Q}(u - i/2)}$$

$$\begin{aligned} H_n^+(u) &\equiv H_n(-1/2 + iu) + H_n(-1/2 - iu), \\ iH_n^-(u) &\equiv H_n(-1/2 + iu) - H_n(-1/2 - iu). \end{aligned}$$

formalism counter parts. for $SU(2)$ it is simply

$$\mathcal{A}_\ell = \langle \mathcal{Q}, \mathbf{1} \rangle_\ell, \quad \mathcal{B} = \frac{(2\mathbb{J})!}{(\mathbb{J}!)^2} \langle \mathcal{Q}, \mathcal{Q} \rangle_L, \quad (A13)$$

and for $SU(2)$ it reads

$$\mathcal{A}_\ell = \Lambda_{\mathcal{A}} \langle \langle \mathcal{Q}, \mathbf{1} \rangle \rangle_{\ell,L}, \quad \mathcal{B} = \Lambda_{\mathcal{B}} \frac{(2J)!}{(J!)^2} \langle \langle \mathcal{Q}, \mathcal{Q} \rangle \rangle_{L,L}. \quad (A14)$$

where

$$\Lambda_{\mathcal{A}} = e^{g^4 (\alpha - \delta_{\ell=2} \pi^2) ((Q_1^-)^2 + Q_2^+)} \quad (A15)$$

$$\Lambda_{\mathcal{B}} = \prod_{i,j}^J \left(1 - \frac{g^2}{x^+(v_i) x^+(v_j)} \right) \prod_{i,j}^J \left(1 - \frac{g^2}{x^-(v_i) x^-(v_j)} \right)$$

$$\begin{aligned} \mathcal{M}_{\text{NNLO-b}} &= \frac{1}{2} (2(Q_1^+)^2 - iQ_2^- - Q_1^- Q_1^+) \partial_1^2 + iQ_1^+ \partial_1^3 + \\ &\quad \frac{1}{2} (Q_1^- Q_1^+ - (Q_1^+)^2) \partial_2^2 - \frac{1}{2} iQ_1^+ \partial_2^3 - \frac{i}{2} Q_1^+ \partial_1 \partial_2^2 + \\ &\quad \frac{1}{2} (iQ_2^- - (Q_1^+)^2) \partial_1 \partial_2 - (\partial_1 \leftrightarrow \partial_L, \partial_2 \leftrightarrow \partial_{L-1}). \end{aligned} \quad (C3)$$

$$\left(\sum_{\ell} C_{\ell}^{\bullet\bullet\bullet} \right)^2 = \frac{(\mathbb{J}_1 + \mathbb{J}_2)!^2}{(2\mathbb{J}_1)!(2\mathbb{J}_2)!} \frac{\langle \mathcal{Q}_1, \mathcal{Q}_2 \rangle_{\ell}^2}{\langle \mathcal{Q}_1, \mathcal{Q}_1 \rangle_{L_1} \langle \mathcal{Q}_2, \mathcal{Q}_2 \rangle_{L_2}}. \quad (26)$$

$$\begin{aligned} \mathcal{A}_\theta &= \sum_{\mathbf{u}=\alpha\cup\bar{\alpha}} (-1)^{|\bar{\alpha}|} \prod_{n=1}^{\ell} \prod_{i \in \bar{\alpha}} \frac{u_i - \theta_n + i/2}{u_i - \theta_n - i/2} \prod_{j \in \alpha} \frac{u_j - u_i + i}{u_j - u_i}, \\ \mathcal{B}_\theta &= \det \left[\partial_{u_i} \log \left(\prod_{n=1}^L \frac{u_j - \theta_n + i/2}{u_j - \theta_n - i/2} \prod_{k \neq j} \frac{u_j - u_k - i}{u_j - u_k + i} \right) \right] \\ &\quad \times \prod_{i < j} \frac{(u_i - u_j)^2}{1 + (u_i - u_j)^2}, \end{aligned} \quad (16)$$

$$\mu_1(u) = \frac{\pi}{2c_u^2} (1 + g^2 \pi^2 (3t_u^2 - 1) + \dots)$$

$$\mu_2(u, v) = \frac{\pi(u-v) s_{u-v}}{c_u c_v} (1 + g^2 \pi^2 ((t_u + t_v)^2 - \frac{4}{3}) + \dots)$$