Testing the Equivalence Principle with the Distortion of Time









Swiss Cosmology Days 2025 ETH Zurich, June 6th, 2025

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Density fluctuation

MATTER FIELD

GRAVITATIONAL POTENTIALS



Spatial component

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Velocity







Density fluctuation

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Spatial component

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Velocity





Y

Time component



Density fluctuation

MATTER FIELD

GRAVITATIONAL **POTENTIALS**



Spatial component

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Velocity



Relations in GR with standard CDM





Density fluctuation

MATTER FIELD

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Relations in GR with standard CDM





Density fluctuation



Spatial component

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Velocity



Relations in GR with standard CDM





Density fluctuation



Spatial component

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Velocity

Continuity

Relations in GR with standard CDM

Time component



Density fluctuation



Spatial component

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Continuity

Velocity

Relations in GR with standard CDM

Euler (Weak Equivalence Principle)

Υ

Time component



Density fluctuation



Spatial component

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Time component

Velocity



Galaxy clustering

Fluctuations in galaxy number counts

 $\Delta(z,\mathbf{n}) = b \delta_m - \frac{1}{\mathcal{H}} \partial_r (\mathbf{V} \cdot \mathbf{n})$

Matter density x galaxy bias

Redshift-space distortions (RSD)

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Galaxy clustering

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 $\Delta(z,\mathbf{n}) = b \delta_m - \frac{1}{\mathcal{H}} \partial_r (\mathbf{V} \cdot \mathbf{n})$

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Redshift-space distortions (RSD)

Two-point correlation function

 $\xi \equiv \langle \Delta(z, \mathbf{n}) \Delta(z', \mathbf{n}') \rangle$

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Gravity modifications



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SC, Grimm and Bonvin (2022) Bonvin & Pogosian (2022) SC, Wang, Dam, Bonvin, Pogosian (2024)



Gravity modifications



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Breaking of the WEP by DM









Gravity modifications



CAN WE DISTINGUISH BETWEEN THE TWO?

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Breaking of the WEP by DM







Gravity modifications



CAN WE DISTINGUISH BETWEEN THE TWO?



<u>Generalised Brans-Dicke</u> Universal coupling

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Breaking of the WEP by DM





Coupled quintessence DM-only coupling





Forecasts for SKA2





Fit with both models (galaxy clustering + CMB + weak lensing)



SC, Wang, Dam, Bonvin, Pogosian (2024)

Generate mock data with one type of modification (e.g. $\beta_1 = 0$, $\beta_2 = 1$)



Forecasts for SKA2





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SC, Wang, Dam, Bonvin, Pogosian (2024)

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Deus ex machina: gravitational redshift



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Deus ex machina: gravitational redshift



Gravity modifications



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Directly measurable null test



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Directly measurable null test



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= 1 WEP valid







Directly measurable null test



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- WEP valid







Directly measurable null test





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- WEP valid
- WEP broken $\neq 1$



No assumptions on

• Primordial power spectrum shape Background cosmological expansion • Time dependence of μ, Θ, Γ • Galaxy bias





Take-home message



Happy to chat live or at <u>sveva.castello@unige.ch</u> :)

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Gravitational redshift is very exciting!
Key probe of the equivalence principle
Necessary to distinguish between modified gravity and non-standard dark matter

Enea's talk: venturing to cluster scales



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Additional slides

MATTER FIELD

Modified Poisson $k^2 \Psi = -4\pi G a^2 \bar{\rho} \,\delta \,\mu(k,z)$

GRAVITATIONAL **POTENTIALS**

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Continuity $\delta' + \frac{k}{\mathcal{H}}V - 3\Phi' = 0$

> What happens if DM violates the WEP?

$$V'_{\rm DM} + V_{\rm DM} - \frac{k}{\mathscr{H}} \Psi = 1$$

SC, Grimm and Bonvin (2022)

Φ

V

 $\underline{} = i$





Gravity modifications



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SC, Grimm and Bonvin (2022) Bonvin & Pogosian (2022)

Breaking of the WEP by DM



Two-point correlation function

Extract information through correlations:

$$\xi \equiv \left< \Delta(\mathbf{n}, z) \Delta(\mathbf{n}', z') \right>$$

 $\begin{array}{l} \longrightarrow & \text{Expansion in Legendre polynomials:} \\ \text{With } \Delta = \delta + \text{RSD}, & \overset{\text{Kaiser (1987)}}{\text{Hamilton (1992)}} \\ \xi = C_0(z, d) P_0(\cos \beta) & \text{Monopole} \\ + C_2(z, d) P_2(\cos \beta) & \text{Quadrupole} \\ + C_4(z, d) P_4(\cos \beta) & \text{Hexadecapole} \\ \end{array}$

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Credits: M.Blanton, SDSS

Relation with gravity modifications

 $C_0(z,d) = \tilde{b}^2(z) + \frac{2}{3}\tilde{b}(z)$ <u>Monopole</u>

 $C_2(z,d) = - \left| \frac{4}{3} \tilde{f}(z) \tilde{b}(z) + \right|$ Quadrupole

<u>Hexadecapole</u>

 $C_4(z,d) = \frac{8}{25}\tilde{f}^2(z)\,\mu_4(z_*,d)$





 $\tilde{f}(z) = f(z)\sigma_8(z)$ and $\tilde{b}(z) = b(z)\sigma_8(z)$

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$$\tilde{f}(z) + \frac{1}{5}\tilde{f}^2(z) \bigg] \mu_0(z_*, d)$$

$$\left[-\frac{4}{7}\tilde{f}^{2}(z)\right]\mu_{2}(z_{*},d)$$



constrained by CMB







$$\Delta(\mathbf{n},z) = b \delta_m - \frac{1}{\mathcal{H}} \partial_r (\mathbf{V} \cdot \mathbf{n})$$

Gravitational lensing

$$+ (5s - 2) \int_{0}^{r} dr' \frac{r - r'}{2rr'} \Delta_{\Omega}(\Phi + \Psi) \qquad \left\{ \begin{array}{l} \text{Subdominant} \\ + \left(\frac{5s - 2}{r \mathcal{H}} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - 5s + f^{\text{evol}} \right) \mathbf{V} \cdot \mathbf{n} + \mathbf{V} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \partial_{r} \Psi \\ + \frac{2 - 5s}{r} \int_{0}^{r} dr' (\Phi + \Psi) - (3 - f^{\text{evol}}) \mathcal{H} \nabla^{-2} (\nabla \mathbf{V}) + \Psi + (5s - 2) \Phi \\ + \frac{1}{\mathcal{H}} \dot{\Phi} + \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} + \frac{2 - 5s}{r \mathcal{H}} + 5s - f^{\text{evol}} \right) \left[\Psi + \int_{0}^{r} dr' (\dot{\Phi} + \dot{\Psi}) \right] \qquad \left\{ \begin{array}{c} \text{Subdom} \\ \text{Subdom} \end{array} \right\}$$

$$+ (5s - 2) \int_{0}^{r} dr' \frac{r - r'}{2rr'} \Delta_{\Omega}(\Phi + \Psi) \qquad \begin{cases} \text{Subdominant} \\ + \left(\frac{5s - 2}{r\mathcal{H}} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - 5s + f^{\text{evol}}\right) \mathbf{V} \cdot \mathbf{n} + \mathbf{V} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \partial_{r} \Psi \\ + \frac{2 - 5s}{r} \int_{0}^{r} dr' (\Phi + \Psi) - (3 - f^{\text{evol}}) \mathcal{H} \nabla^{-2} (\nabla \mathbf{V}) + \Psi + (5s - 2) \Phi \\ + \frac{1}{\mathcal{H}} \dot{\Phi} + \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} + \frac{2 - 5s}{r\mathcal{H}} + 5s - f^{\text{evol}}\right) \left[\Psi + \int_{0}^{r} dr' (\dot{\Phi} + \dot{\Psi})\right] \qquad \end{cases}$$
Subdom

Relativistic effects

$$+ (5s - 2) \int_{0}^{r} dr' \frac{r - r'}{2rr'} \Delta_{\Omega}(\Phi + \Psi) \qquad \begin{cases} \text{Subdominant} \\ + \left(\frac{5s - 2}{r\mathcal{H}} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - 5s + f^{\text{evol}}\right) \mathbf{V} \cdot \mathbf{n} + \mathbf{V} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \partial_{r} \Psi \\ + \frac{2 - 5s}{r} \int_{0}^{r} dr' (\Phi + \Psi) - (3 - f^{\text{evol}}) \mathcal{H} \nabla^{-2} (\nabla \mathbf{V}) + \Psi + (5s - 2) \Phi \\ + \frac{1}{\mathcal{H}} \dot{\Phi} + \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} + \frac{2 - 5s}{r\mathcal{H}} + 5s - f^{\text{evol}}\right) \left[\Psi + \int_{0}^{r} dr' (\dot{\Phi} + \dot{\Psi})\right] \qquad \end{cases}$$
Subdom

$$+ (5s - 2) \int_{0}^{r} dr' \frac{r - r'}{2rr'} \Delta_{\Omega}(\Phi + \Psi) \qquad \left\{ \begin{array}{l} \text{Subdominant} \\ + \left(\frac{5s - 2}{r\mathcal{H}} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} - 5s + f^{\text{evol}} \right) \mathbf{V} \cdot \mathbf{n} + \mathbf{V} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \frac{1}{\mathcal{H}} \partial_{r} \Psi \\ + \frac{2 - 5s}{r} \int_{0}^{r} dr' (\Phi + \Psi) - (3 - f^{\text{evol}}) \mathcal{H} \nabla^{-2} (\nabla \mathbf{V}) + \Psi + (5s - 2) \Phi \\ + \frac{1}{\mathcal{H}} \dot{\Phi} + \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^{2}} + \frac{2 - 5s}{r\mathcal{H}} + 5s - f^{\text{evol}} \right) \left[\Psi + \int_{0}^{r} dr' (\dot{\Phi} + \dot{\Psi}) \right] \qquad \left\{ \begin{array}{c} \text{Subdom} \\ \text{Subdom} \end{array} \right\}$$

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What we really observe

Yoo et al. (2010) Bonvin and Durrer (2011) Challinor and Lewis (2011) Jeong, Schmidt and Hirata (2012)



 $+ \left(\frac{5s - 2}{\mathcal{H}r} - \frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + f^{\text{evol}} \right) \mathbf{V} \cdot \mathbf{n}$



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What we really observe

Gravitational redshift

 $\Delta(\mathbf{n},z) = b \,\delta_m - \frac{1}{\mathscr{H}} \partial_r (\mathbf{V} \cdot \mathbf{n}) + \frac{1}{\mathscr{H}} \partial_r \Psi + \frac{1}{\mathscr{H}} \dot{\mathbf{V}} \cdot \mathbf{n} + \mathbf{V} \cdot \mathbf{n}$





Extracting the signal from observations

Relativistic effects break the symmetry of ξ

$$C_{1}(z,d) = \frac{\mathscr{H}}{\mathscr{H}_{0}} \nu_{1}(d,z_{*}) \left[5\tilde{f} \left(\tilde{b}_{B}s_{F} - \tilde{b}_{F}s_{B} \right) \left(1 - \frac{1}{2} - 3\tilde{f}^{2}\Delta s \left(1 - \frac{1}{r\mathscr{H}} \right) + \tilde{f}\Delta \tilde{b} \left(\frac{2}{r\mathscr{H}} + \frac{\dot{\mathscr{H}}}{\mathscr{H}^{2}} \right) \right) \right]$$
$$+ \Delta \tilde{b} \left(\Theta \quad \tilde{f} - \frac{3}{2} \frac{\Omega_{m,0}}{a} \frac{\mathscr{H}_{0}^{2}}{\mathscr{H}^{2}} \Gamma \mu \sigma_{8} \right) \left[-\frac{2}{5} \frac{2}{5} \left(\Theta - \tilde{f} - \frac{3}{2} \frac{\Omega_{m,0}}{a} \frac{\mathscr{H}_{0}^{2}}{\mathscr{H}^{2}} \right) \right]$$

Compare $\mu(\Gamma + 1)$ term in the evolution equation

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Symmetry breaking by gravitational redshift



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Reproduced from Bonvin, Hui and Gaztañaga (2014)

Survey specifications

 σ_{μ_0} (restricted to WEP va $\sigma_{\mu_0+\Gamma_0}$ (no assumption on

DESI (Bright Galaxy Sample):

- 10 million galaxies up to z=0.5.
- Galaxy bias: $b_{BGS}(z) = b_0 \delta(0) / \delta(z)$. $b_0 = 1.34$ (fiducial value)

Fisher analysis:

- minimum separation $d_{\min} = 20 \,\mathrm{Mpc}/h$.
- include shot noise, cosmic variance, cross-correlations between different multipoles

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	SDSS-IV	DESI	SKA2
alidity)	0.21	0.02	0.004
WEP)	6.05	0.42	0.068

SKA, phase 2: • ~1 billion galaxies up to z=2.0. • Galaxy bias: $b_{SKA}(z) = b_1 \exp(b_2 z)$. $b_1 = 0.554$, $b_2 = 0.783$ (fiducial value)



Deus ex machina: gravitational redshift



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SC, Grimm and Bonvin (2022)



Precision with SDSS data



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Modified gravity vs dark sector interactions

all constituents (μ, η)



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Bonvin and Pogosian (2022)



Forecasts for SKA2



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